

強磁場原始中性子星からの 非等方ニュートリノ放出と関連現象

**Asymmetric Neutrino Emission from Magnetized Proto-Neutron
Stars and Its Related Astronomical Phenomena:
Pulsar Kick and Spin-Deceleration**

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§ 1. Introduction

High Density Matter Study $\Rightarrow$ Exotic Phases inside Neutron Stars

Strange Matter, Ferromagnetism, Meson Condensation, Quark matter

Observable Information $\cdots$ Neutrino Emissions

S.Reddy, et al., PRD58 #013009 (1998) **Influence from Hyperons Λ , Σ**

Magnetar 10^{15}G in surface 10^{17-19}G inside (?) $\rightarrow$ Large Asymmetry of ν ?

P. Arras and D. Lai, PRD60, #043001 (1999) , S. Ando, P.R.D68 #063002 (2003)

(Non-Relativistic)

Our Works $\Rightarrow$ Neutrino Reactions on High Density Matter with Strong Mag. Fields
in the Relativistic Mean-Field (RMF) Approach

TM et al., PRD83, 081303(R) (2011)

Application to Phenomena related with ν — Asymmetric Emission

Pulsar Kick in Poloidal Magnetic Field

Spin Deceleration in Toroidal Magnetic Field

§ 2. Formulation

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Magnetic Field : $\vec{B} = B\hat{z}.$

Lagrangian : $\mathcal{L} = \mathcal{L}_{RMF} + \mathcal{L}_{lep.} + \mathcal{L}_{mag} + \mathcal{L}_{int}$

Baryon

Lepton

B & L – Mag.

1. Proto-Neutron-Star (PNS) Matter without Mag. Field
2. Baryon Wave Function under Mag. Field in Perturbative Way
3. Cross-Sections for ν reactions

Weak Interaction

$$\mathcal{L}_{int} = G_F \{ \bar{\psi}_l \gamma_\mu (1 - \gamma_5) \psi_l \} \{ \bar{\psi}_{B'} \gamma^\mu (c_V - c_A \gamma_5) \psi_B \}$$

$\nu_e + B \rightarrow \nu_e + B$: scattering

$\nu_e + B \rightarrow e + B'$: absorption

§ 2-1 EOS of Proto Neutron-Star-Matter in RMF $N, \Lambda, \sigma, \omega, \rho$

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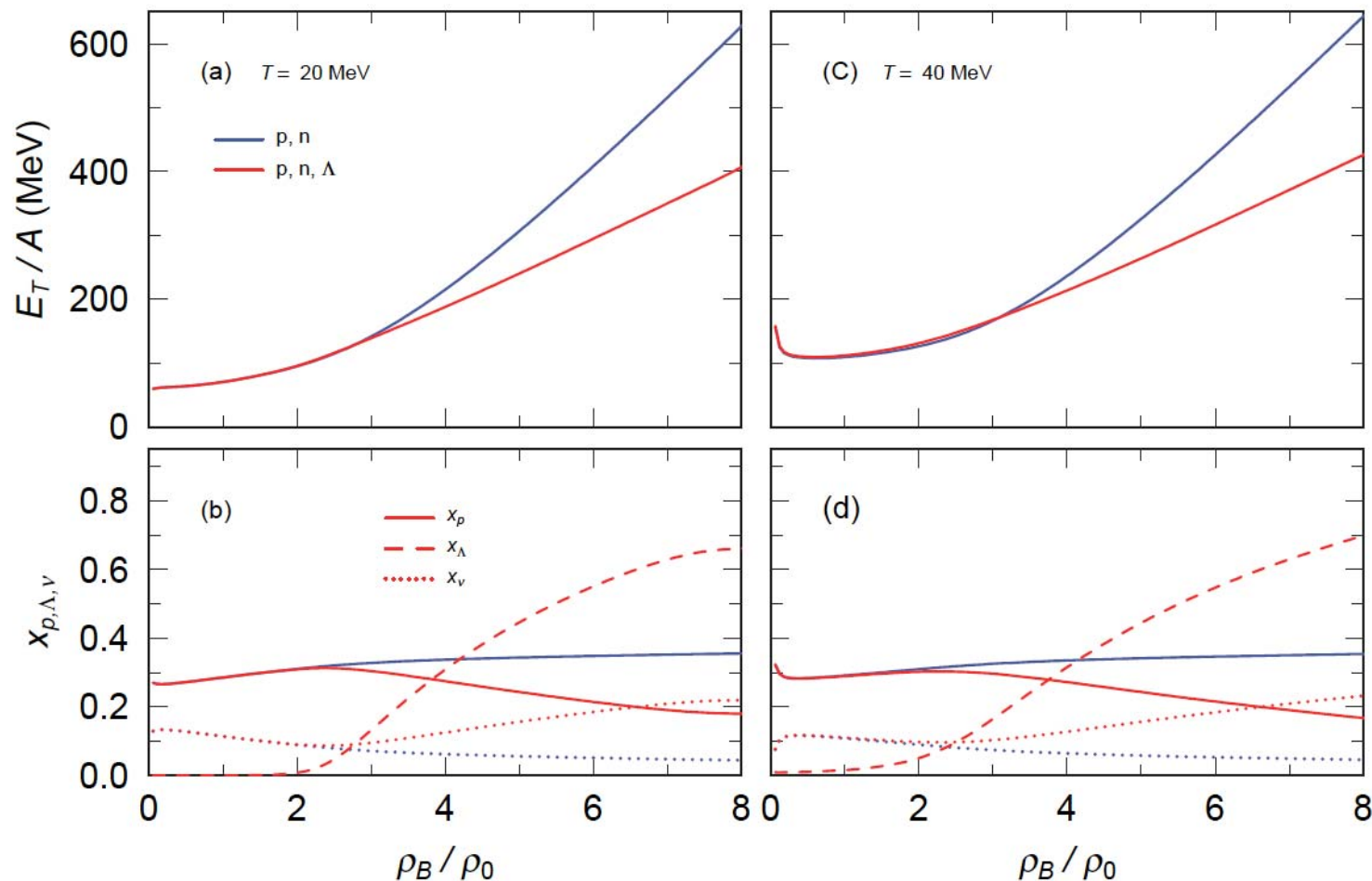
PM1-L1

$$BE = 16 \text{ MeV}, M_N^* / M_N = 0.7, K = 200 \text{ MeV at } \rho_0 = 0.17 \text{ fm}^{-3}$$

$$g_{\sigma,\omega}^\Lambda = \frac{2}{3} g_{\sigma,\omega} \quad \text{SU(3)}$$

T.M, et al.
PTP. 102, p809
(1999)

Charge Neutral ($\rho_p = \rho_e$) & Lepton Fraction : $Y_L = 0.4$



§ 2-2 Dirac Equation under Magnetic Fields

$\mu_N B \ll \varepsilon_N$ (Chem. Pot) $\rightarrow B$ can be treated perturbatively
 $B \sim 10^{17}$ G Landau Level can be ignored

Lagrangian

$$\mathcal{L}_{mag} \approx - \sum_b \mu_n B \bar{\psi}_b \sigma_z \psi_b$$

Dirac Eq.

$$\hat{K}(p)u(p) = [\gamma_\mu p^\mu - M^* - \mu B \sigma_z] u(p) = 0$$

Single Part. Eng.

$$e(\mathbf{p}, s) = \left[\left(\sqrt{p_x^2 + M^{*2}} + s\mu B \right)^2 + p_z^2 \right]^{\frac{1}{2}} \approx E_p^* + s\mu B \frac{\sqrt{p_T^2 + M^{*2}}}{E_p^*}$$

Dirac Spinor

$$u(\mathbf{p}, s) \bar{u}(\mathbf{p}, s) \approx \frac{(\not{p} + M)(1 + \gamma_5 \not{d}(p)s)}{4E_p^*}$$

Spin Vector

$$a_z = \frac{E_p}{\sqrt{p_T^2 + M^2}} \quad \mathbf{a}_T = 0 \quad a_0 = \frac{p_z}{\sqrt{p_T^2 + M^2}}$$

Relativistic Effects of Magnetic Contributions

- 1) Mom.-Dep. of Spin Vector
- 2) Deformation of Fermi Distr.

The Cross-Section of Lepton-Baryon Scattering

$$\frac{d^2\sigma}{dk' d\Omega'_k} = \frac{G_F^2}{8\pi^2} k'^2 \sum_{s_i, s_f} \int \frac{d^3p}{(2\pi)^3} \tilde{W}_{BL} (2\pi) \delta(|\mathbf{k}| - |\mathbf{k}'| + e_i(\mathbf{p}) - e_f(\mathbf{k} + \mathbf{p} - \mathbf{k}')) \\ \times [1 - f'_l(\mathbf{k}')] n_B(e_i) [1 - n_{B'}(e_f)]$$

$$\tilde{W}_{BL} = \text{Tr} \left\{ \frac{(\not{\mathbf{k}}' + m_f)(1 + \gamma_5 \not{\phi}_{l'})}{4|\mathbf{k}'|} \gamma^\mu (1 - \gamma_5) \frac{\not{\mathbf{k}}'}{2|\mathbf{k}|} \gamma^\nu (1 - \gamma_5) \right\} \\ \times \text{Tr} \left\{ \frac{(\not{\mathbf{p}}' + M_f^*)(1 + \gamma_5 \not{\phi}_f(\mathbf{p}'))}{4E_f^*(\mathbf{p}')} \gamma_\mu (c_V - c_A \gamma_5) \frac{(\not{\mathbf{p}} + M_i^*)(1 + \gamma_5 \not{\phi}_i(\mathbf{p}))}{4E_i^*(\mathbf{p})} \gamma_\nu (c_V - c_A \gamma_5) \right\}$$

$$m_f = 0 \quad \text{when } l_f = \nu \quad m_f = m_e \quad \text{when } l_f = e$$

Fermi
Distribution

$$n(e(\mathbf{p}), s) \approx n(\varepsilon(\mathbf{p}, s)) + n'(\varepsilon(\mathbf{p}, s)) \frac{\sqrt{p_T^2 + M^{*2}}}{E_p^*} \mu B s.$$

Deformed Distribution

Perturbative
Treatment

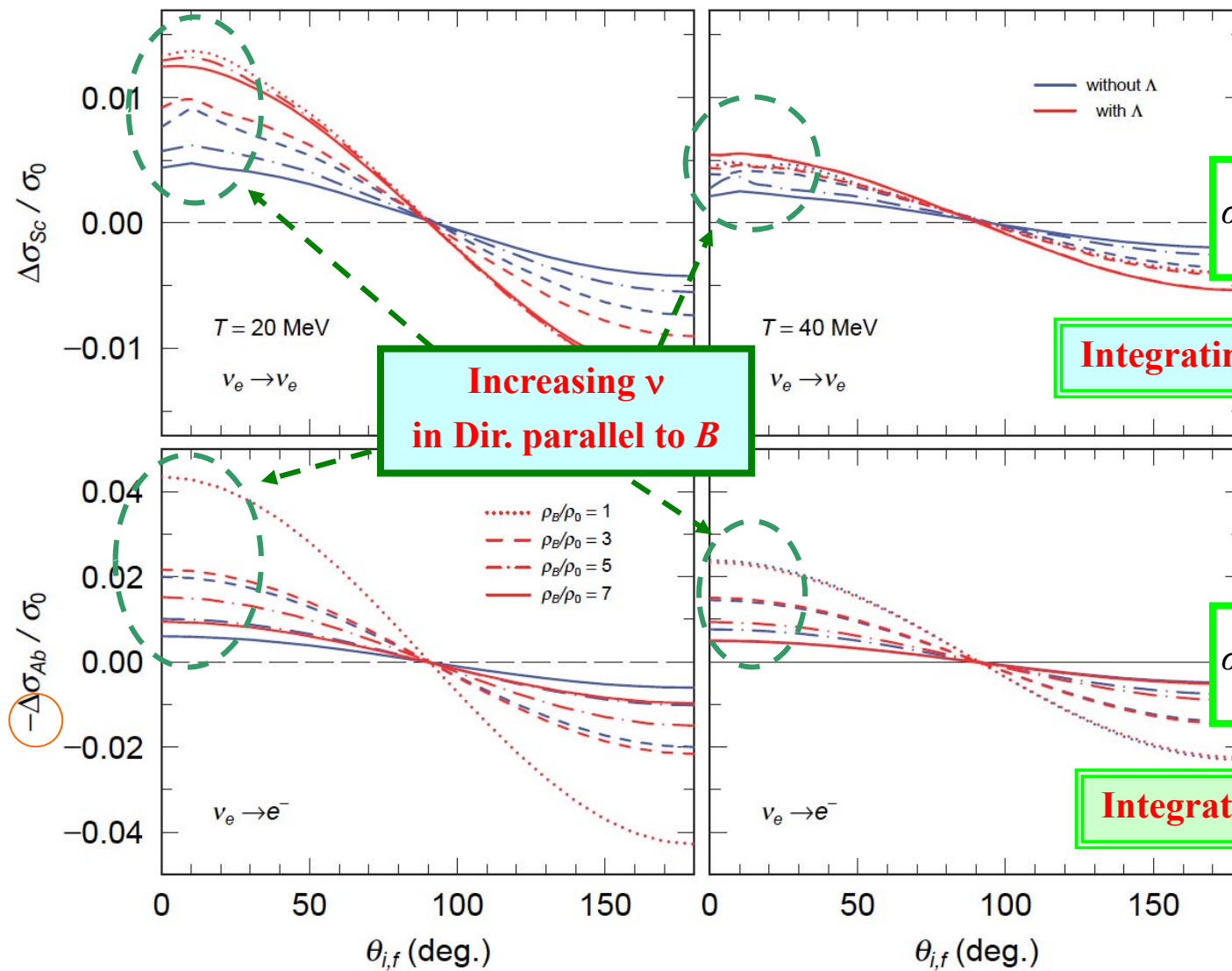
$$\sigma = \sigma_0 + \Delta\sigma \quad \Delta\sigma \propto B$$

Non-Magnetic Part

Magnetic Part

§ 2-3 Magnetic parts of Cross-Sections

$$\sigma = \sigma_0 + \Delta\sigma \quad \Delta\sigma \propto B$$



$$k_i = \varepsilon_\nu \text{ (neutrino chem. pot.)}, \quad B = 2 \times 10^{17} \text{ G} \quad \text{and} \quad \theta_i = 0^\circ$$

§ 2-4 Neutrino Transportation

Neutrino Phase Space Distribution Function

$$f(\mathbf{p}, \mathbf{r}) \approx f_0(\mathbf{p}, \mathbf{r}) + \Delta f(\mathbf{p}, \mathbf{r}), \quad f_0(\mathbf{p}, \mathbf{r}) = 1 / \{1 + \exp [(|\mathbf{p}| - \varepsilon_\nu) / T]\}$$

Equib. Part

Non-Equib. Part

Neutrino Propagation $\Rightarrow$ Boltzmann Eq.

$$c \frac{\partial}{\partial \mathbf{r}} f_0(\mathbf{p}, \mathbf{r}) \approx c \frac{\partial}{\partial \mathbf{r}} f_0(\mathbf{p}, \mathbf{r}) + c \frac{\partial}{\partial \mathbf{r}} \Delta f(\mathbf{p}, \mathbf{r}) = I_{coll} \approx -c b_\nu \Delta f(\mathbf{p}, \mathbf{r}), \quad b_\nu = \frac{\sigma_{ab}}{V}$$

Neutrinos Propagate on **Strait Line**

only absorption

Solution $\Rightarrow$

$$\Delta f(\mathbf{p}, \mathbf{r}_T, z) = \int_0^z dx \left[-\frac{\partial}{\partial x} f_0(p, r_T, x) \right] \exp \left[-\frac{1}{c} \int_x^z dy b_\nu(y) \right],$$

$$z = \mathbf{r} \cdot \hat{\mathbf{p}}, \quad \frac{\partial}{\partial z} f_0(p, r_T, z) = \frac{d\varepsilon_\nu}{dz} \frac{\partial}{\partial \varepsilon_\nu} f_0(p, r_T, z)$$

§ 3 Estimating Pulsar Kick Velocities of Proto-Neutron Star

A.G.Lyne, D.R.Lomier, Nature 369, 127 (94)

Asymmetry of Supernova Explosion

kick and translate Pulsar with

Kick Velocity: Average ... 400km/s,

Highest ... 1500km/s

Explosion Energy $\sim 10^{53}$ erg
(almost Neutrino Emissions)

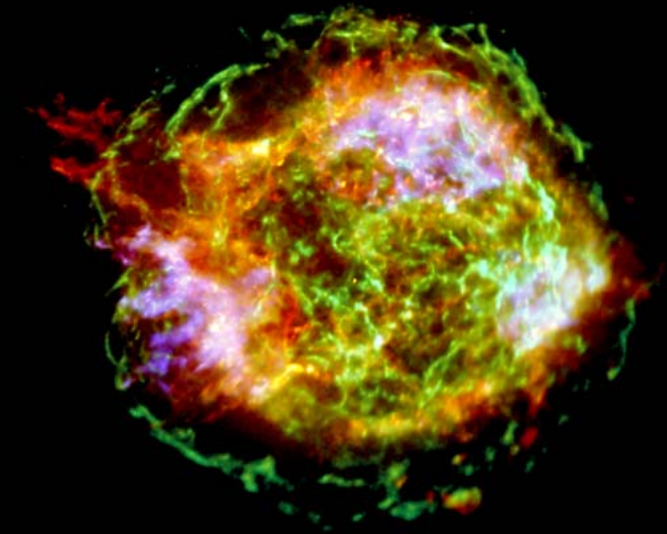
1% Asymmetry is sufficient to explain the Pulsar Kick

D.Lai & Y.Z.Qian, Astrophys.J. 495 (1998) L103

Estimating Kick Velocity of PNS with $T = 20$ MeV and $B = 2 \times 10^{17}$ G

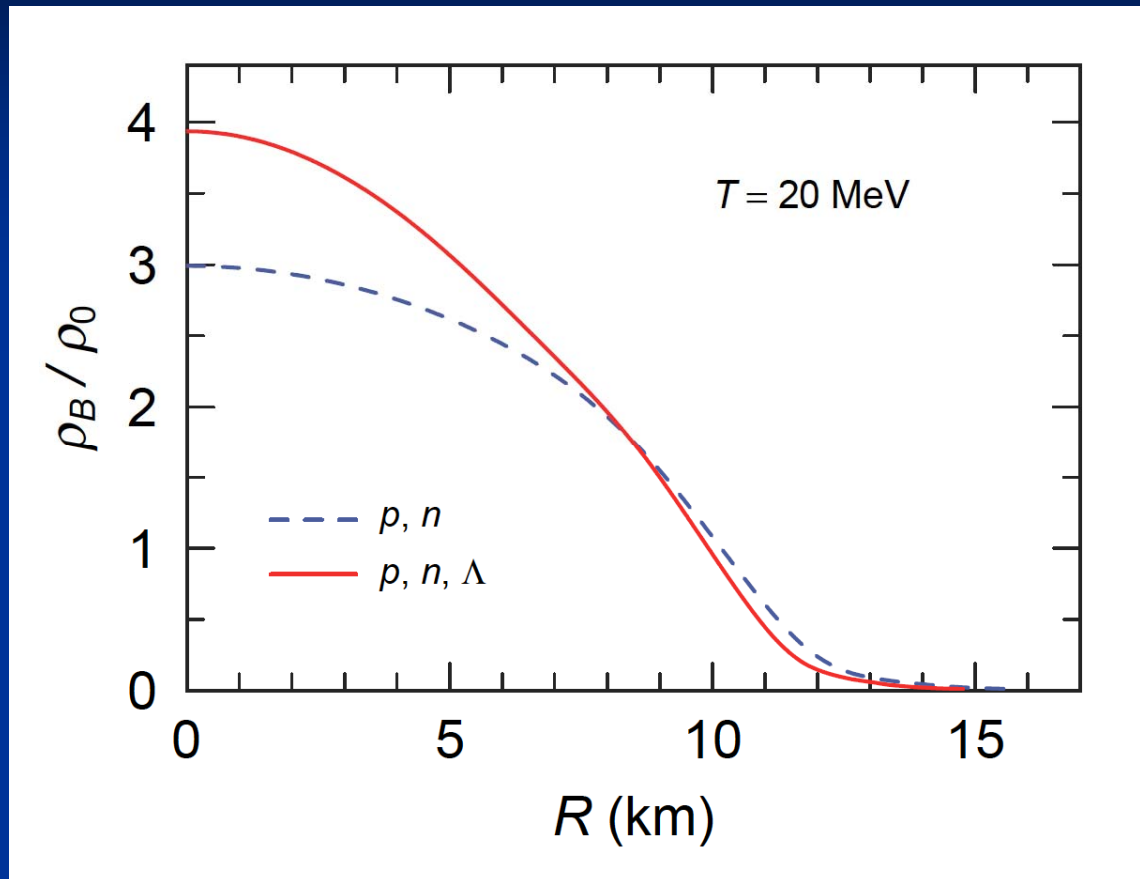
Poloidal Magnetic Field 2 - 3% Asymmetry in Absorption

CasA



http://chandra.harvard.edu/photo/2004/casa/casa_xray.jpg

Baryon density in Proto-Neutron Star



$$M = 1.68 M_{\text{solar}}$$

$$Y_L = 0.4$$

$$P_\alpha = \int_{S_N} dr \int dp p_\alpha \Delta f(r, p_L, n) / \int_{S_N} dr \int dp \Delta f(r, p_L, n)$$

Calculating Neutrino Propagation above $\rho_B = \rho_0$

Angular Dep. of Emitted Neutrinos in Uniform Poloidal Mag. Field

$$M_{NS} = 1.68 M_{solar} \text{ [g]},$$

$$E_T = 3 \times 10^{53} \text{ [erg]}$$

$$\langle p(\mathbf{n}) \rangle \approx P_0 + P_1 \cos \theta$$

$$B = 2 \times 10^{17} \text{ [G]}$$

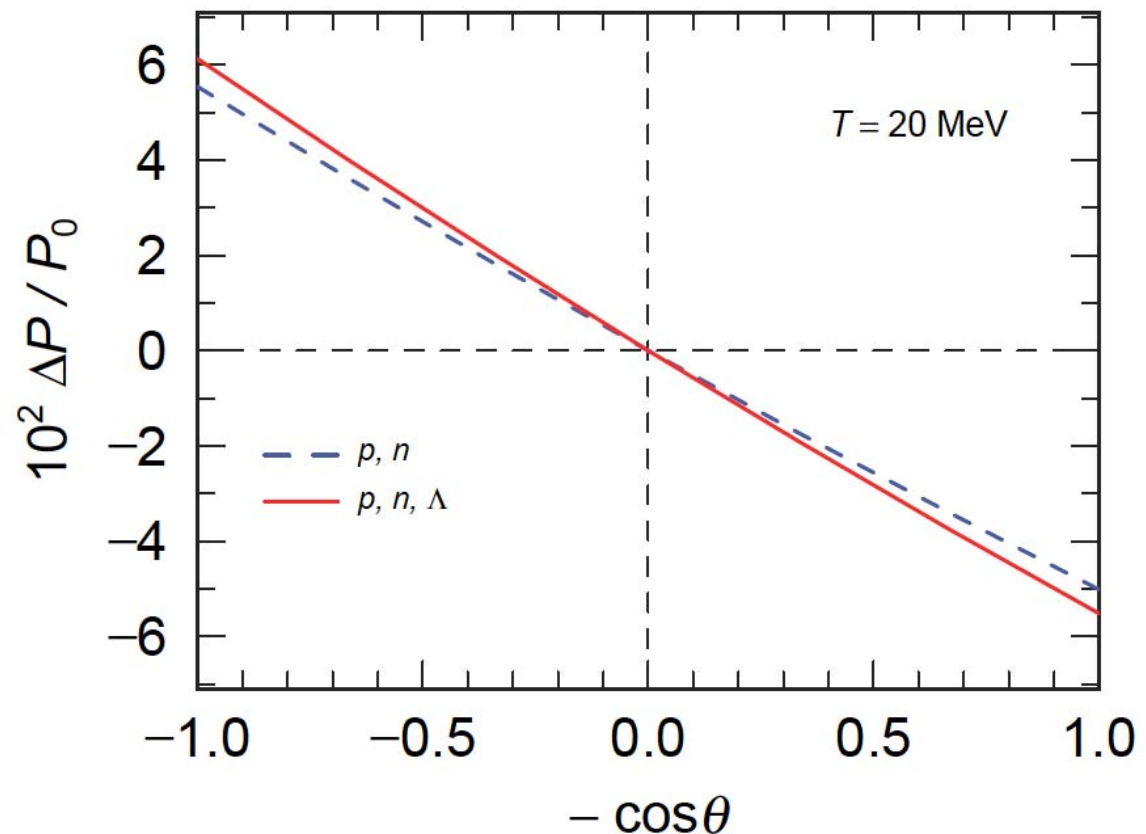
$$v_{kick} = \frac{\langle p_z \rangle}{M} \approx 600 \text{ [km/s]}$$

$$T = 20 \text{ MeV}$$

Observable Average
... 400km/s,

$$\frac{\langle p_z \rangle}{E_T} = \frac{P_1}{3P_0} = 1.76 \times 10^{-2} \quad p, n$$

$$= 1.94 \times 10^{-2} \quad p, n, \Lambda$$



§ 4 Angular Deceleration in Toroidal Magnetic Field

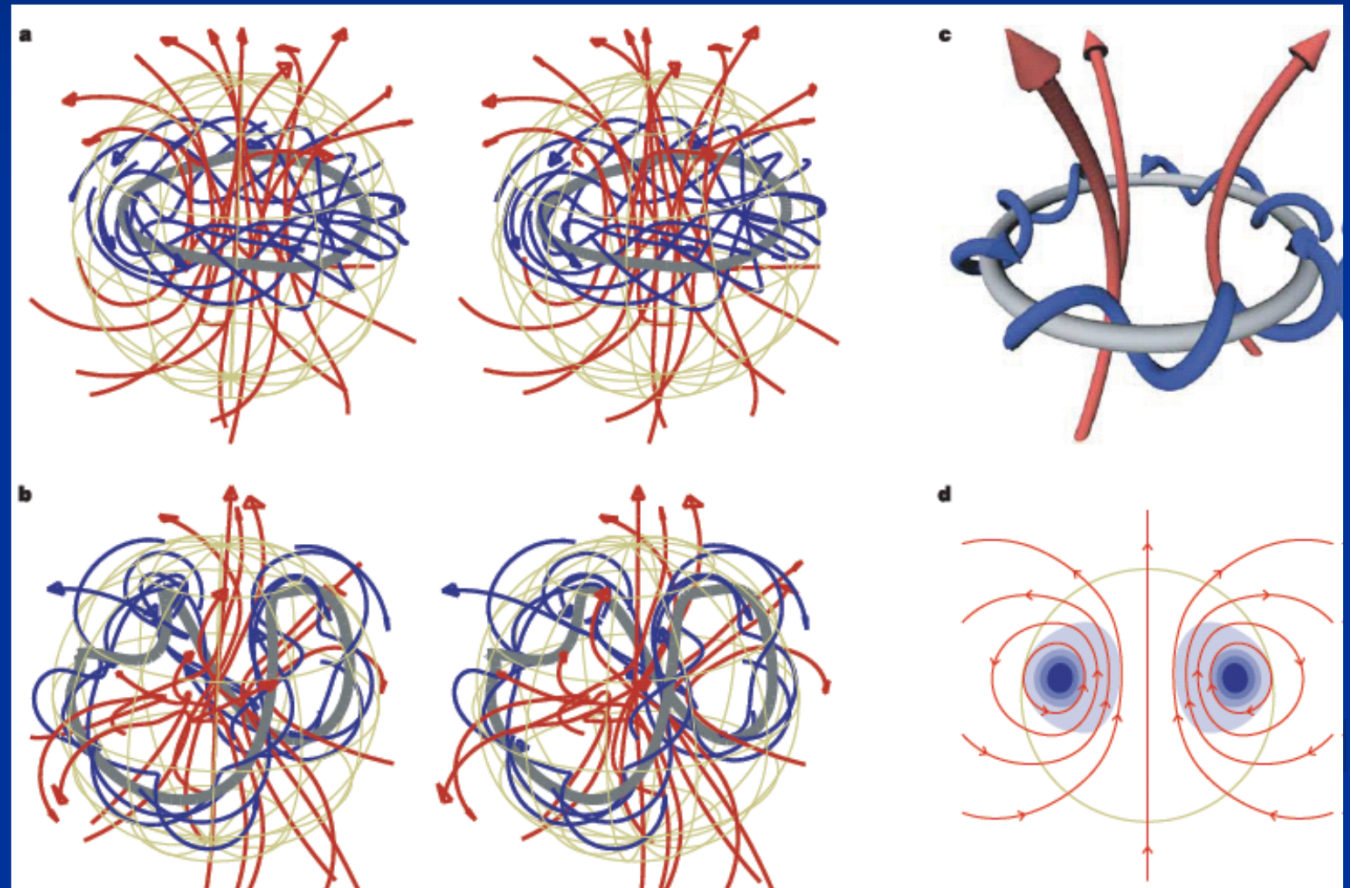
Stability of Magnetic Field in Compact Objects

(Braithwaite & Spruit 2004)

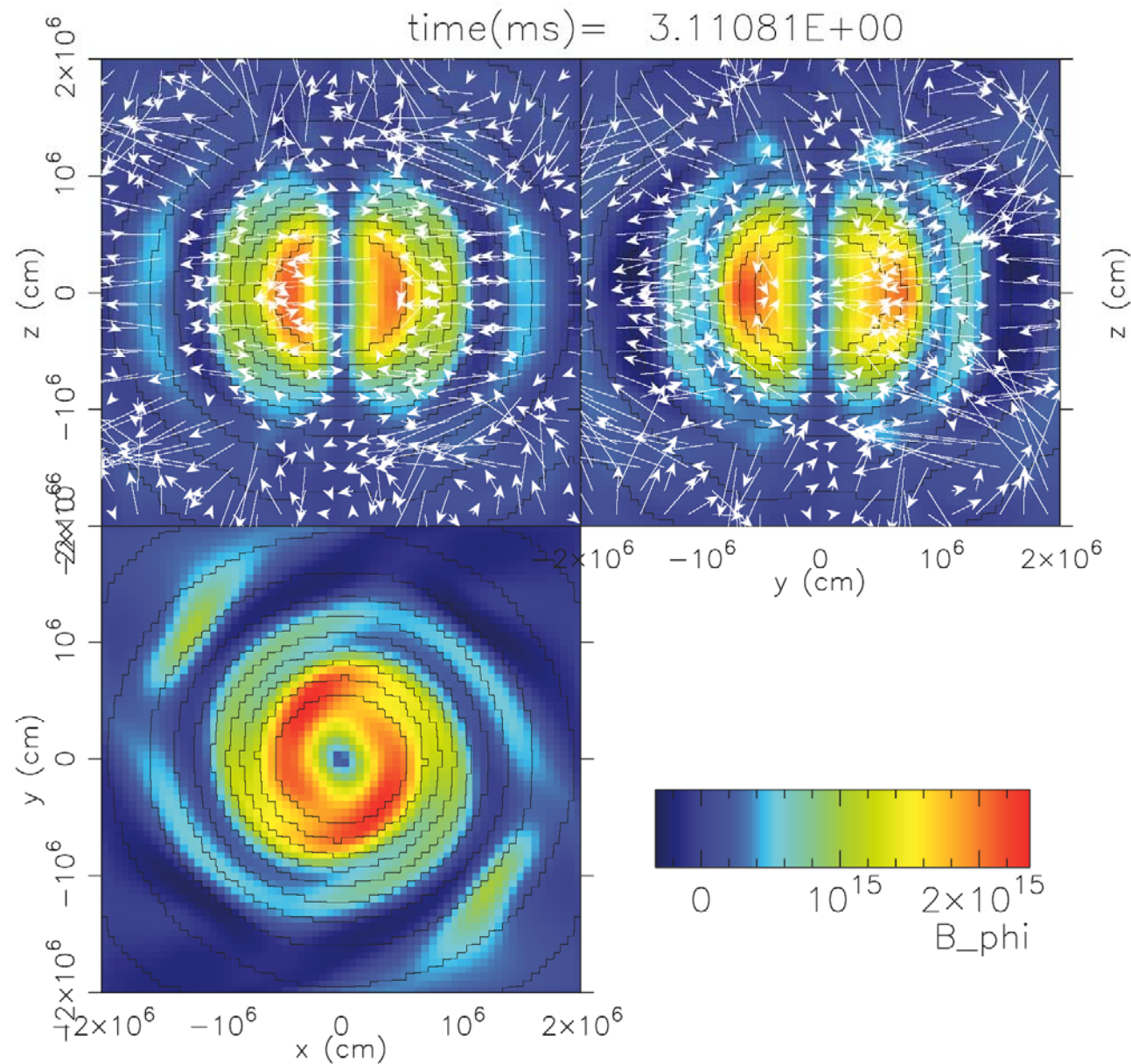
Toroidal Magnetic
Field is stable !!

Mag. Field
Parallel
to Baryonic Flow

Assym. of ν -Emit.
must decelerate
PNS Spin



No poloidal Magnetic Field at the beginning



**Single
Toroidal**

by T. Kuroda

Toroidal Magnetic Field

$$B(r_T, z) = B_0 G_T(r_T) G_L(z) \hat{e}_\phi$$

$$G_T(r_T) = \frac{16 \exp[-(r_T - R_0) / \Delta r]}{\{1 + \exp[-(r_T - R_0) / \Delta r]\}^2}$$

$$G_L(z) = \frac{\exp[-z / \Delta r]}{\{1 + \exp[-z / \Delta r]\}^2}$$

$$\hat{e}_\phi = (-\sin \phi, \cos \phi, 0)$$

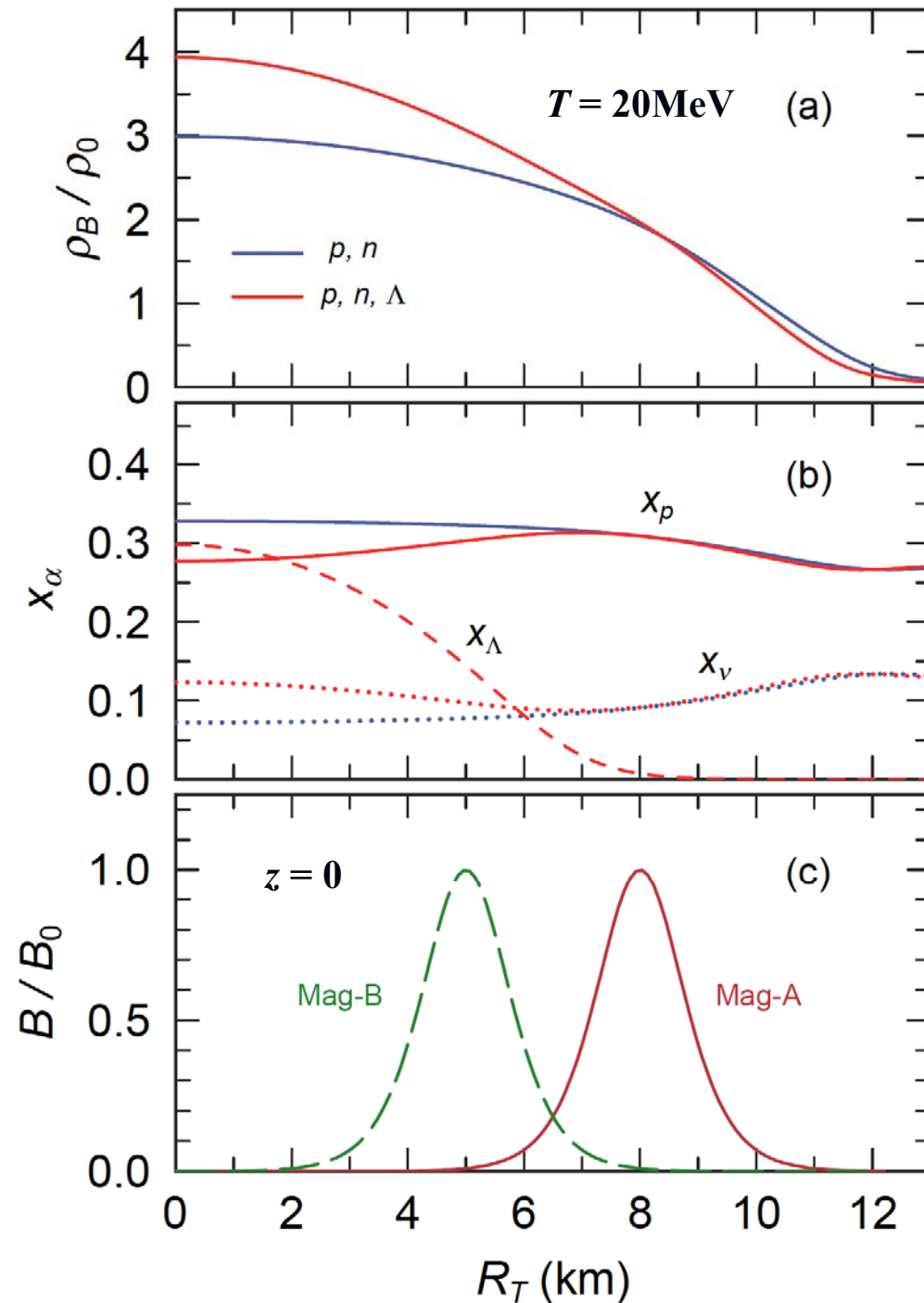
$$\Delta r = 0.5 \text{ (km)}$$

$$R_0 = 8 \text{ (km)}$$

Mag. - A

$$R_0 = 5 \text{ (km)}$$

Mag. - B



$$\frac{dL_z}{dt} = c \int_{S_N} d\mathbf{r} \int d\mathbf{n} \int dp_L \Delta f(\mathbf{r}, p_L, \mathbf{n}) (\mathbf{r} \times \mathbf{p})_z$$

$$\dot{\omega} = \frac{d\omega}{dt} = \frac{1}{I_{NS}} \frac{dL_z}{dt} = \frac{1}{I_{NS}} \left(\frac{dE_T}{dt} \right)_\nu \frac{cdL_z / dt}{dE_T / dt}$$

$$\frac{-\dot{P}}{P} = \frac{\dot{\omega}}{\omega} = \frac{P}{2\pi I_{NS}} \left(\frac{dE_T}{dt} \right)_\nu \left(\frac{cdL_z / dt}{dE_T / dt} \right) \propto P$$

Neutrino Luminosity

$$(dE_T/dt)_\nu \sim 3 \times 10^{52} \text{ erg/s}$$

Magnetic Dipole Radiation (MDR)

$$P\dot{P} = B_{pol}^2 \left(\frac{125\pi^2 I_{NS}^2}{3M_{NS}^2 c^3} \right)$$

$$B_{tro} \approx 100 B_{pol}$$

Results in Asymmetric Neutrino Emmission

Neutrino Luminosity $(dE_\nu/dt)_\nu \sim 3 \times 10^{52} \text{ erg/s}$

$M_{NS} = 1.68 M_{solar}$

Poloidal ($B_{pol} = 10^{14} \text{G}$) + **Toroidal** ($B_0 = 10^{16} \text{G}$)

T. Takiwaki, K.Katake and K. Sato Astro. J 691, 1360 (2009)

Period $P = 10 \text{ms}$

Mag Distr.	Bary.	$-\frac{cdL_z / dt}{dE_T / dt}$ (cm)	$\dot{P}/P$		MDR
			$(\nu \text{ emis.})$		
			$\rho_s = \rho_0$	$\rho_s = \rho_0 / 10$	
p,n	Mag-A	3.34	3.45×10^{-6}	7.25×10^{-7}	9.86×10^{-8}
	Mag-B		4.97×10^{-7}	3.16×10^{-7}	
p,n, Λ	Mag-A	5.45	6.39×10^{-6}	1.02×10^{-6}	7.76×10^{-8}
	Mag-B		4.57×10^{-7}	2.01×10^{-7}	

In Early Stage ($\sim 10 \text{ ms}$) ν **Asymmetric** Emission must affect **PNS Spin**

More Significantly than Magnetic Dipole γ -Radiation

§ 5 *Summary*

■ Asymmetry of Neutrino Absorption

4.3 % at $\rho_B = \rho_0$, 2.2 % at $\rho_B = 3\rho_0$ when $T = 20$ MeV and $B = 10^{17}$ G

■ Pulsar Kick in Poloidal Mag. Field $B = 2 \times 10^{17}$ G $\Rightarrow$ Perturbative Cal.

$v_{\text{kick}} = 580$ km/s (p,n), 610 km/s (p,n, Λ) at $T = 20$ MeV

400 km/s (Average of Observed Values)

■ Estimating Spin-Down Rate of PNS with Toroidal Magnetic Field Config.

Mag. Field Poloidal 10^{14} G, Toroidal Max: 10^{16} G

Spin-Down Ratio $\dot{P}/P \approx 10^{-6} \sim 10^{-7}$ (1/s) for Asym. ν -Emit

$\approx 10^{-8}$ (1/s) for MDR

Extension

Other Effects: ν -Scattering & ν -Production

Iso-Temp. $\Rightarrow$ Iso-Entropy

Neutrino Propagation in Low Density $e^- + p \rightarrow \nu_e + n$

Examining Density-Dependence of Symmetry Energy

Exact Solution of Dirac Eq. in Non-Perturbative Cal.

$\rightarrow$ **Landau Level** at least for Electron

Gamma-Ray Emission in Strong Magnetic Field

$$p \rightarrow p + \pi^0$$