

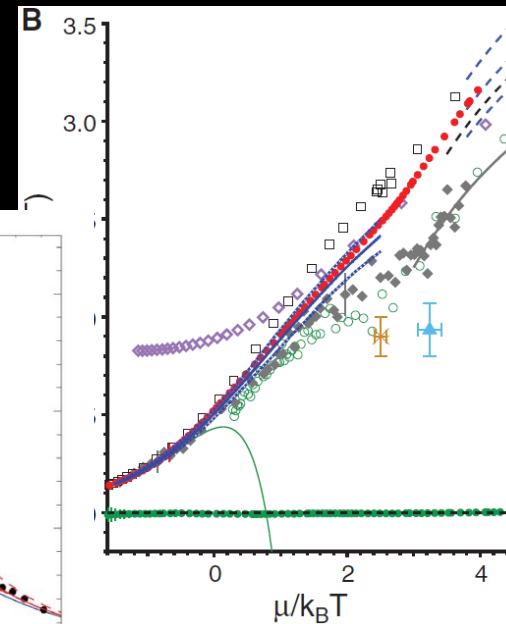
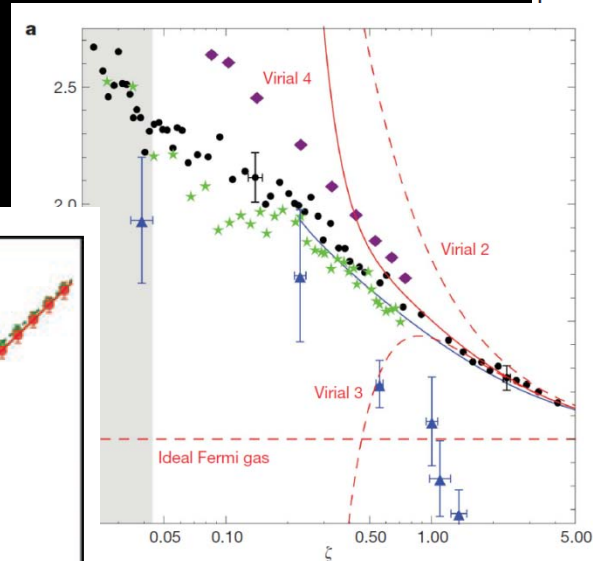
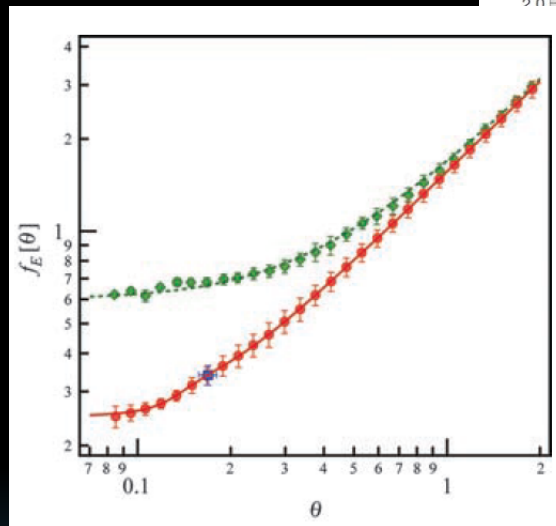
Measurement of p-wave interaction energy
in an ultracold spin-polarized Fermi gas of ${}^6\text{Li}$



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University of Electro-Communications

Experimental determination of Equation of State (EOS)

unitary Fermi gas (s-wave)



M. Ku, et al.
Science **335**, 563 (2012).

S. Nascimbene, et al.
Nature **463**, 1057 (2010).

M. Horikoshi, et al.
Science **327**, 442 (2010).

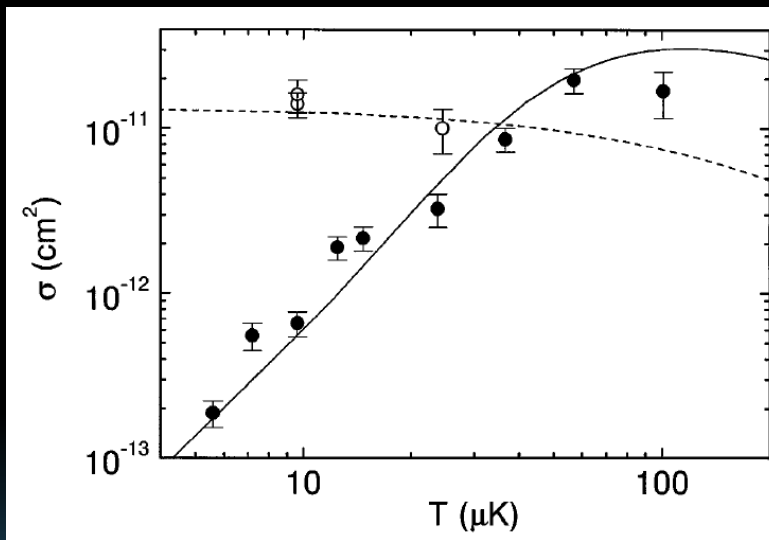
Our goal:
EOS for fermions with p-wave
interactions

Scattering of identical particles

Bosons : s -wave ($\ell=0$), d -wave ($\ell=2$), ...

Fermions : p -wave ($\ell=1$), f -wave ($\ell=3$), ...

Higher partial wave contribution decreases with temperature.



B. DeMarco et al.
Phys. Rev. Lett. 82, 4208 (1999).

○ s -wave collision cross section
between distinguishable fermions

● p -wave collision cross section
between identical fermions

Threshold energy for ${}^6\text{Li}$ is on the
order of mK.
(Typical temperature range $\sim 1\mu\text{K}$)

p -wave Feshbach resonance

- narrow Feshbach resonance

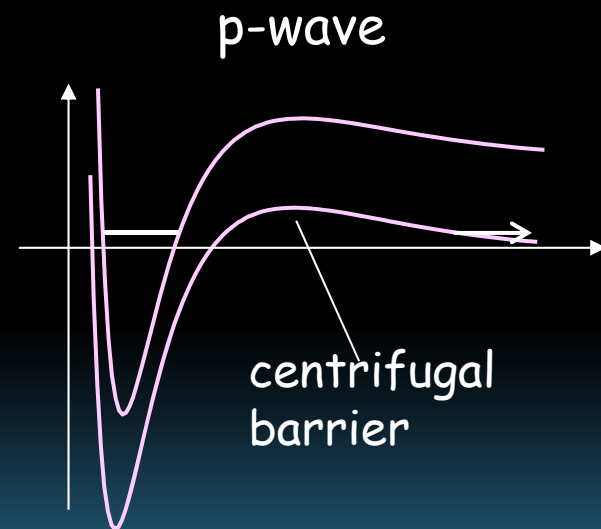
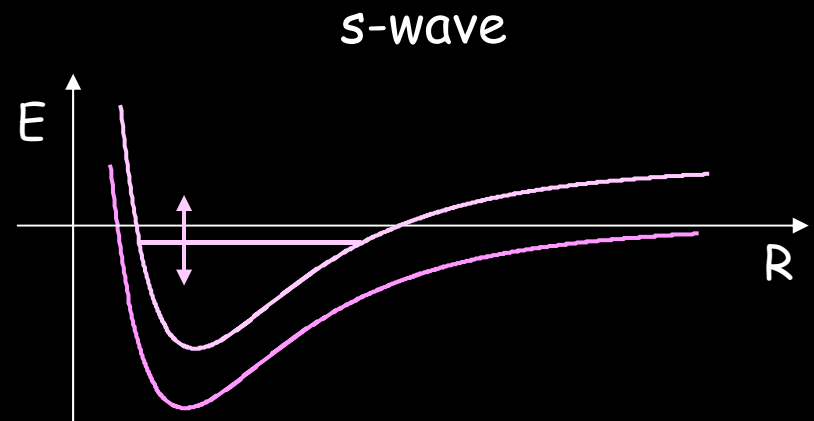
Atoms need to tunnel through the barrier

- quasi-bound state

molecular state has finite lifetime even at higher energy than the dissociation limit

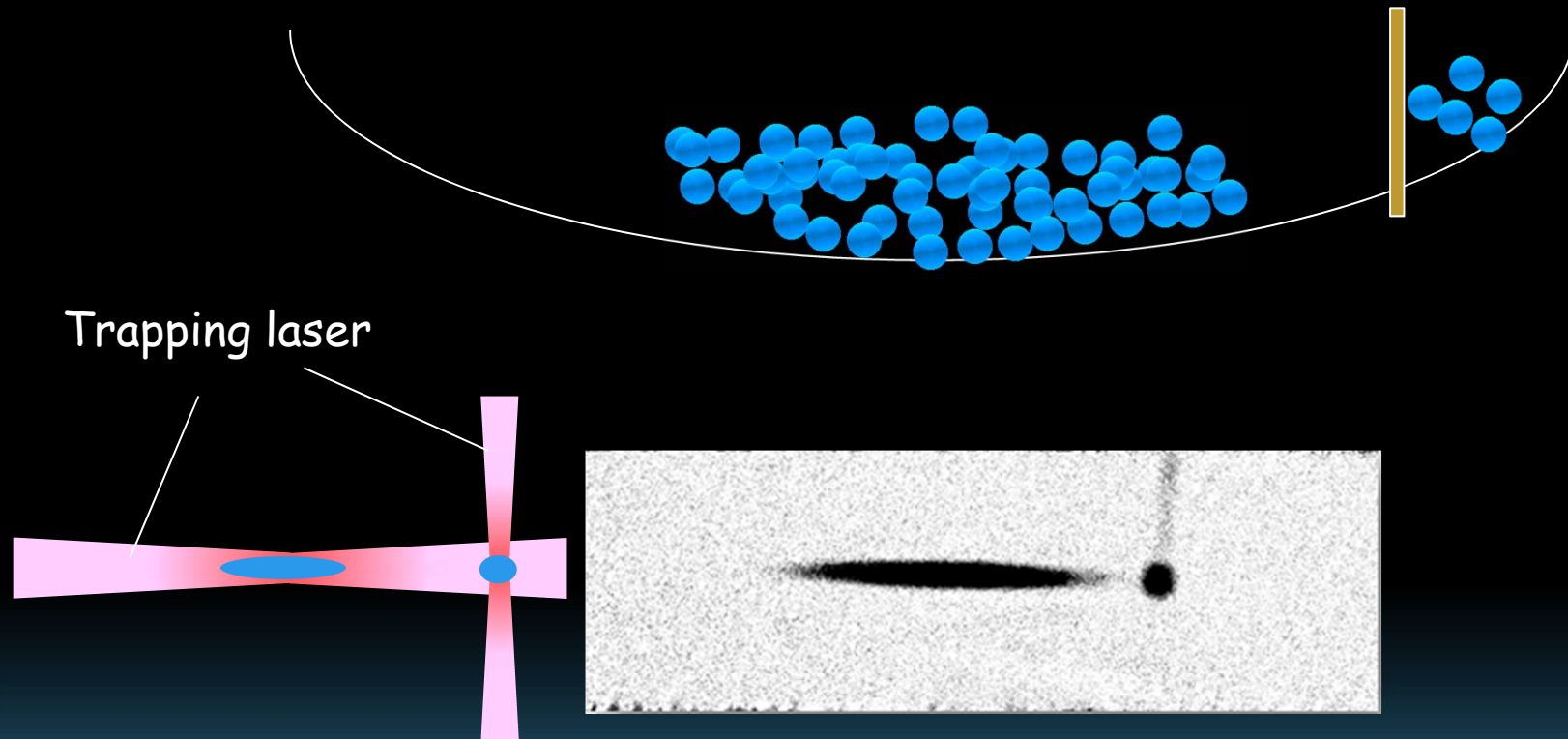
- doublet structure

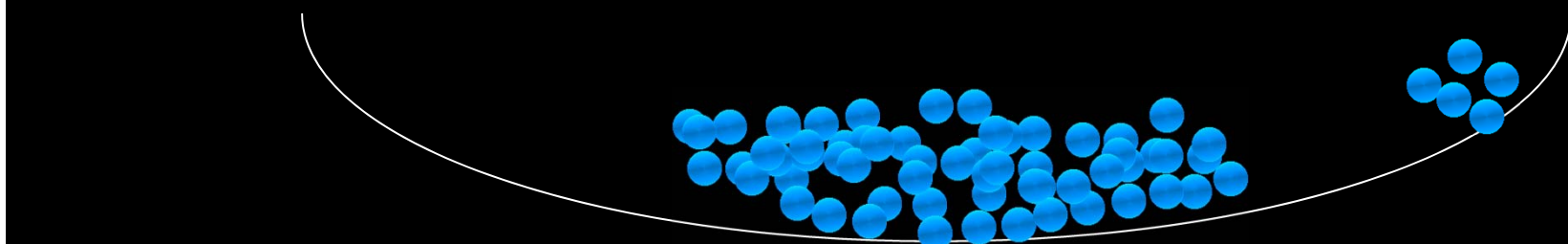
spin-spin interaction



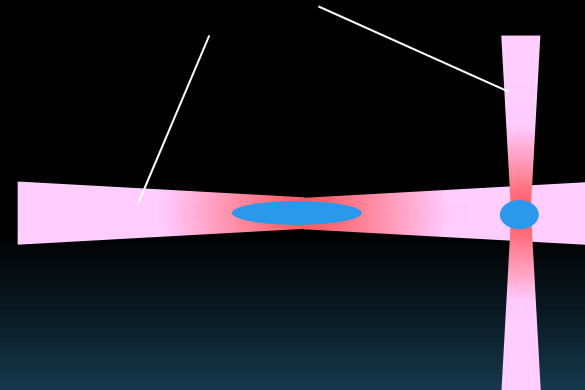
When the scattering length is tuned to nearly zero,

→ Atoms would not collide with each other

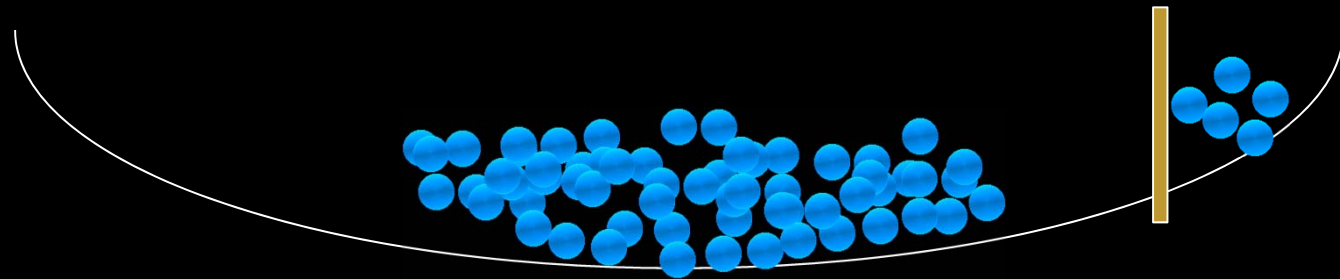




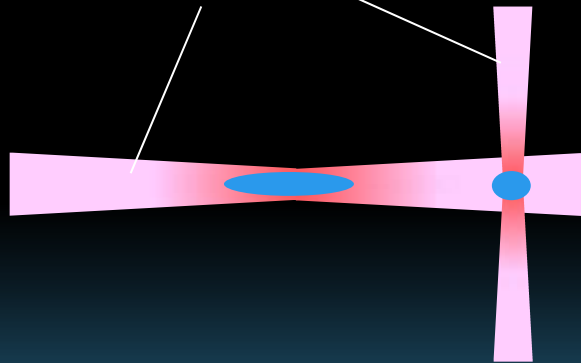
Trapping laser

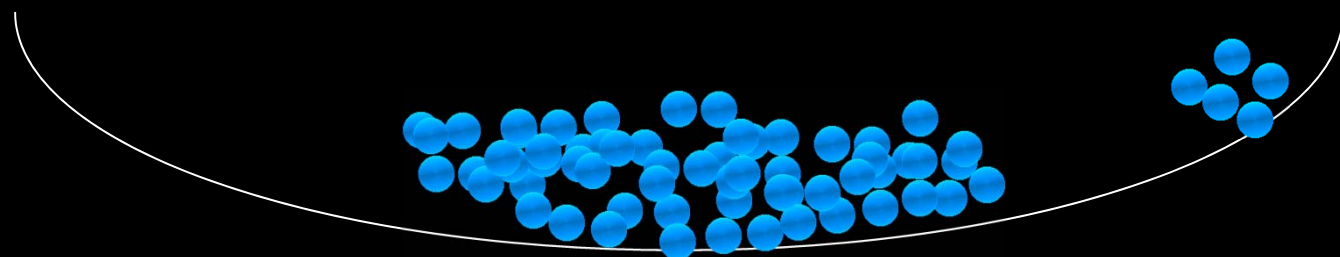


At the unitarity limit,

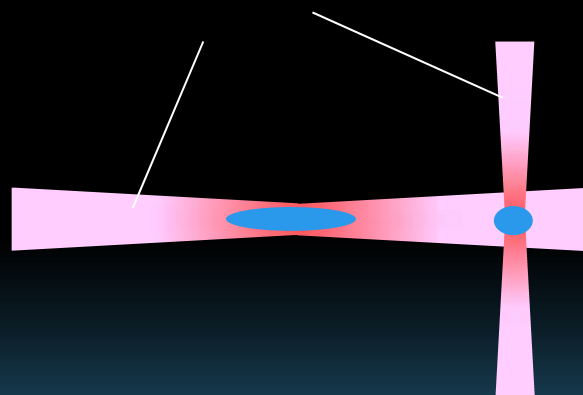


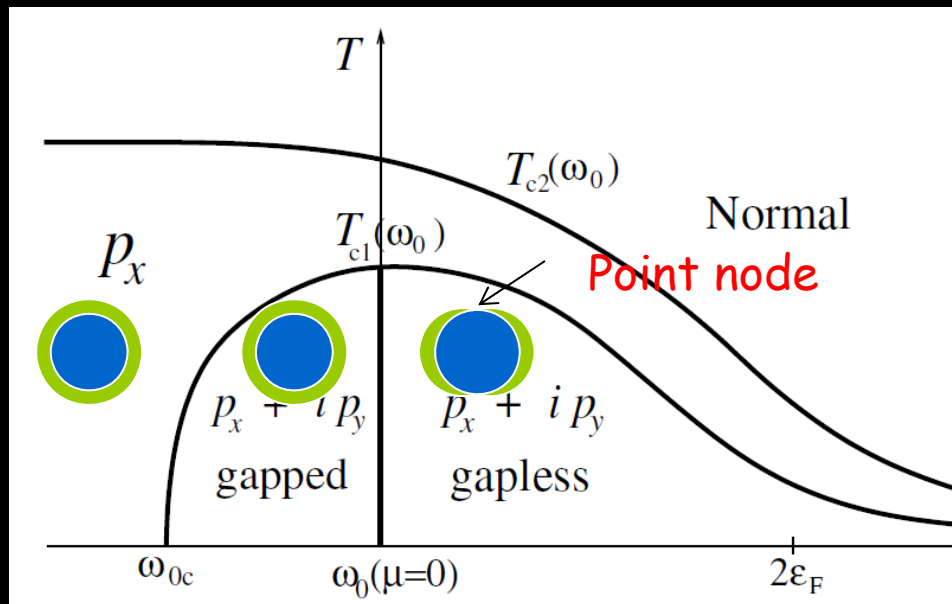
Trapping laser





トラップレーザー





Expected p-wave superfluid
phase diagram :

V. Gurarie et al.
PRL **94**, 230403 (2005)

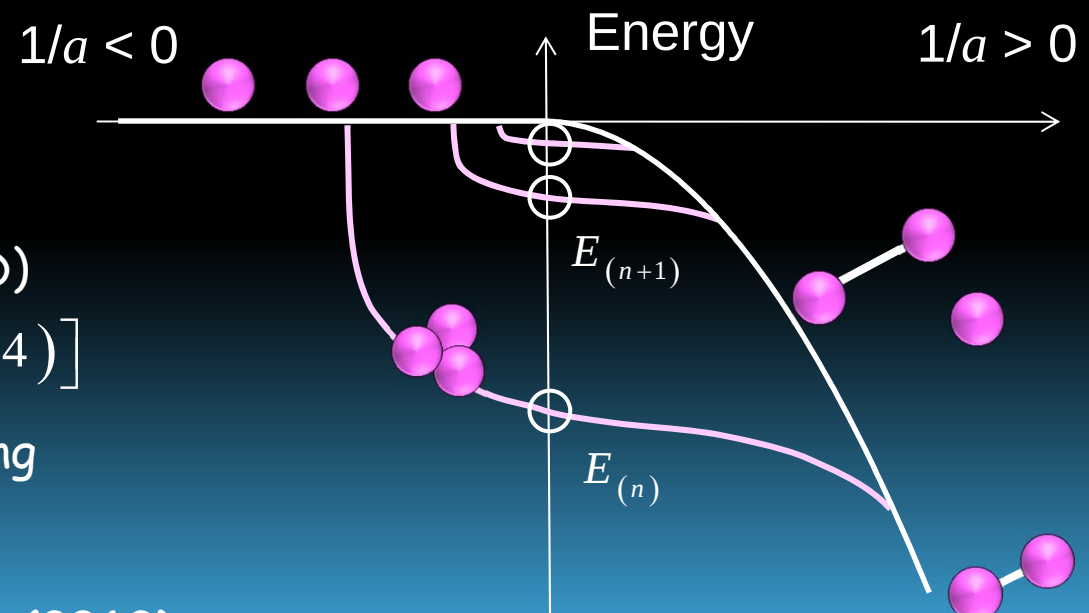
"Super" Efimov effect

p-wave resonance
in identical fermions (2D)

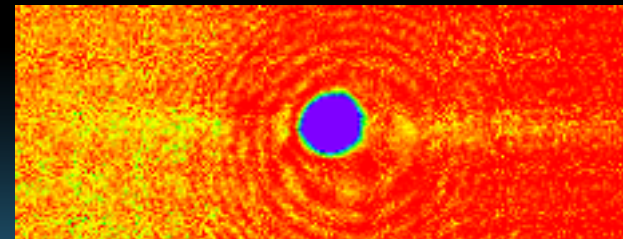
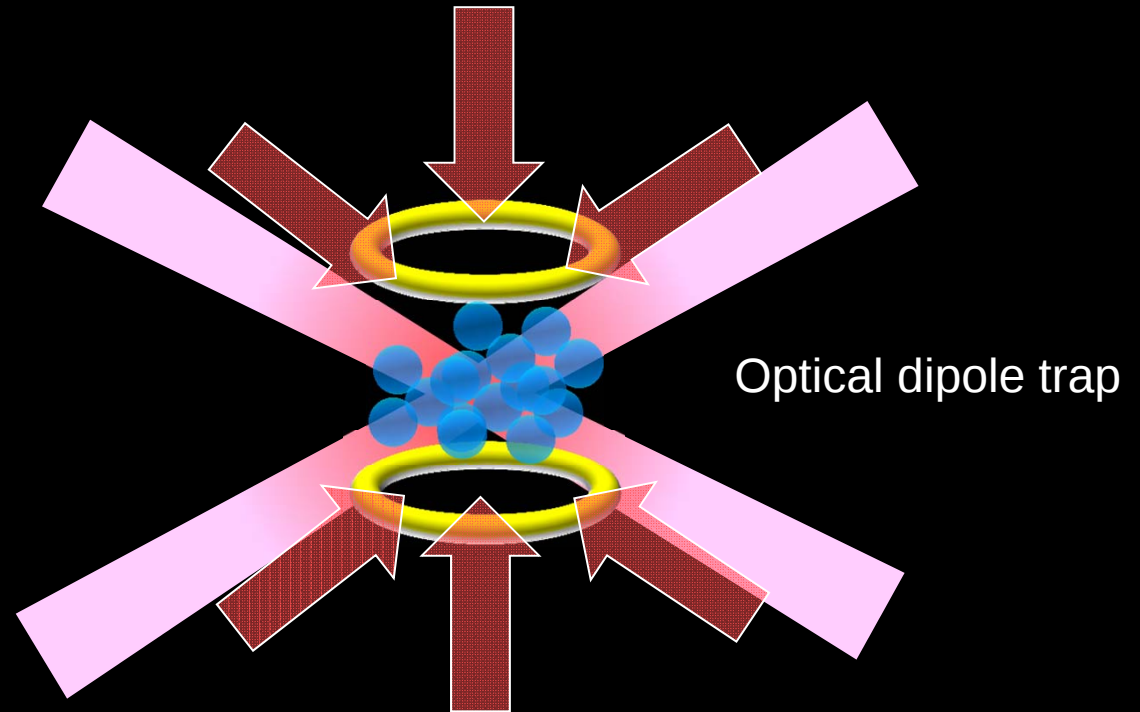
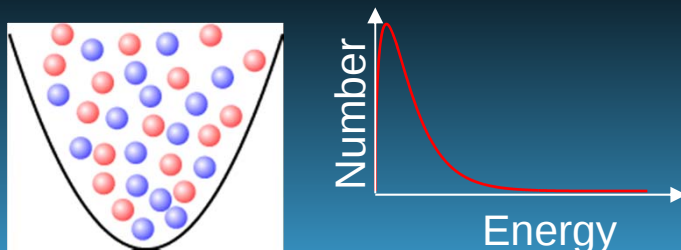
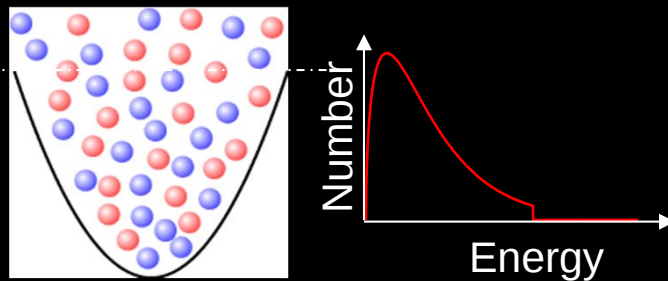
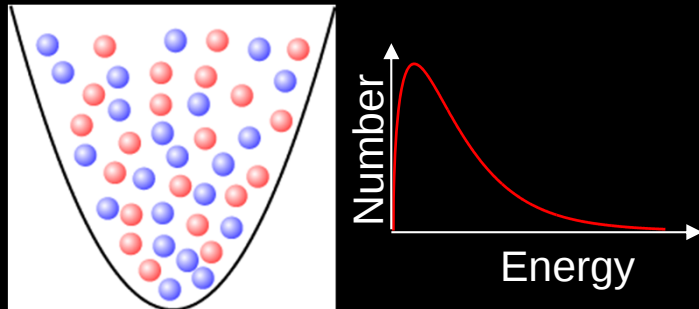
$$E_{(n)} \propto \exp \left[-A \exp \left(3\pi n/4 \right) \right]$$

doubly-exponential scaling

Y. Nishida et al.,
Phys. Rev. Lett. **110**, 235301 (2013)

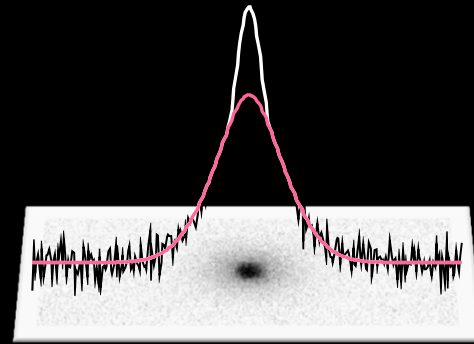


Evaporative cooling in an optical dipole trap



Size $\sim 10\text{-}100\mu\text{m}$
Aspect ratio $\sim 5:4:1$
Number of atoms $\sim 10^5\text{-}10^6$
Temperature $\sim 1\mu\text{K}$

${}^6\text{Li}$ in quantum degenerate regime

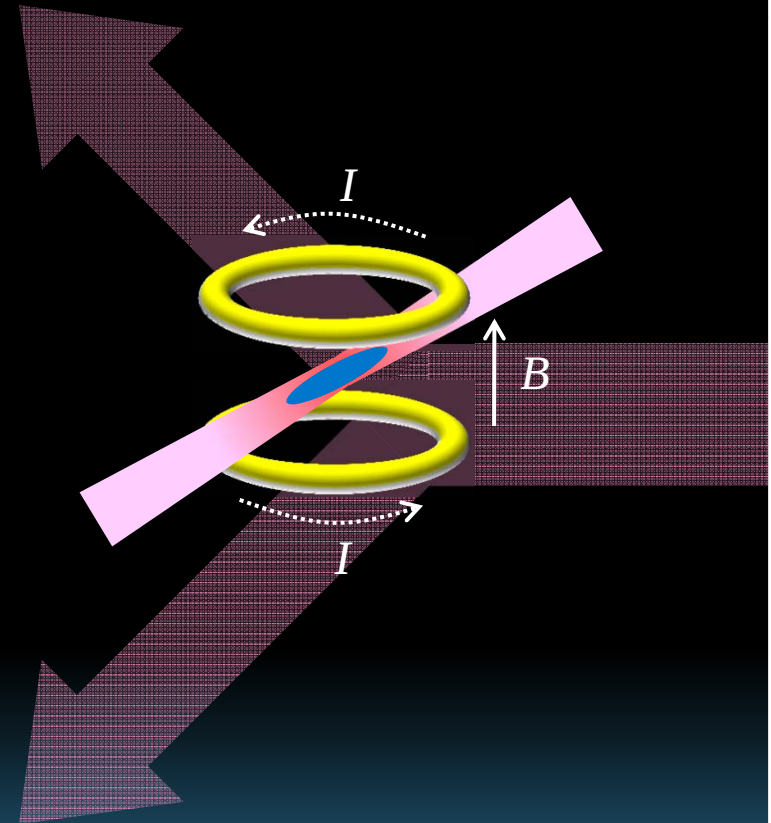
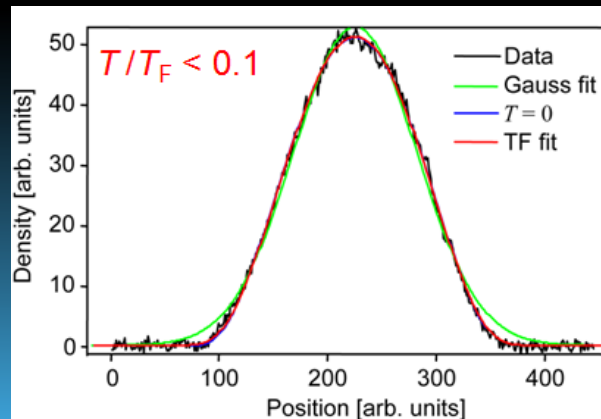
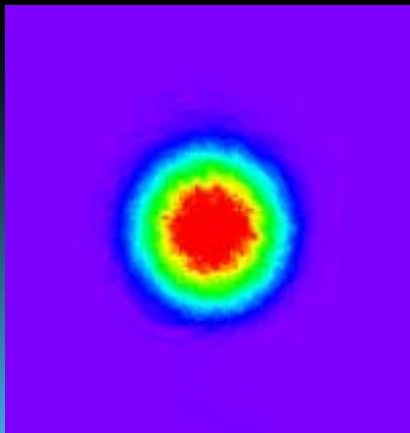


molecular BEC

$B = 650 \sim 800 \text{ G}$

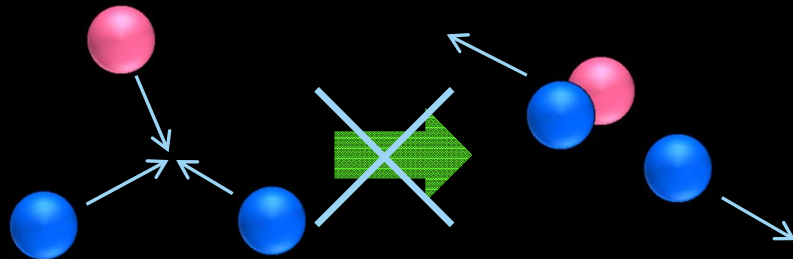
degenerate Fermi gas

$B = 0 \sim 500, 800 \sim \text{G}$



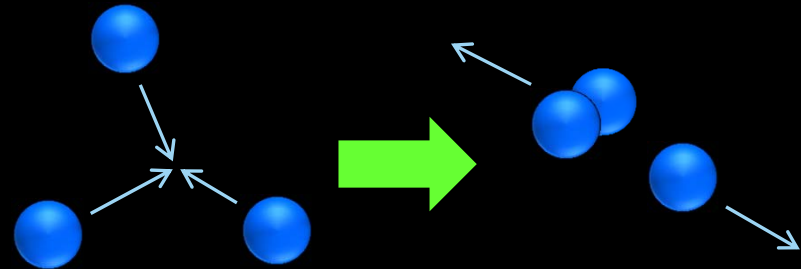
Interaction and loss are the opposite sides of the same coin...

s-wave case



two-component fermions with
s-wave interaction
→ **stable against 3-body loss**

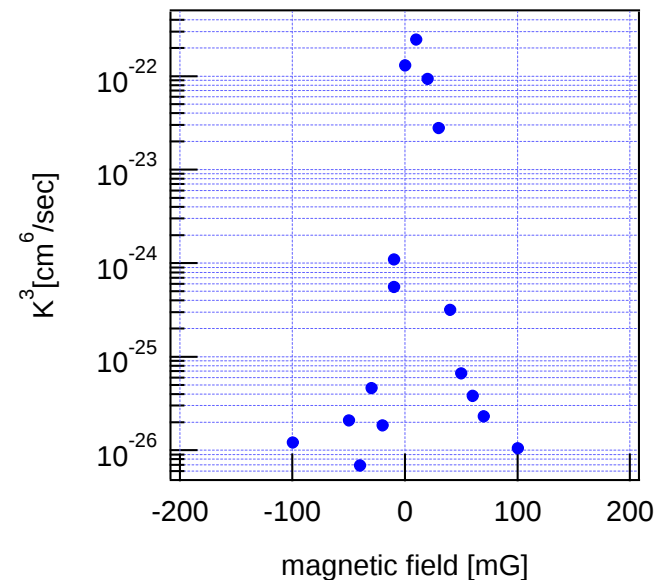
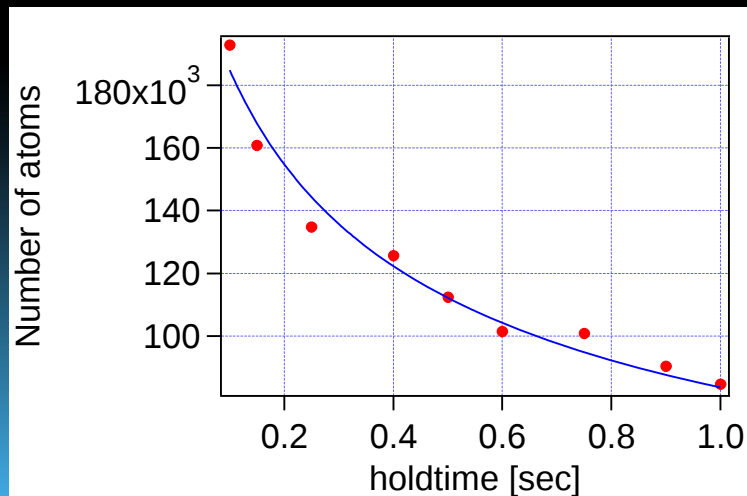
p-wave case



3-body loss is crucial

$$\dot{n} = -K_3 n^3$$

decay plot



Work done so far with p-wave Feshbach resonances

⁴⁰K

- ✓ p-wave Feshbach resonance [1]
- ✓ splitting between $m_f = 0$ and $|m_f| = 1$ resonances[2]
- ✓ p-wave molecule creation [3]

[1] C. A. Regal et. al. PRL **90**, 053201 (2003)

[2] C. Ticknor et. al. PRA **69**, 042712 (2004)

[3] J. P. Gaebler et. al. PRL **98**, 200403 (2007)

⁶Li

- ✓ observation of p-wave Feshbach resonances[4-6]
- ✓ "loss" and "recovery" of atoms near the p-wave Feshbach resonance[4]
- ✓ determination of the magnetic moment of p-wave molecular state[6]
- ✓ p-wave molecule creation [7,8]

[4] J. Zhang et al. PRA **70**, 030702(R) (2004)

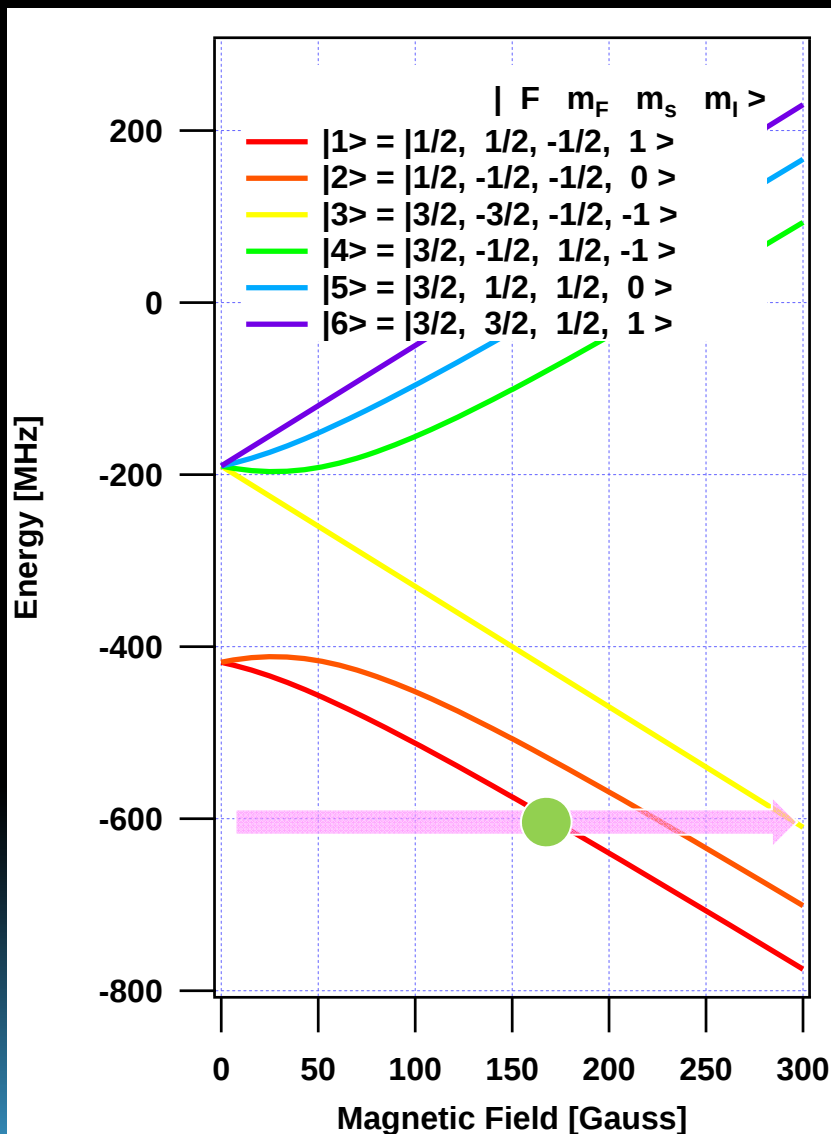
[5] C. H. Schunck et al. PRA **71**, 045601 (2005)

[6] J. Fuchs et al. PRA **77**, 053616 (2008)

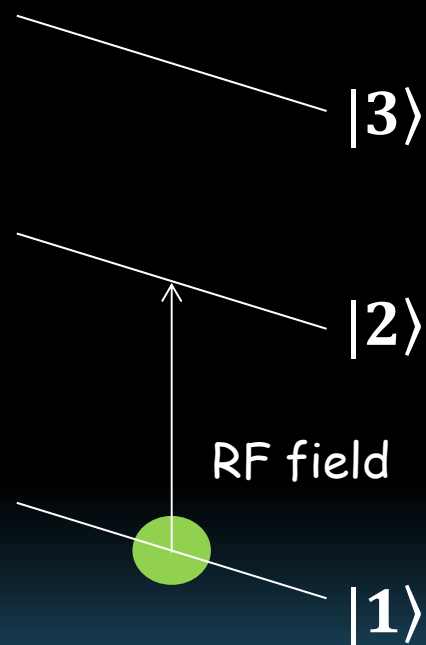
[7] Y. Inada et al. PRL **101**, 100401 (2008)

[8] R. A. W. Maier et al. PRA **81**, 064701 (2010)

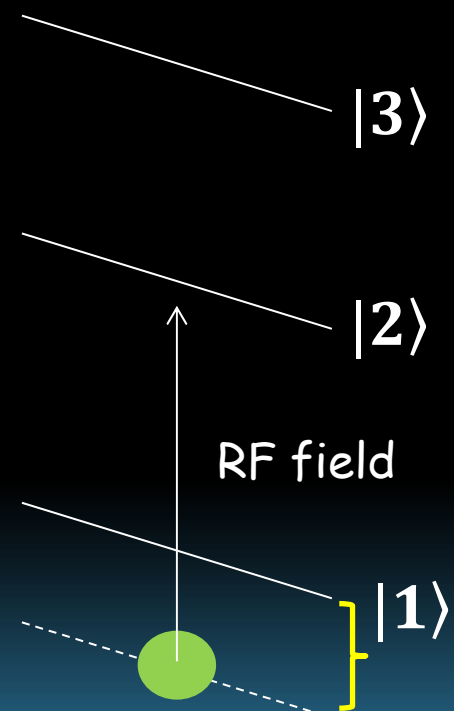
The way to probe interaction - RF spectroscopy



without
interaction



with the presence
of interaction



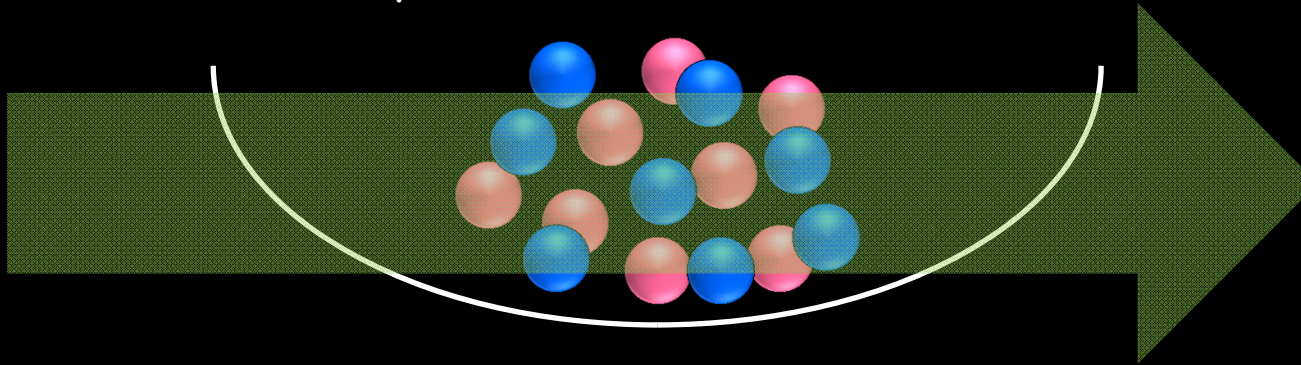
Energy shift due to
p-wave interaction

Stern-Gerlach measurement

Turn off the trap

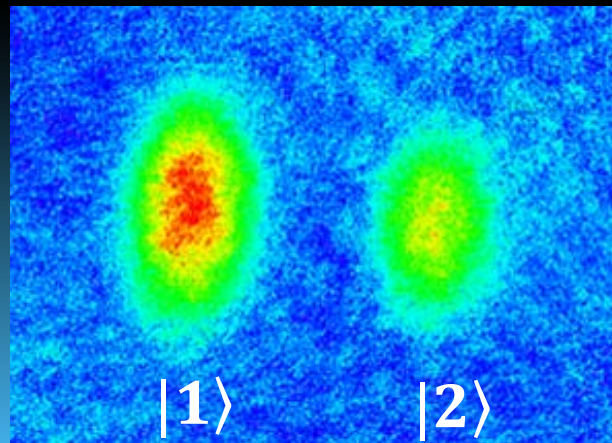


Apply magnetic field gradient



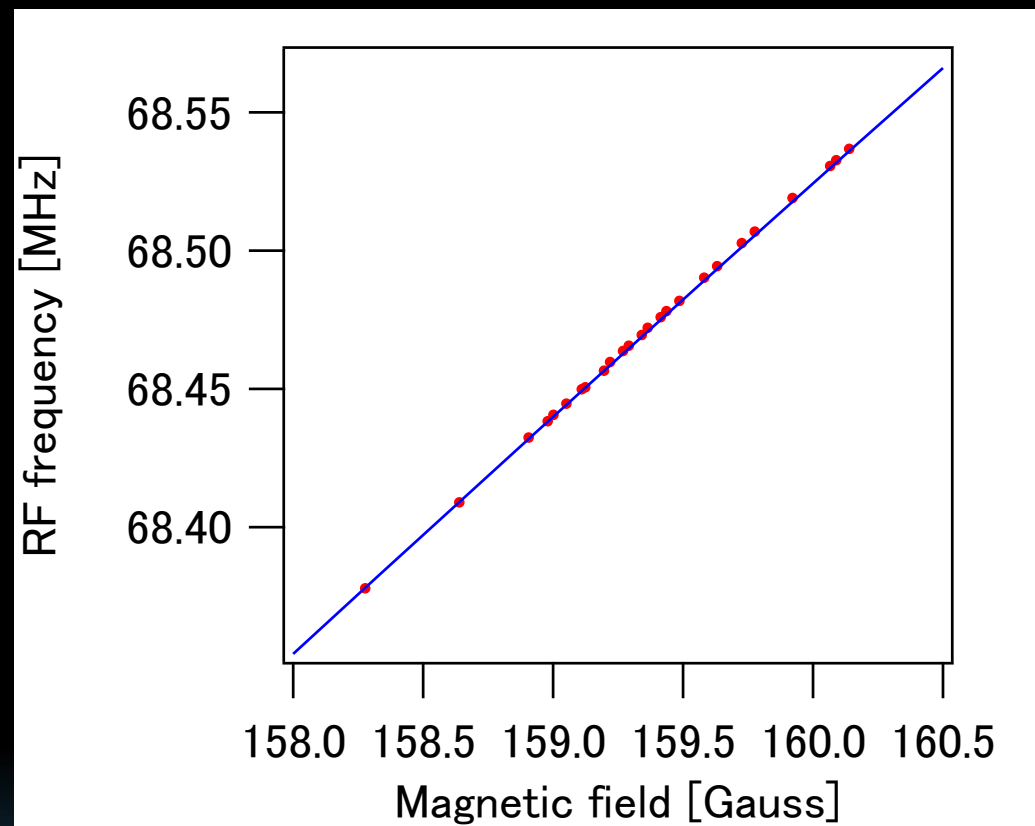
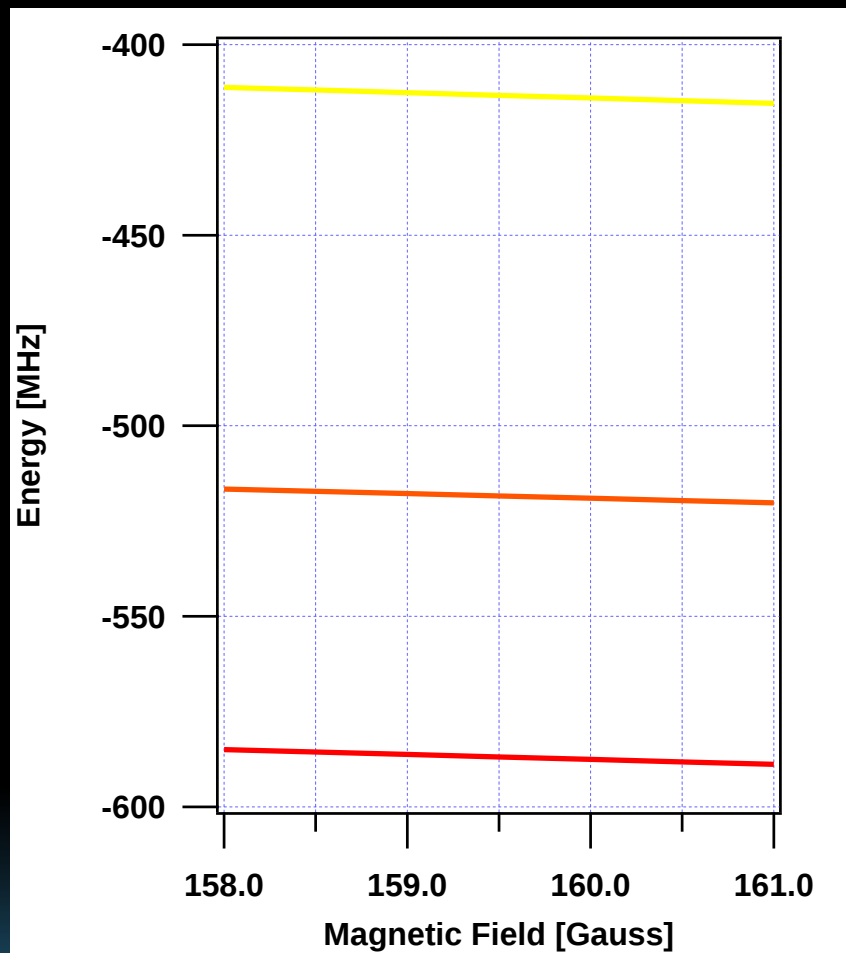
$|1\rangle$

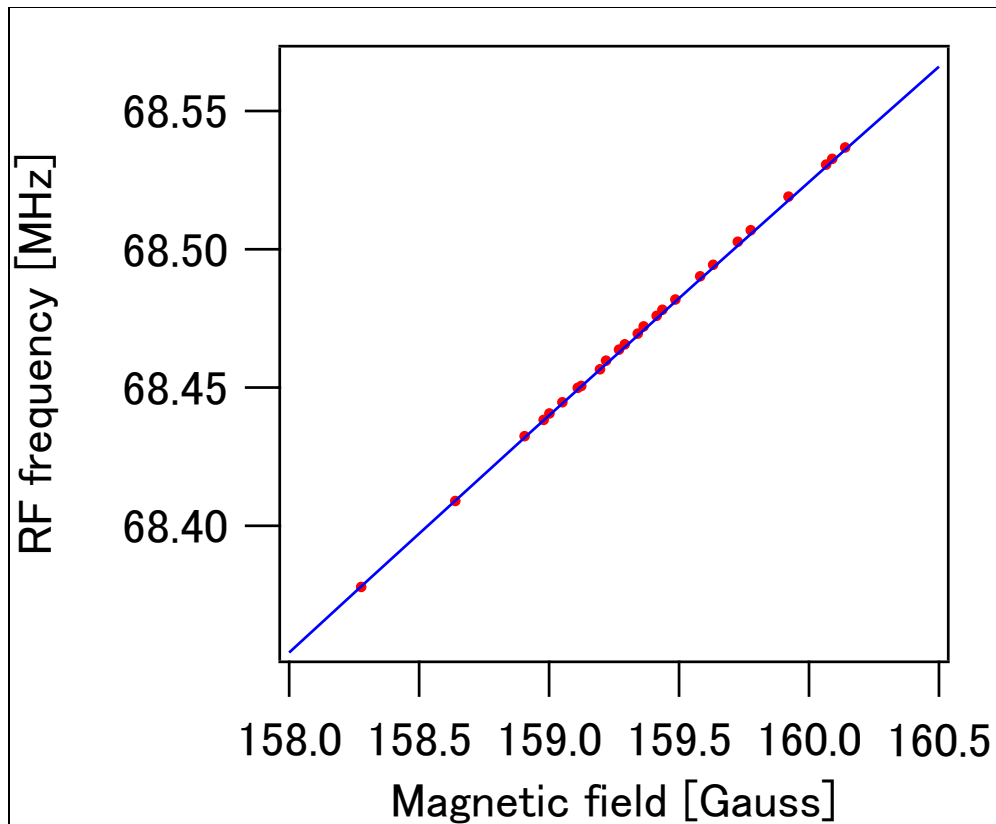
$|2\rangle$



$|1\rangle$

$|2\rangle$



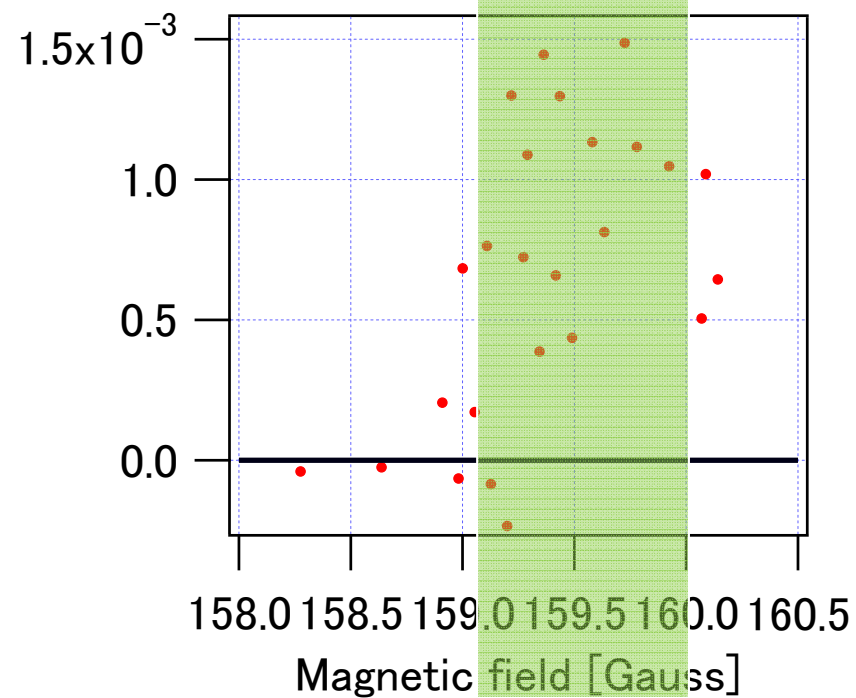


Preliminary!

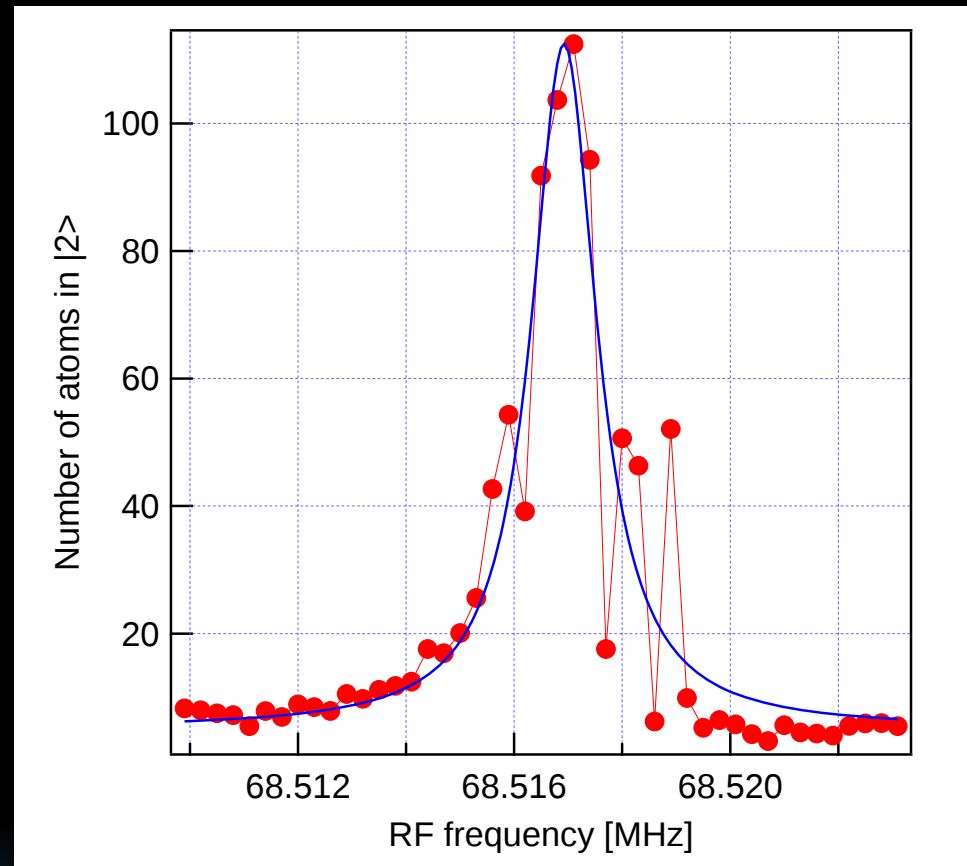
Region where atomic loss is prominent

Subtract obvious
energy shift from data

RF frequency [MHz]

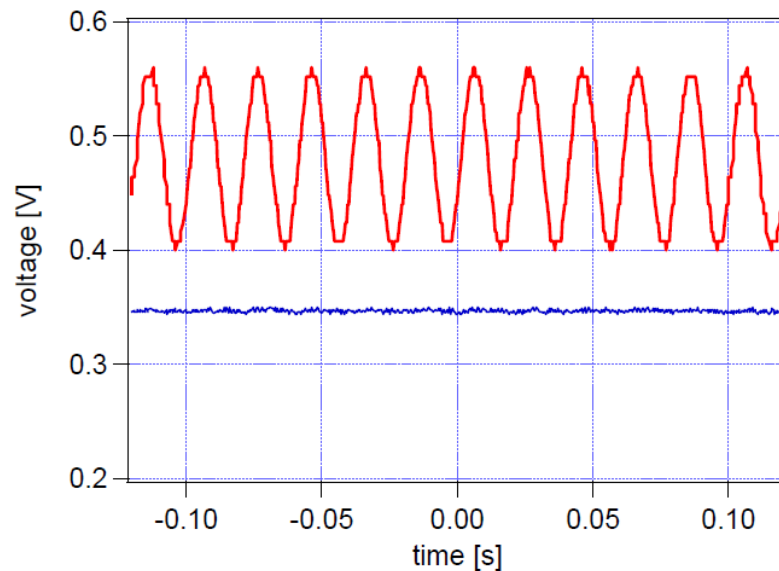
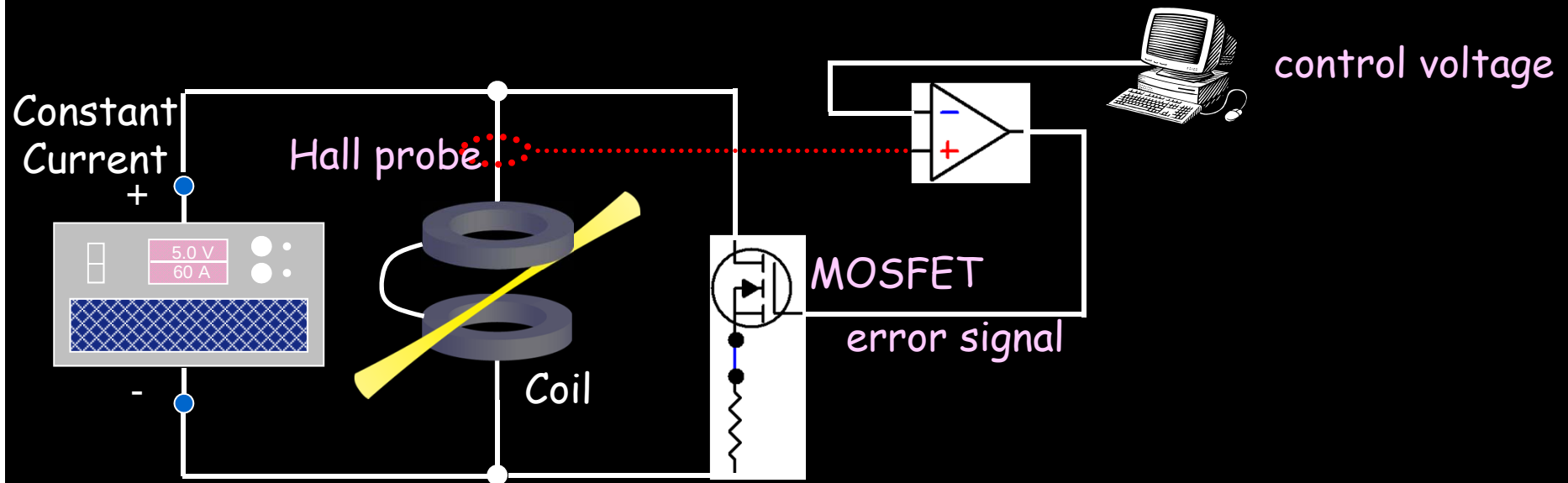


Typical RF spectrum



- 12 mG drift corresponds to 1 kHz (half div) shift of the resonance
- Slight asymmetry observed

current stabilization



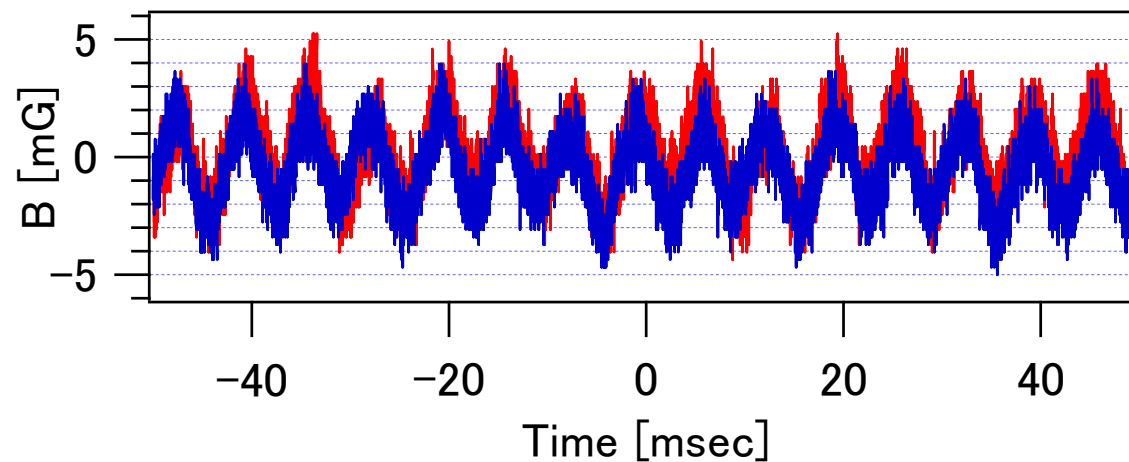
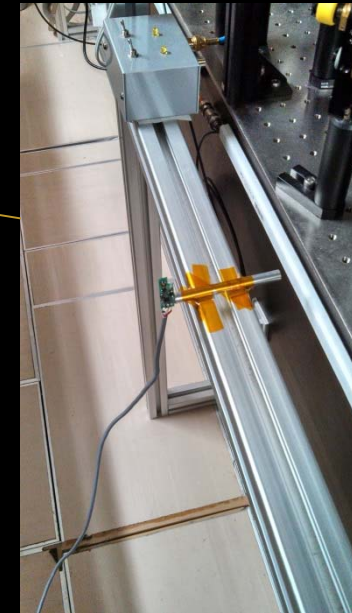
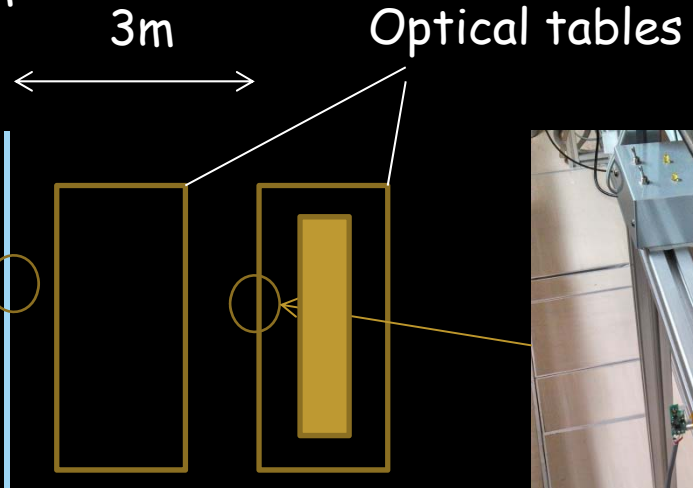
Without stabilization

B fluctuation ~ 300 mG

With stabilization

B fluctuation ~ 4 mG@160G

Magnetic field fluctuation



red: near the wall
blue: on the machine table

High-temperature expansion of a grand partition function

$$\Xi = \text{Tre}^{-(H-\mu N)/k_B T} = \Xi^{(0)} + 2\sqrt{2} \left(\frac{Vz^2}{\lambda^3} \right) b_2 + O(z^3)$$

$$z = e^{\mu/k_B T} \ll 1$$

$$b_2 = \sum_{\text{bound states}} e^{|E_b|/k_B T} + \sum_{l=0}^{\infty} (2l+1) \int_0^{\infty} \frac{dk}{\pi} \frac{d\delta_l(k)}{dk} e^{-\lambda^2 k^2 / 2\pi}$$

$$\lambda \equiv h / \sqrt{2\pi m k_B T}$$

"Statistical Mechanics"
Kerson Huang

Interaction energy

$$\varepsilon_{\text{int}} = \frac{3nk_B T}{2} (n\lambda^3) \left[-\frac{b_2}{\sqrt{2}} + \frac{\sqrt{2}}{3} T \frac{\partial b_2}{\partial T} \right]$$

T. -L. Ho and E. J. Muller, Phys. Rev. Lett.
92, 160404 (2004)

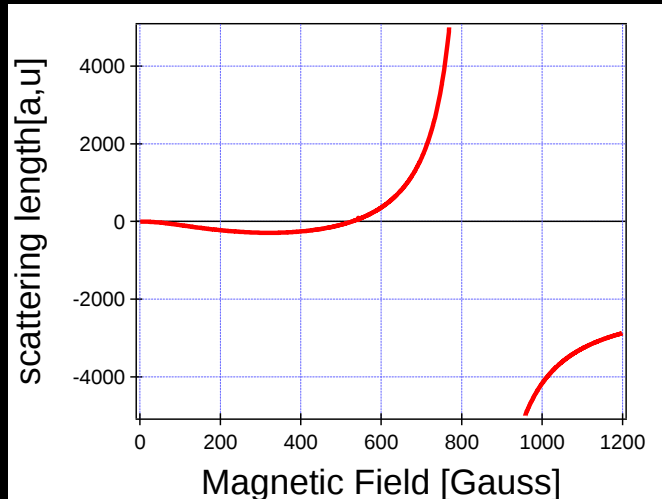
p-wave scattering phase shift

$$\cot \delta_p(k) = \frac{1}{k^3} \left(-\frac{1}{V} + k_e k^2 \right)$$

V and k_e have to be determined from a separate experiment.

Scattering parameters

s-wave, ${}^6\text{Li}$, $|1\rangle \rightarrow |2\rangle$



$$f_s(k) = \frac{1}{-1/a(B) + r_s k^2/2 - ik}$$

$$a(B) = a_{bg} \frac{1}{1 + \frac{DB}{B - B_{res}}} \dots$$

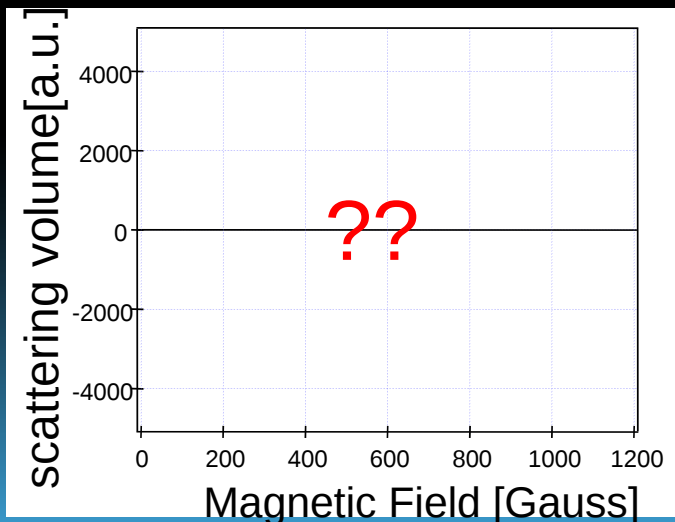
$$a_{bg} = 1582 a_0$$

$$DB = -262.3 [\text{G}]$$

$$B_{res} = 832.18 [\text{G}]$$

G. Zurn et al.
Phys. Rev. Lett.
110, 135301 (2013)

p-wave, ${}^6\text{Li}$, $|1\rangle \rightarrow |1\rangle$

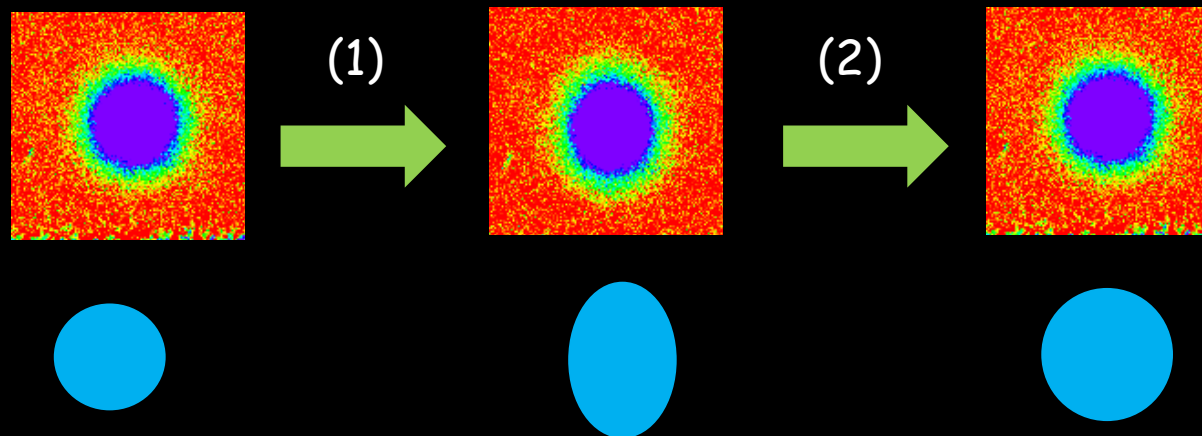


$$f_p(k) = \frac{k^2}{-1/V(B) + k_e k^2/2 - ik^3}$$

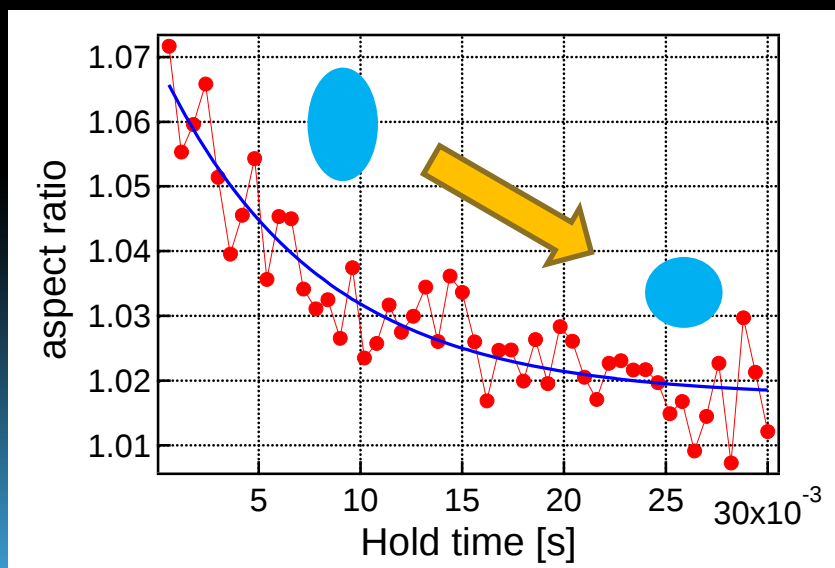
$$V(B) = a_{bg}^3 \frac{1}{1 + \frac{DB}{B - B_{res}}} \dots$$

 ... unknown parameters

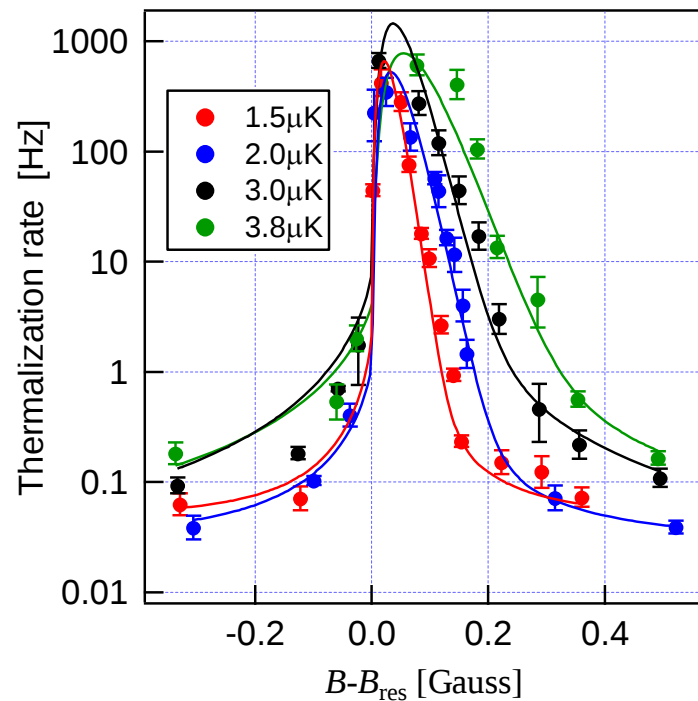
Images taken after ballistic expansion



- (1) modulate trap laser intensity at 3.4 kHz to add energy in one direction
- (2) cross-dimensional thermalization due to elastic collisions



Elastic collision cross-section can be derived from the timescale of the thermalization.



T. Nakasuji, J. Yoshida, and T. Mukaiyama,
Physical Review A **88**, 012710 (2013).

$$V_{\text{bg}} \pm DB = 2.8 \pm 10^6 a_0^3 [\text{Gauss}]$$

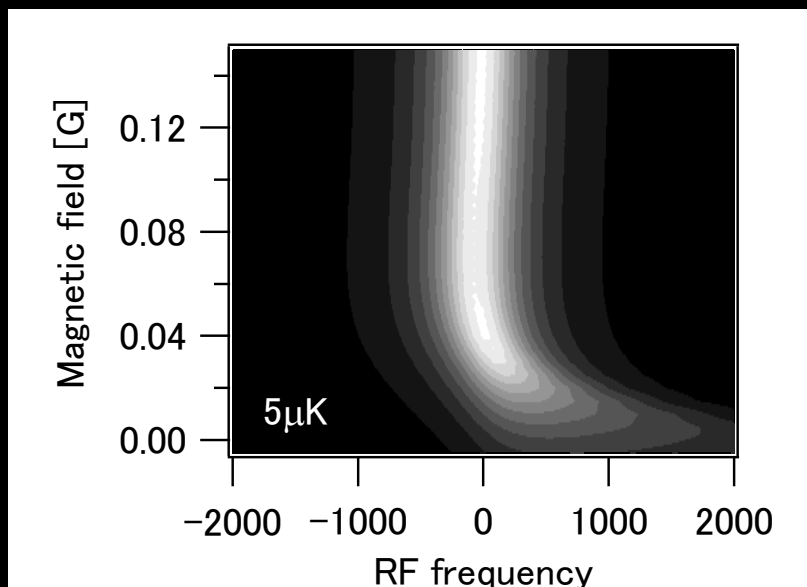
$$k_e = -0.058 a_0^{-1}$$

$$\langle n^2 \rangle \pm \sqrt{\frac{2}{mp}} (k_B T)^{-5/2} \cdot 4p(2l+1) |f_l(E)| E^2 \exp(-E/k_B T) dE$$

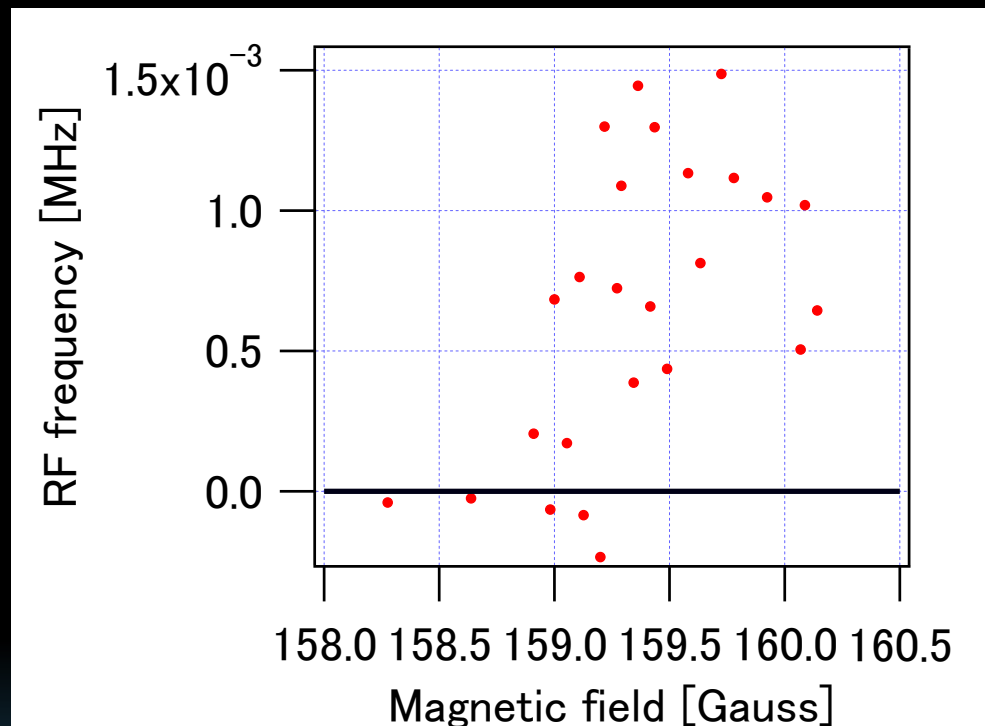
$$f_{l=1}(k, B) = \frac{k^2}{-\frac{1}{V(B)} + k_e k^2 - ik^3} \quad \frac{1}{V(B)} = V_{\text{bg}}^{-1} \begin{pmatrix} 1 \\ \vdots \end{pmatrix} + \frac{DB}{B - B_{\text{res}}} \begin{pmatrix} \vdots \\ -1 \end{pmatrix}$$

Calculation result

$$T=5\mu\text{K} \quad n\lambda^3 = 0.03$$



Order of energy shift ~ 1 kHz
Magnetic field range ~ 40 mG



Order of energy shift ~ 1 kHz
Magnetic field range ~ 1 G

Summary

RF spectroscopy to probe p-wave interaction in a spin-polarized Fermi gas

- Systematic study on temperature dependence
- Reduction of magnetic field fluctuation
- Further integration to improve S/N ratio

are needed to claim the first measurement of p-wave interaction energy.

研究室メンバー

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松田	優衣	M1
張	志琪	M1
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向山	敬	PI

