

2013年7月7日

和州閣

# $S=-1$ から $S=-2$ の分光へ

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# 内容

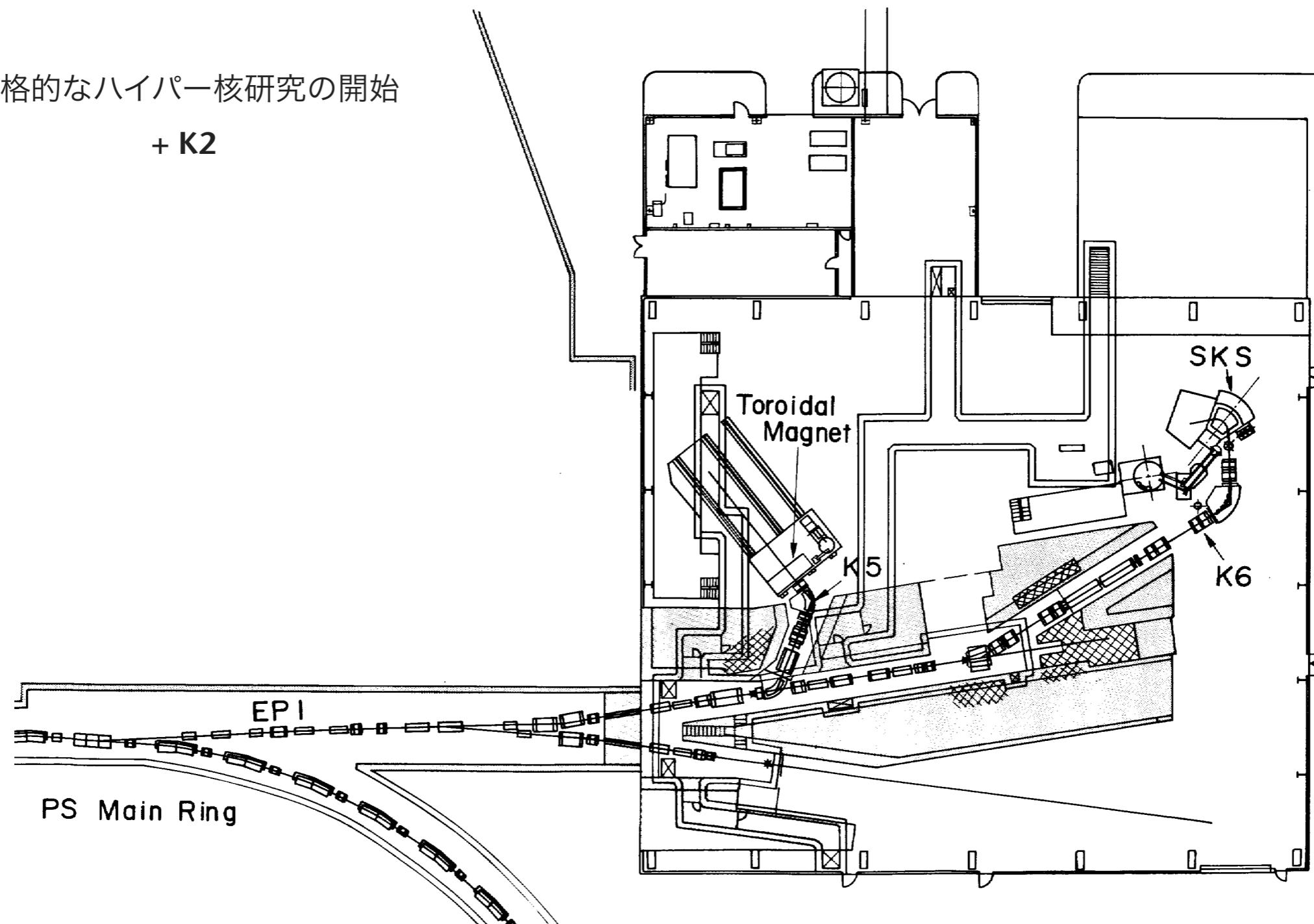
- SKSスペクトロメーター
- $(\pi^+, K^+)$ スペクトロスコピー
- E140a, E336, E369
- S-2Sの建設と $(K^-, K^+)$ スペクトロスコピー

# 北カウンターホール

KEK 12 GeV PS  
New Experimental Hall

本格的なハイパー核研究の開始

+ K2



# SKS建設の歴史

- 1987: 建設開始
- 1990: 超伝導磁石完成 → 北カウンターホールへ据え付け
- 1991: 冷凍システム完成
- 1991.9: 3T励磁に成功！ → 磁場測定
- 1992.2: SKS+K6ビームコミッショニング開始
- 1992.10-11: E269実験( $\pi$ -C弹性散乱) データ取得
- 1993.1-7: E140a実験データ取得開始 (最初の( $\pi$ , K)実験)

# 2人の巨人逝く：1990,91

T. Motoba, NPA547 (1992) 379c-392c.

In memory of professor Hiroharu Bandō , I quote a part of the note added by Žofka and myself in our recent review article [3].

"Hiro started his scientific work at Kyoto University with an investigation on nuclear collective motion in 1961 and did many noticeable studies with more than 110 papers which are related to three groups: (i) the many-body theory of the effective interactions based on the realistic nuclear force, (ii) the cluster model of nuclei and clustering mechanism, and (iii) the theory of hyper-nuclear physics covering the various stages (production, structure and decay) of hypernuclei."



Dr. Hiroharu Bandō (1938–1990)

In memory of professor Jan Žofka, let me simply quote several sentences from a recent article by G. Ripka[5].

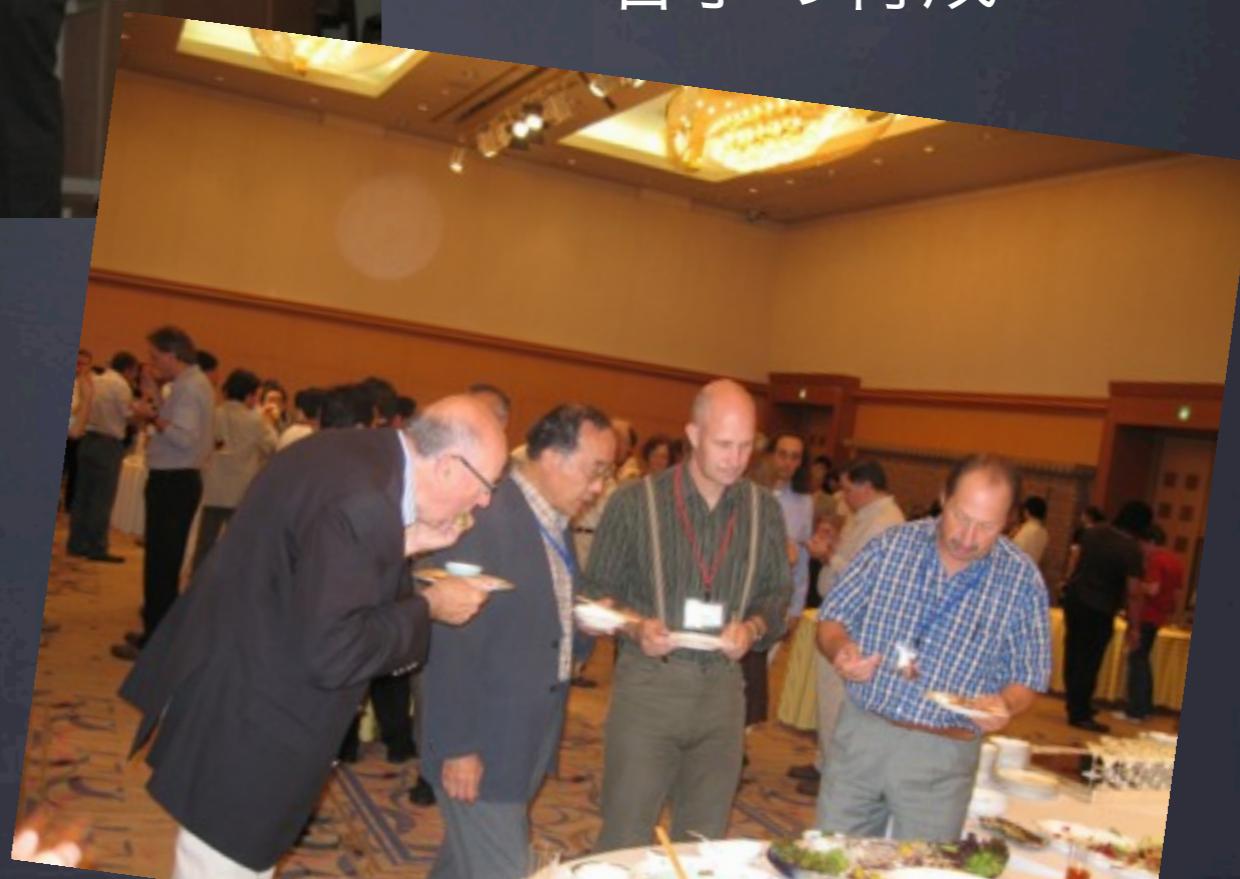
"Jan Žofka, the internationally known expert on hypernuclei, died unexpectedly last May (5th, 1991) in Prague. Žofka graduated in 1966 from the Czech Technical Faculty of Nuclear Physics and entered the Institute of Nuclear Physics in Řež near Prague, where he worked until his death. He started as a student of Professor Trlifaj and after a six month visit to the Service de Physique at CEN Saclay, France, he completed, in 1971 in Prague, his PhD on the Hartree-Fock theory applied to nuclear shape." [5]



Dr. Jan Žofka (1943 – 1991)

# 元場さんの支援

- 現実的な理論計算
- ハイパー核コミュニティの形成
- 若手の育成



# SKS建設の歴史（続き）

- 1987: 建設開始
- 1990: 超伝導磁石完成→北カウンターホールへ据え付け
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# SKSでの実験のまとめ

- 実験期間：1992～2005（14年間）
- K2Kライン建設のため1年間シャットダウン
- 実施ビームタイム：～1700シフト（570日）
- 実験数：15課題
- Ph.D取得数：20人

# ( $\pi^+$ , $K^+$ ) スペクトロスコピー

（主に、 $\pi^+$  の観測）

（主に、 $K^+$  の観測）

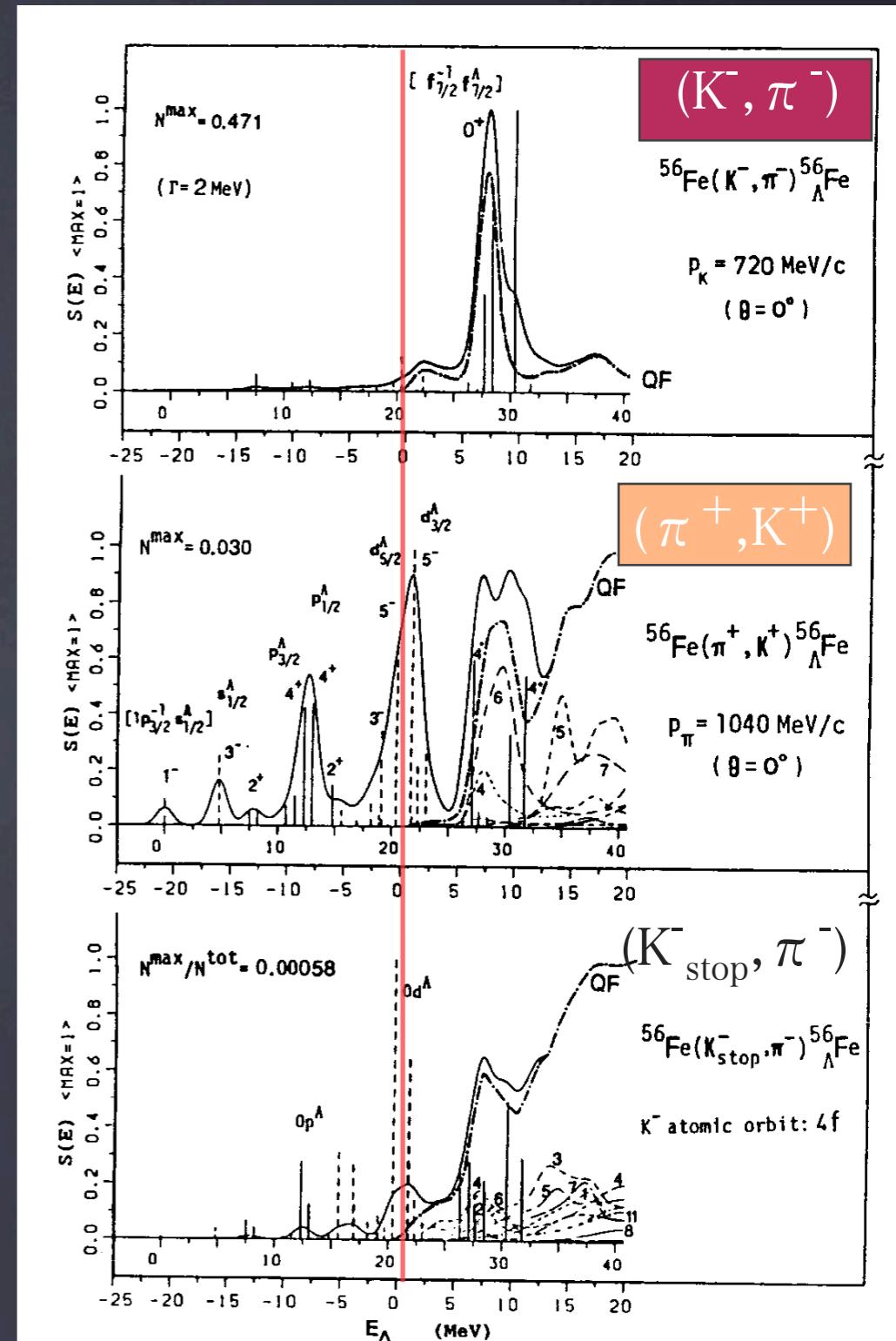
（主に、 $\pi^+$  の観測）

（主に、 $K^+$  の観測）

（主に、 $\pi^+$  の観測）

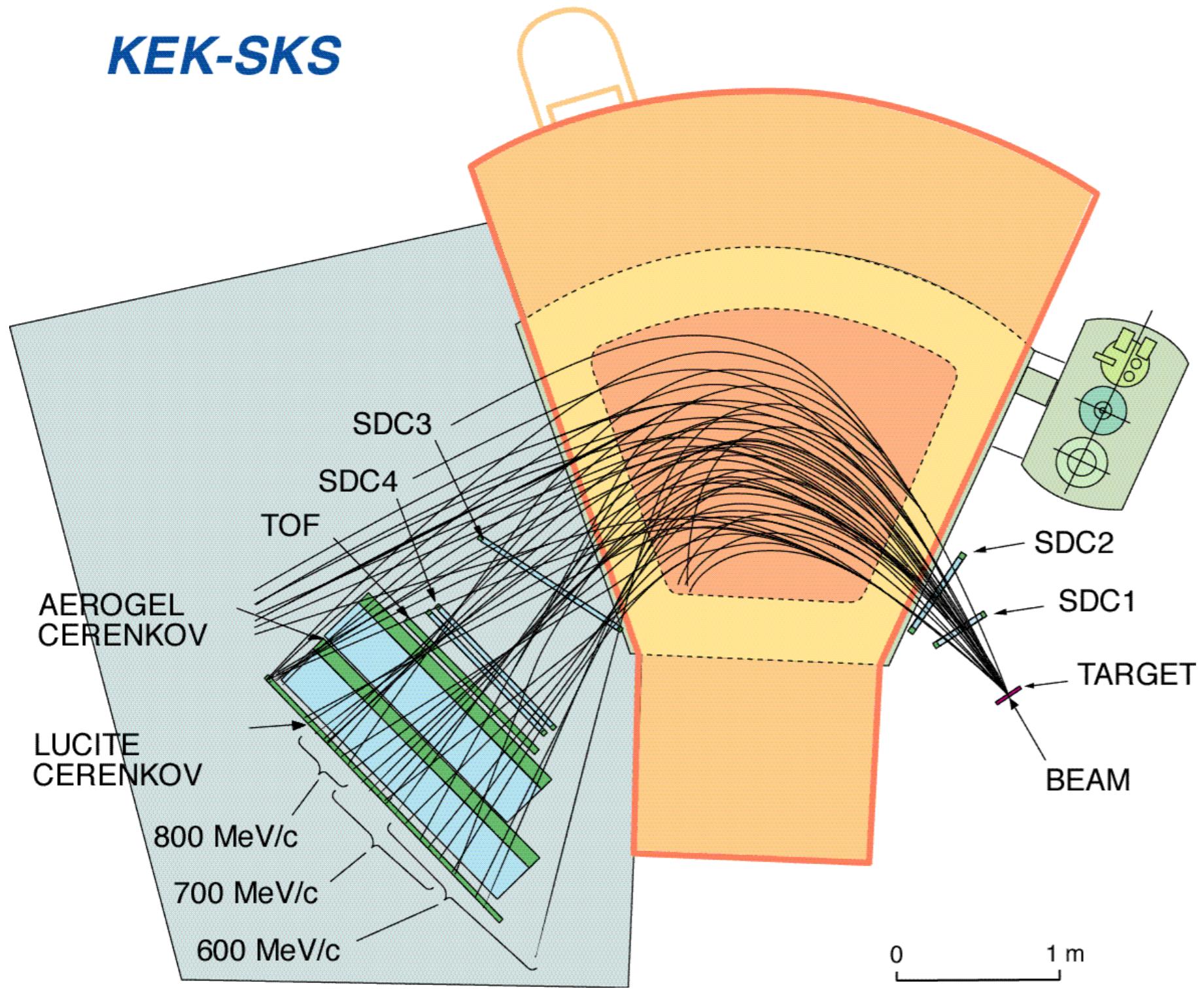
# $(\pi^+, K^+)$ Spectroscopy

- Merits
  - Large momentum transfer  $q \sim 350$  MeV/c
  - Efficiently produces deeply-bound states
  - Low backgrounds:  $\gamma, n$
- Demerits
  - No difference in angular distributions



# SKS Spectrometer at KEK-PS

**KEK-SKS**



- Superconducting Kaon Spectrometer for the ( $\pi^+$ ,  $K^+$ ) reactions
- Constructed by INS, Univ. of Tokyo, from 1987 to 1990
- In operation since 1992
- $B_{max}=3T(500A)$
- Pole Gap=50 cm
- 10.6 MJ stored
- Cold Mass ~4.5 t
- ~280 tons

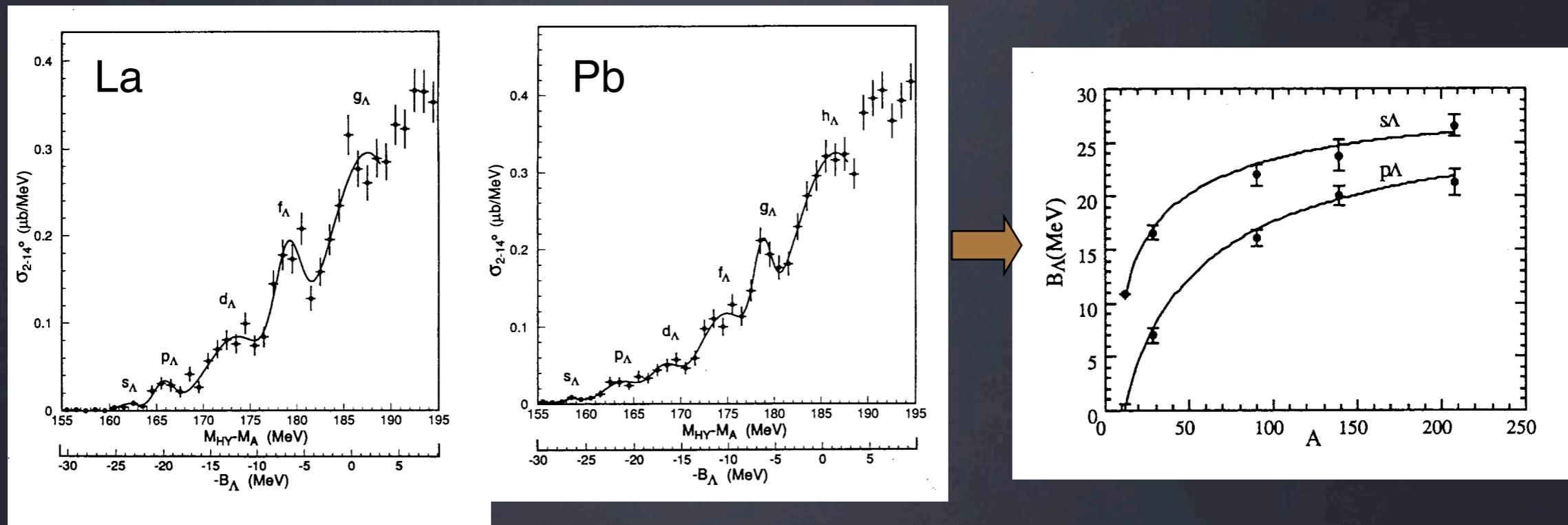
# $(\pi, K^+)$ experiments with SKS

- E140a:  $^{10}\text{B}$ ,  $^{12}\text{C}$ ,  $^{28}\text{Si}$ ,  $^{89}\text{Y}$ ,  $^{139}\text{La}$ ,  $^{208}\text{Pb}$ 
  - Phys. Rev. C 53 (1996) 1210.
- E336:  $^7\text{Li}$ ,  $^9\text{Be}$ ,  $^{13}\text{C}$ ,  $^{16}\text{O}$ 
  - Nucl. Phys. A 639 (1998) 93c, Nucl. Phys. A 691 (2001) 123c.
  - Prog. Part. and Nucl. Phys. 57 (2006) 564.
- E369:  $^{89}\text{Y}$ ,  $^{51}\text{V}$ ,  $^{12}\text{C}$  in high-resolution
  - Phys. Rev. C 64 (2001) 044302.
- E521:  $^{10}\text{B}(\pi^-, K^+)$ 
  - Phys. Rev. Lett. 94 (2005) 052502.

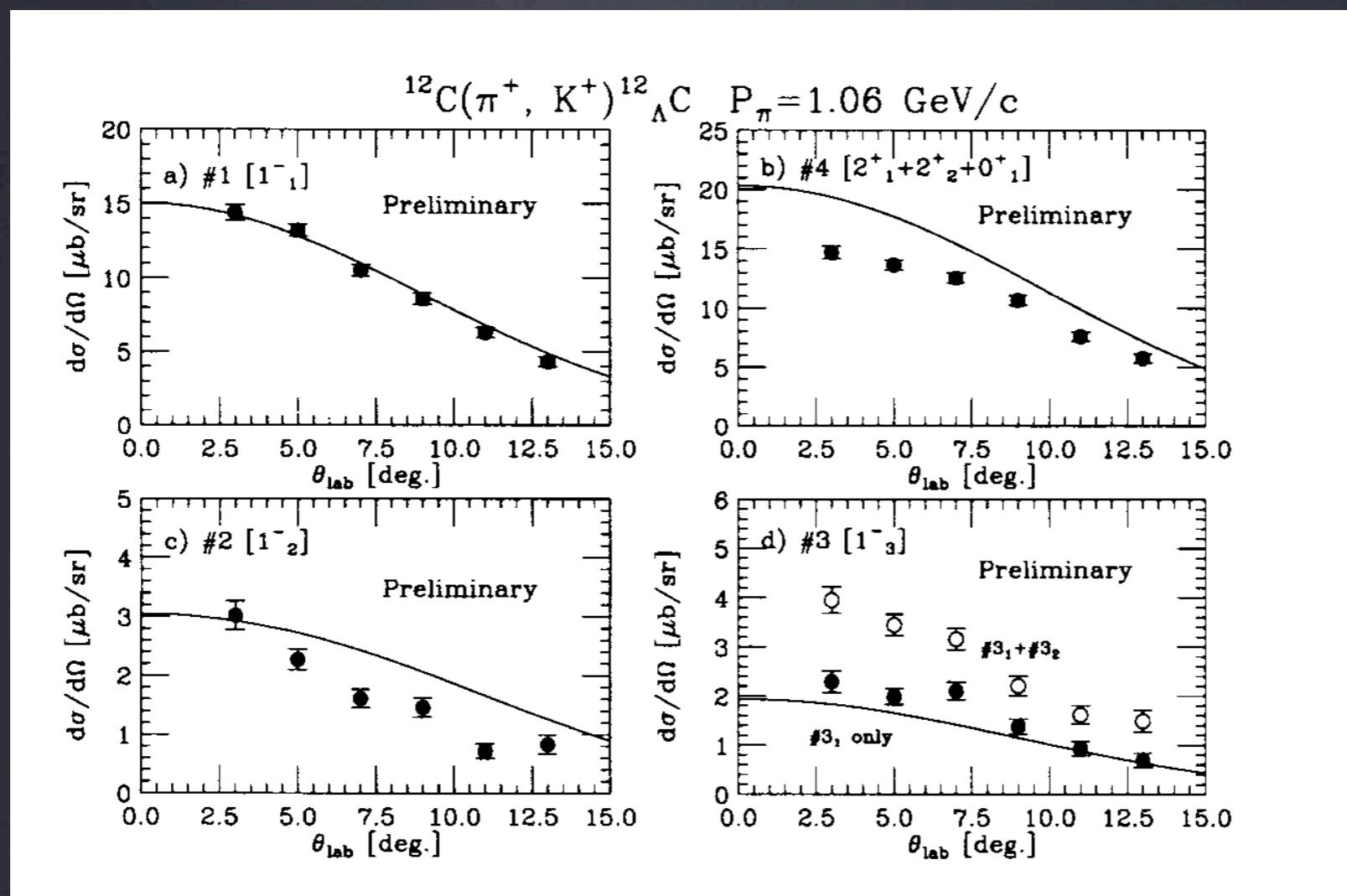
# E140a:

## First ( $\pi^+, K^+$ ) exp. with the SKS

- Targets:  $^{10}\text{B}$ ,  $^{12}\text{C}$ ,  $^{28}\text{Si}$ ,  $^{89}\text{Y}$ ,  $^{139}\text{La}$ ,  $^{208}\text{Pb}$ 
  - $^{12}\Lambda\text{C}$ : First observation of core-excited states
  - Confirmed  $\Lambda$  Shell Structures up to  $^{208}\Lambda\text{ Pb}$



# E336: $^{12}\Lambda\text{C}$ : 角度分布

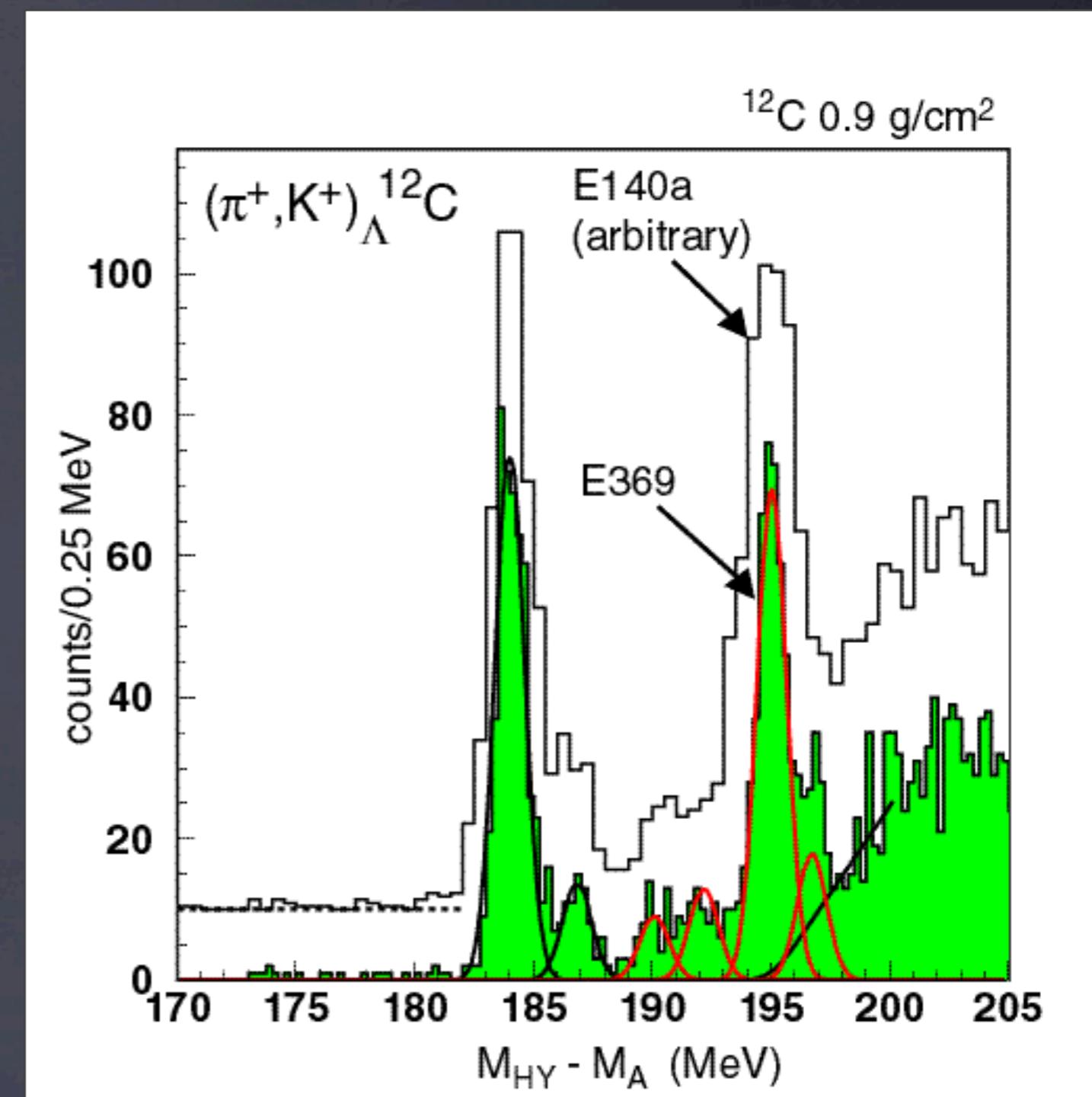


T.Takahashi et al., Nucl. Phys. A670 (2000) 265c.  
Cal. by K.Itonaga et al., Phys. Rev. C49 (1994) 1045.

# E369: $\Lambda^{12}\text{C}$

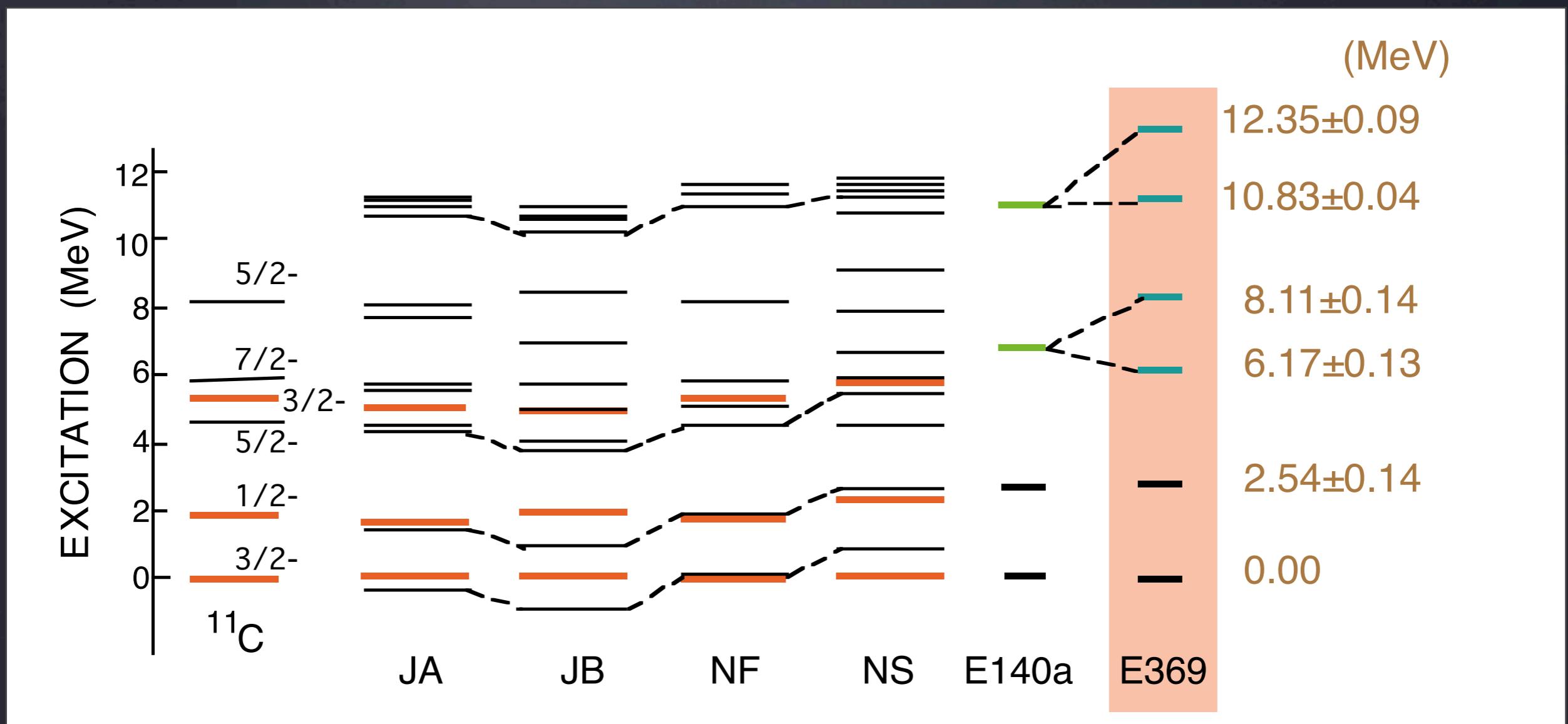
- Best energy resolution
- $\Delta E(\text{FWHM}) = 1.45 \text{ MeV}$

↑  
2.0 MeV



# Core-excited states of $^{12}_{\Lambda}\text{C}$

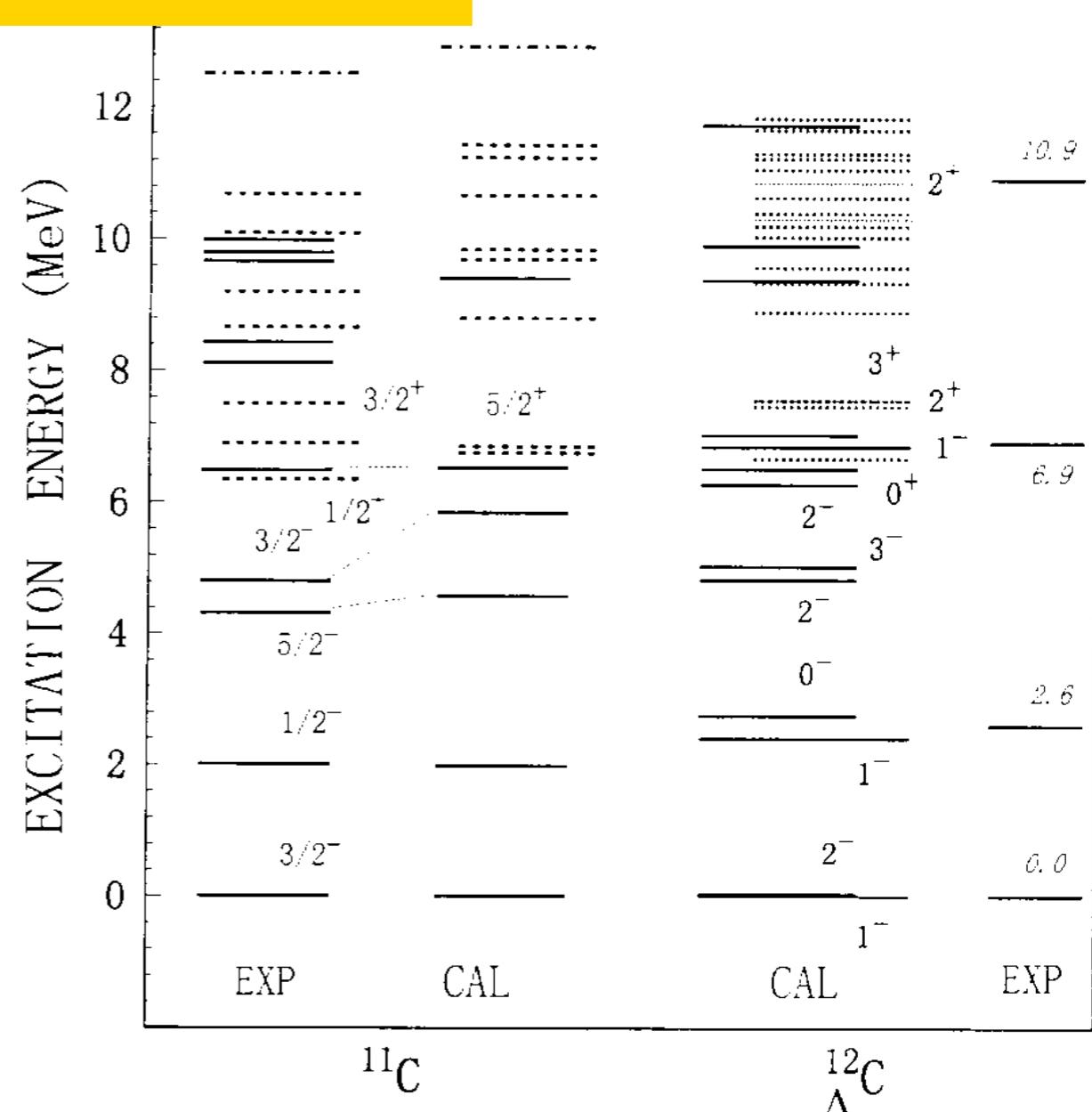
New states are resolved.  
Effects of  $\Lambda\text{N}$  spin-dependent forces



# Parity-mixing intershell coupling

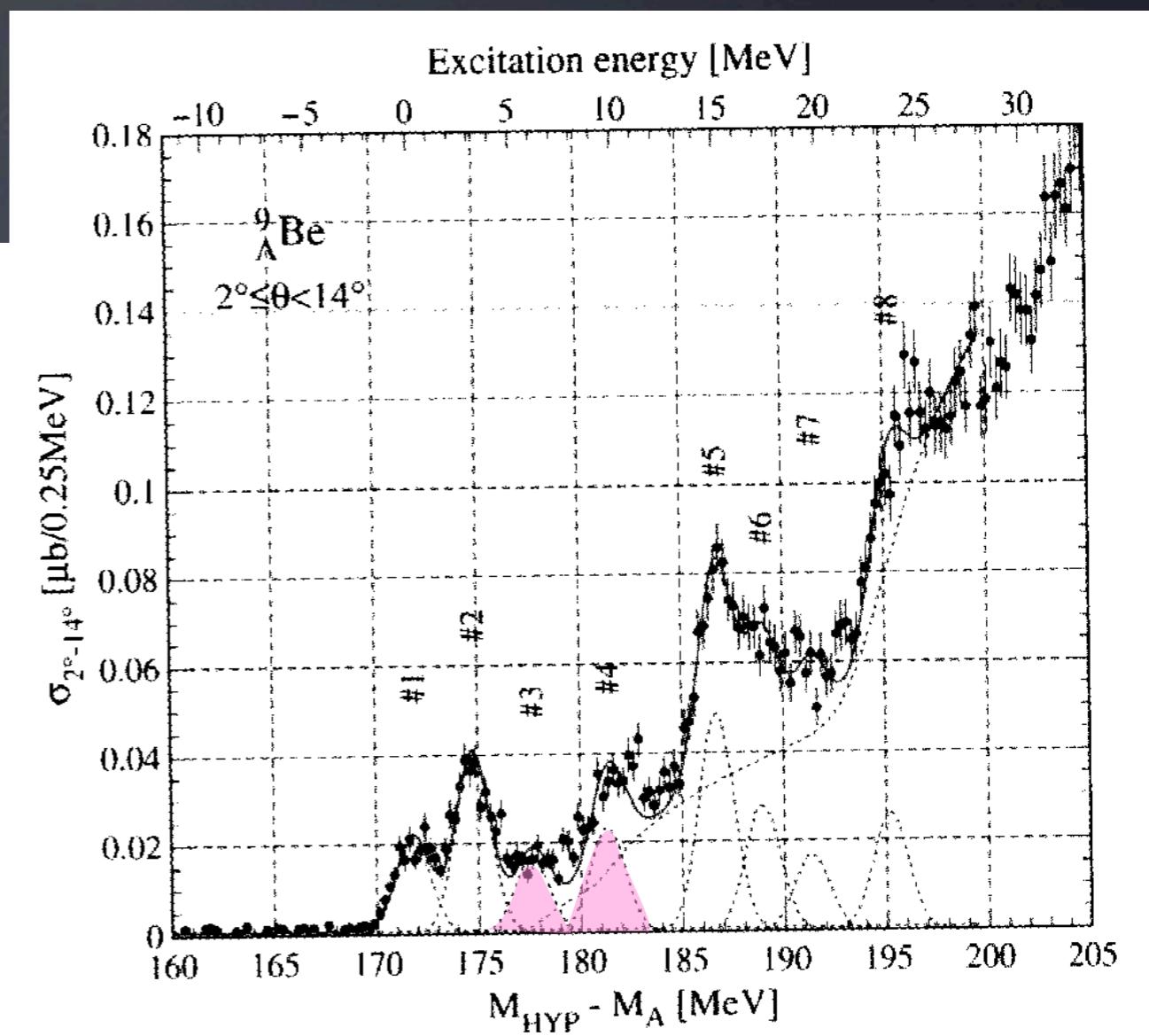
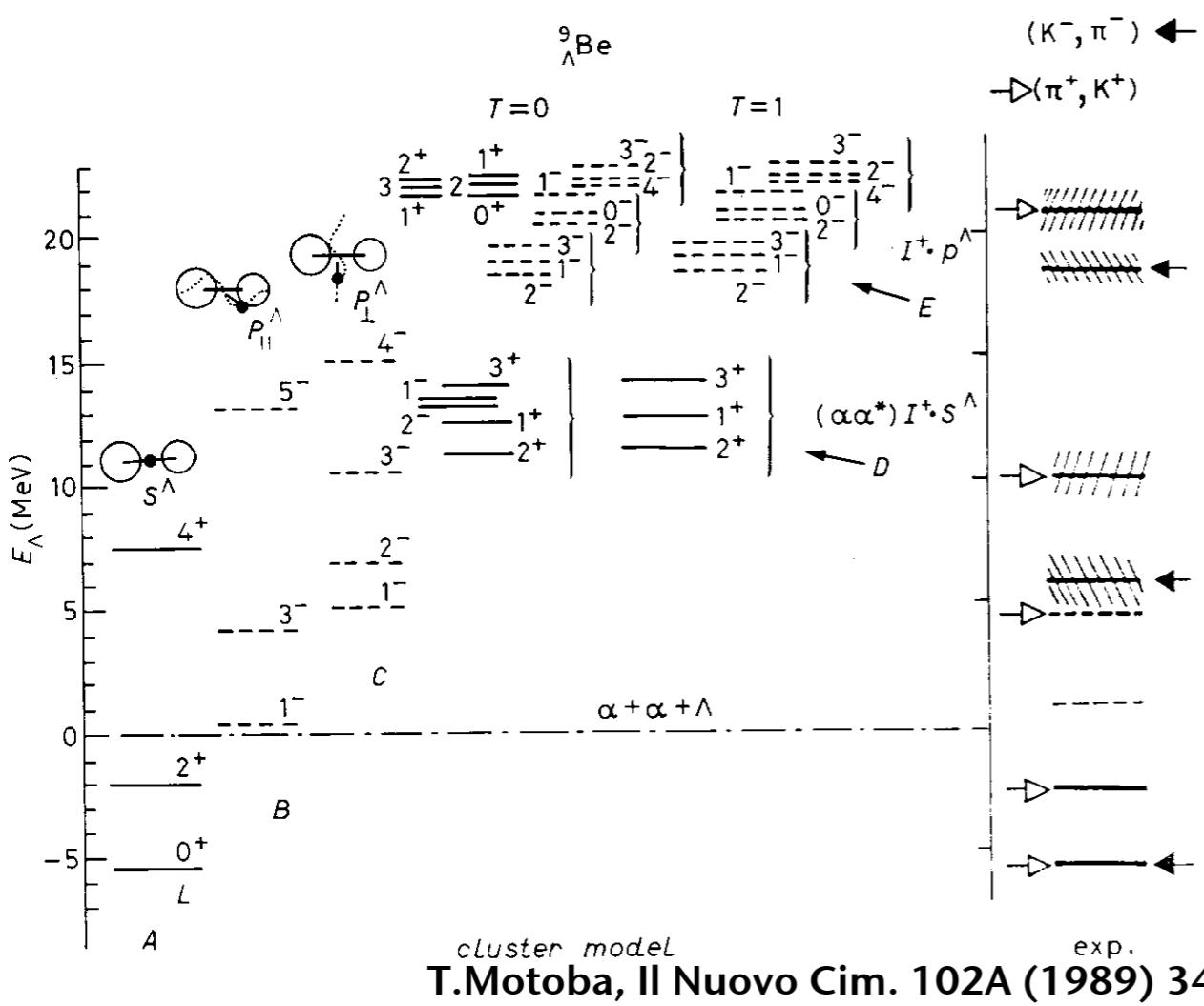
- T. Motoba, in HYP97(Nucl. Phys. A 639 (1998) 135c.)

$$|_{\Lambda}^{12}\text{C}; J^+ \rangle = [s^4 p^7]_- \otimes 1p^{\Lambda} + \{ [s^4 p^6 (sd)^1]_+ + [s^3 p^8]_+ \} \otimes 1s^{\Lambda}$$

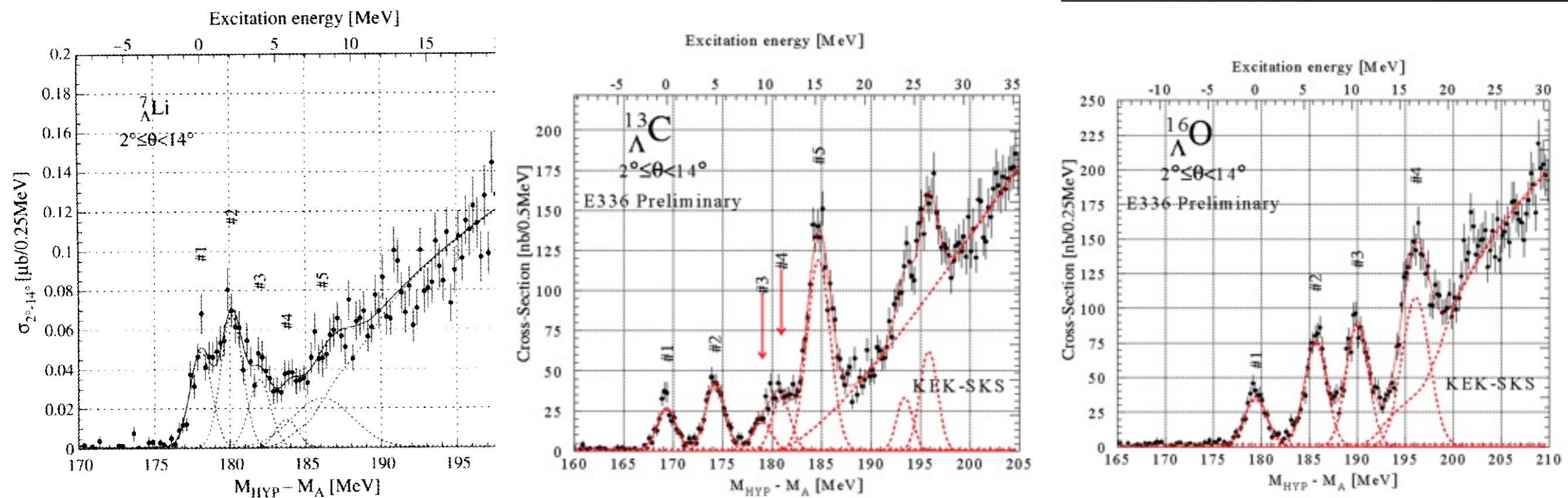


# E336: ${}^9\Lambda\text{Be}$

- Observation of “genuine” hypernuclear states or “Supersymmetric” states



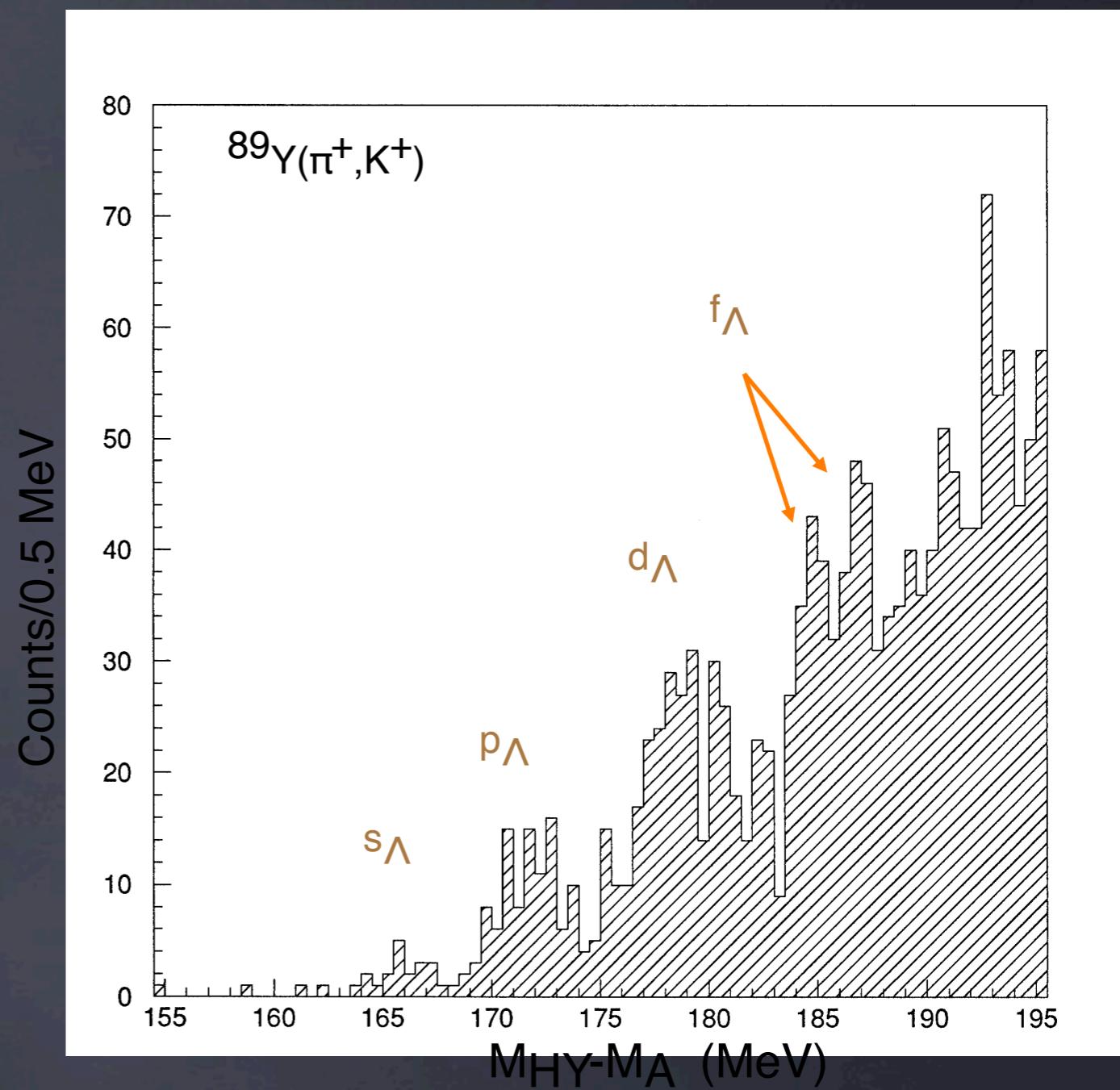
# E336: $\Lambda$ <sup>7</sup>Li, $\Lambda$ <sup>13</sup>C, $\Lambda$ <sup>16</sup>O



$$\Delta E(2^+ - 0^+) = 0.04 \pm 0.32 \text{ MeV}$$

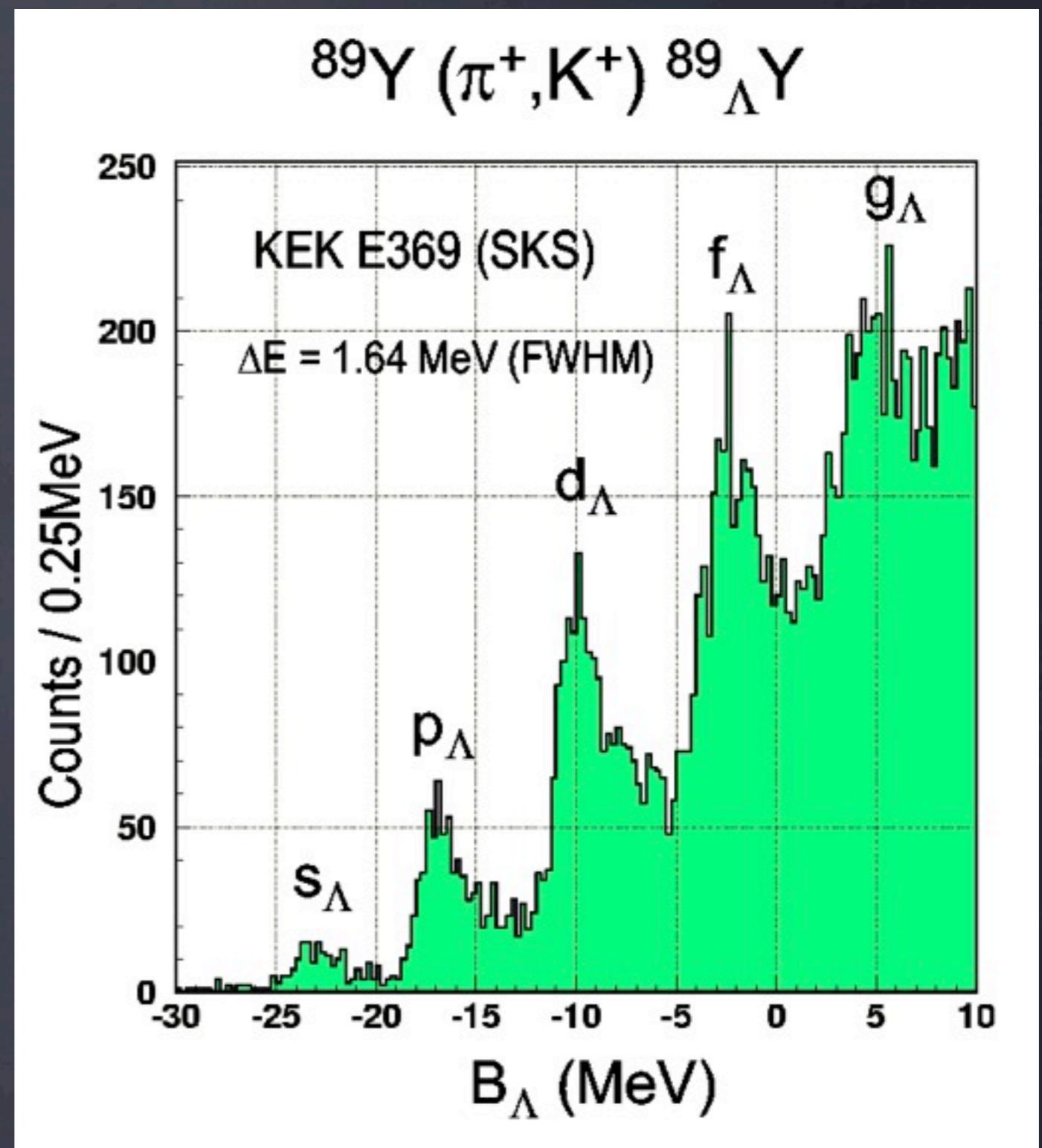
# $^{89}\Lambda Y$ : in E140a

■ KEK-PS E140a  
 $\Delta E = 2.2 \text{ MeV}$



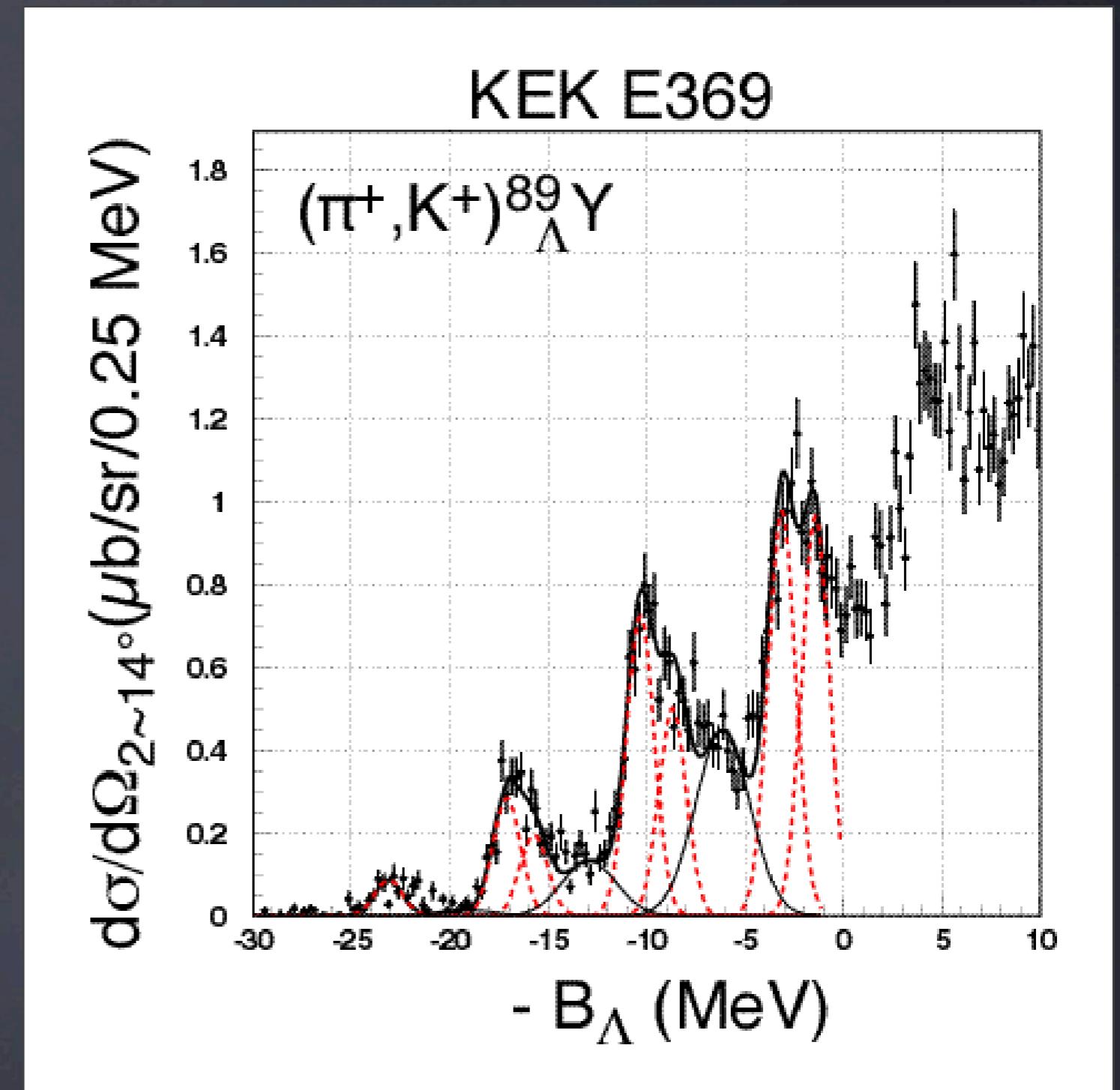
# $^{89}\Lambda Y$ : in E369

- Single Gaussian( $\sigma$ )
- p:  $2.4 \pm 0.2$  MeV
- d:  $3.0 \pm 0.2$  MeV
- f:  $4.6 \pm 0.5$  MeV



# E369: $\Lambda^{89}\text{Y}$

- $B_{\Lambda s} = 23.1 \pm 0.1$  MeV
- Energy Splitting
  - $\Delta E_f = 1.70 \pm 0.10$  MeV
  - $\Delta E_d = 1.63 \pm 0.14$  MeV
  - $\Delta E_p = 1.37 \pm 0.20$  MeV
- Peak Ratio
  - $R/L_f = 0.99 \pm 0.07$
  - $R/L_d = 0.69 \pm 0.06$
- Extra n-hole at  $+4.1 \pm 0.1$  MeV,  
width= $3.2 \pm 0.2$  MeV



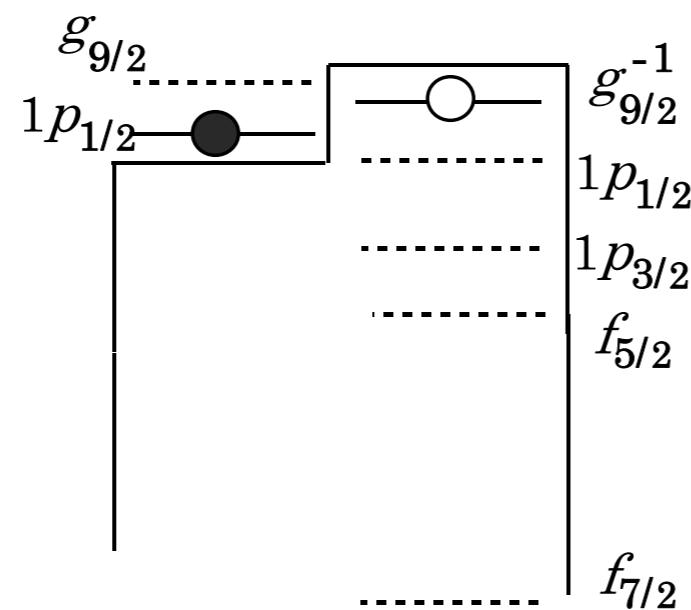
- SplittingがLSによるものとすると、大きすぎる。
  - Splittingは $(2\ell+1)$ に比例するべきだが、これと異なる。
  - L/Rのピーク比が、Kinematical Factorからずれている。

Table 1

Experimental energy spacings ( $\Delta E_i$ ) and cross sections ( $\bar{\sigma}$ , averaged over  $\theta_K^{\text{Lab}} = 2^\circ - 14^\circ$ ) for a series of major peaks (Left and Right subpeaks) in the  $^{89}\text{Y}(\pi^+, K^+)^{89}\text{Y}_{\Lambda}$  reaction [11]. The cross sections and their ratios are compared with theoretical predictions using the DWIA [4]. See the text for NSC97f.

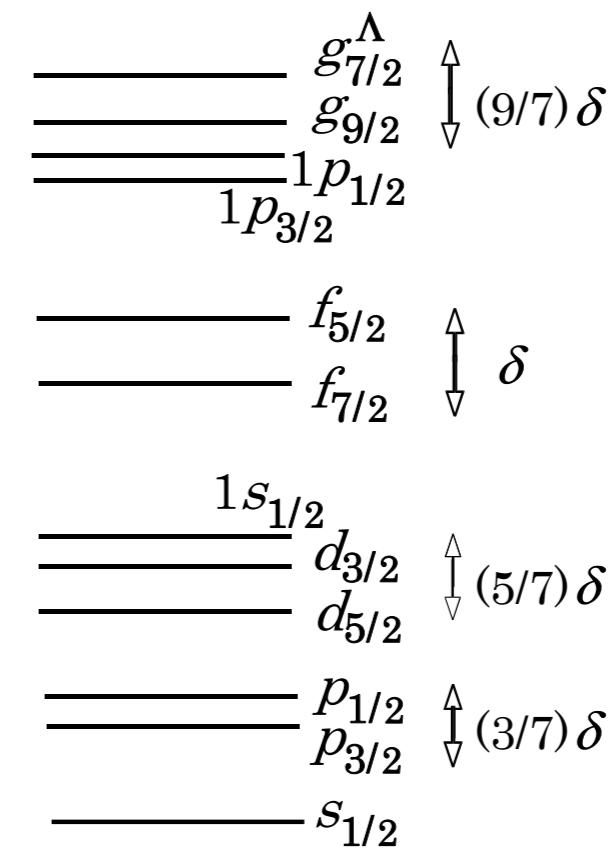
peak	$E_{\Lambda}^{\text{exp}}$	$\bar{\sigma}^{\text{exp}}$	$\Delta E_i$ (MeV)		$\Delta E_i/\Delta E_1$ ratio			$\bar{\sigma}_L/\bar{\sigma}_R$	
	(MeV)	( $\mu\text{b}/\text{sr}$ )	exp	NSC97f	exp	NSC97f	WS	exp	calc
$s_{\Lambda}$	-23.11	0.60							
$p_{\Lambda}$	1L -17.10	2.00							
	1R -15.73	1.38	1.37	0.46	1.00	1.00	1.00	1.45	1.00
$d_{\Lambda}$	2L -10.32	5.10							
	2R -8.69	3.52	1.63	0.72	1.19	1.57	2.16	1.45	0.55
$f_{\Lambda}$	3L -3.13	6.87							
	3R -1.43	6.79	1.70	0.92	1.24	2.00	3.36	1.01	0.44

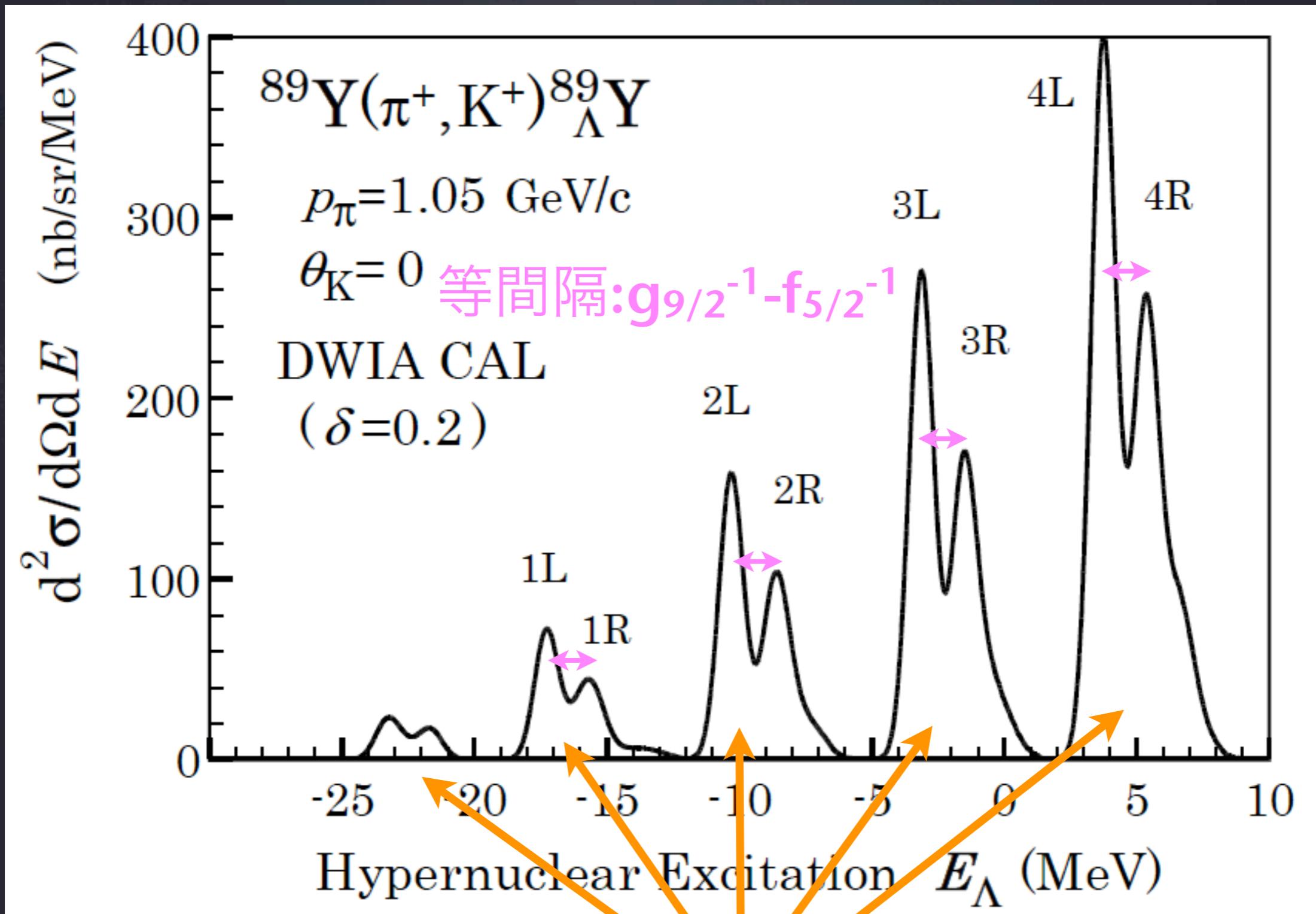
$^{88}\text{Y} (j_{\text{p}} j_{\text{n}}^{-1})_{J_{\text{c}}}$



$[(j_{\text{p}} j_{\text{n}}^{-1})_{J_{\text{c}}} \; j^{\Lambda}]_{J_{\text{H}}}$

$\Lambda$  orbits ( $j^{\Lambda}$ )





形が相似形

# E369: $\Lambda^{51V}$

Splitting in d-(and p-) orbit(s)

$$B_{\Lambda} = (20 \pm 0.13) + 0.56 \text{ MeV}$$

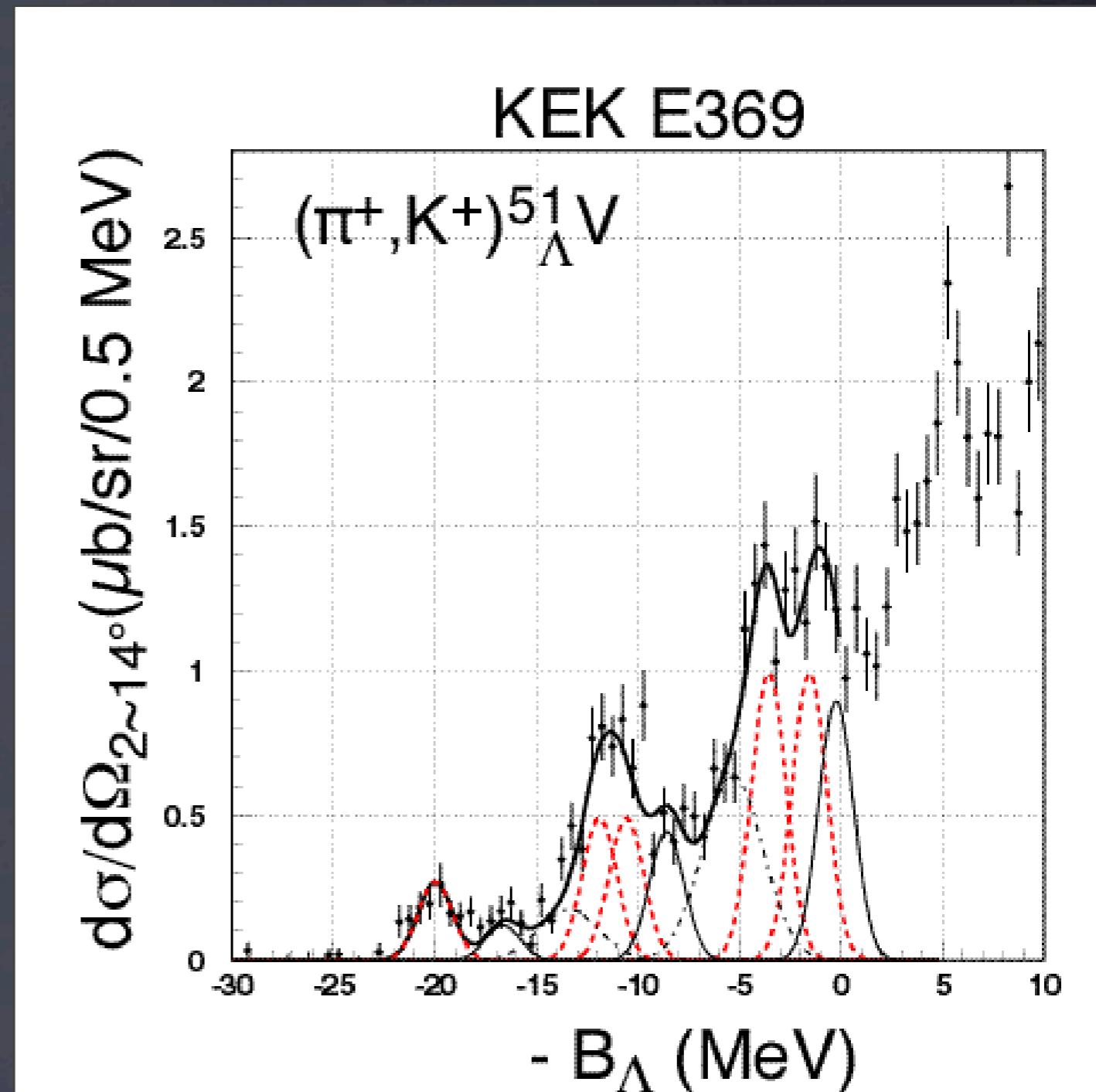
Width=1.95 MeV

Peak Ratio=1(fixed)

Extra n-holes

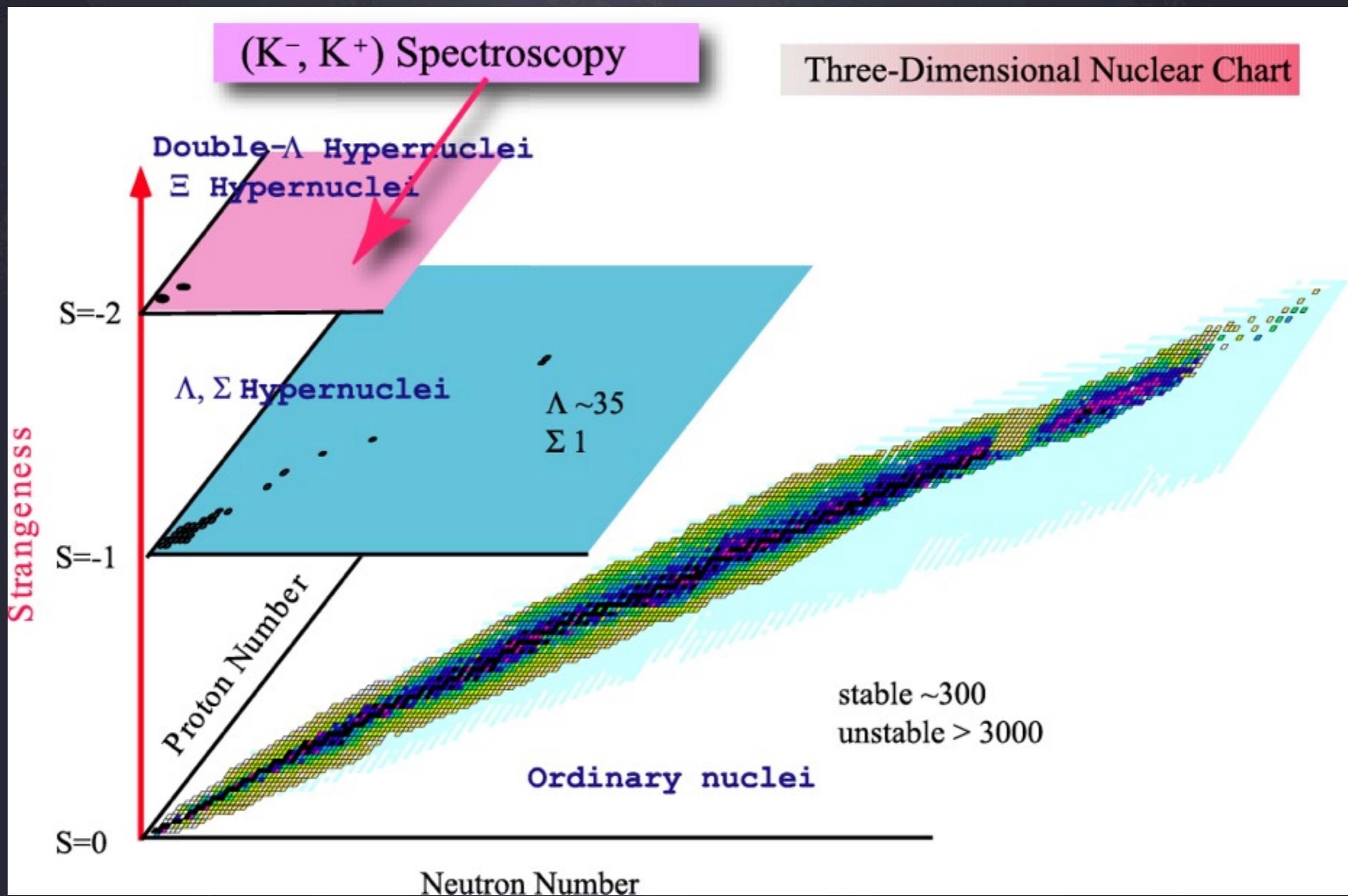
At  $+3.3 \pm 0.2$  MeV, width=1.95 MeV

At  $+6.6 \pm 0.2$  MeV, width=3.46 MeV



$S=2\wedge : (K^-, K^+)$

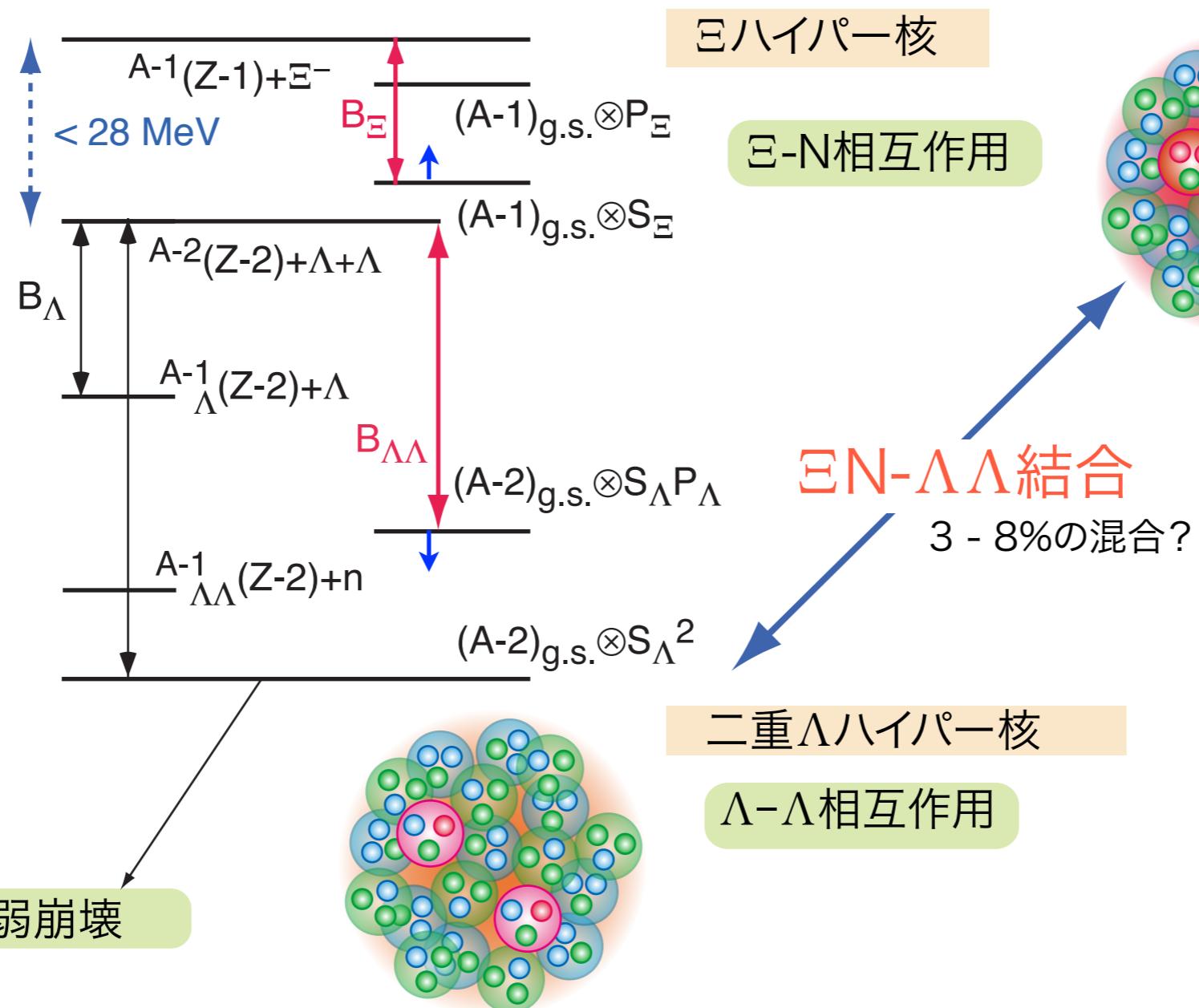
# Strangeness Nuclear Physics



# $S=-2$ 多体系

- バリオン混合の重要性

$S=-2$ のバリオン多体系のエネルギースペクトル

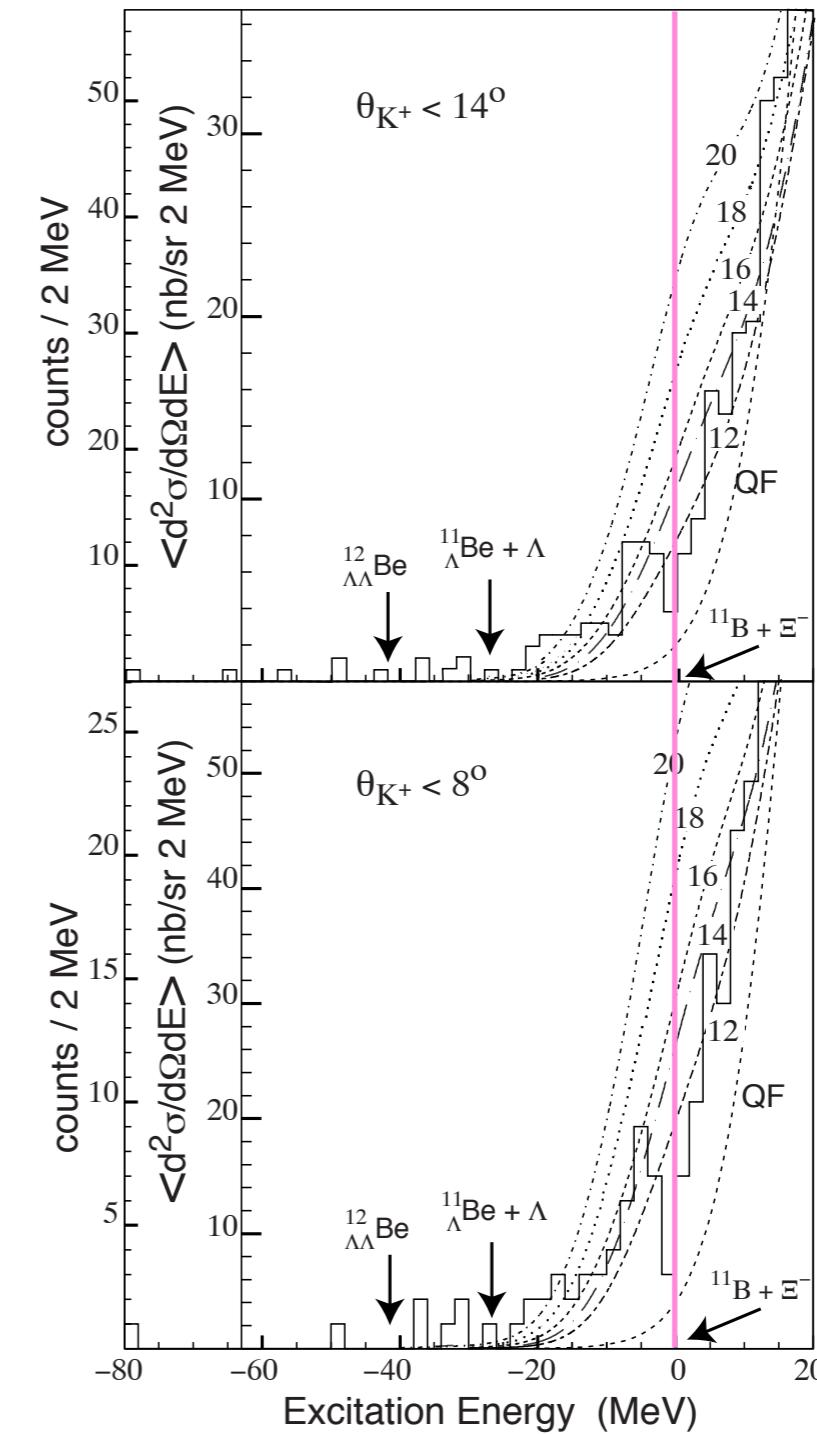


# 過去の実験データ

$^{12}\text{C}(\text{K}^-, \text{K}^+)$

- BNL-E885
  - Evidence ?
  - $U_{\Xi} = -14$  MeV
  - $89 \pm 14$  nb/sr
- $\Delta E \sim 14$  MeV(FWHM)
- $\Delta \Omega \sim 19$  msr

P. Khaustov et al., PRC61(2000) 054603.

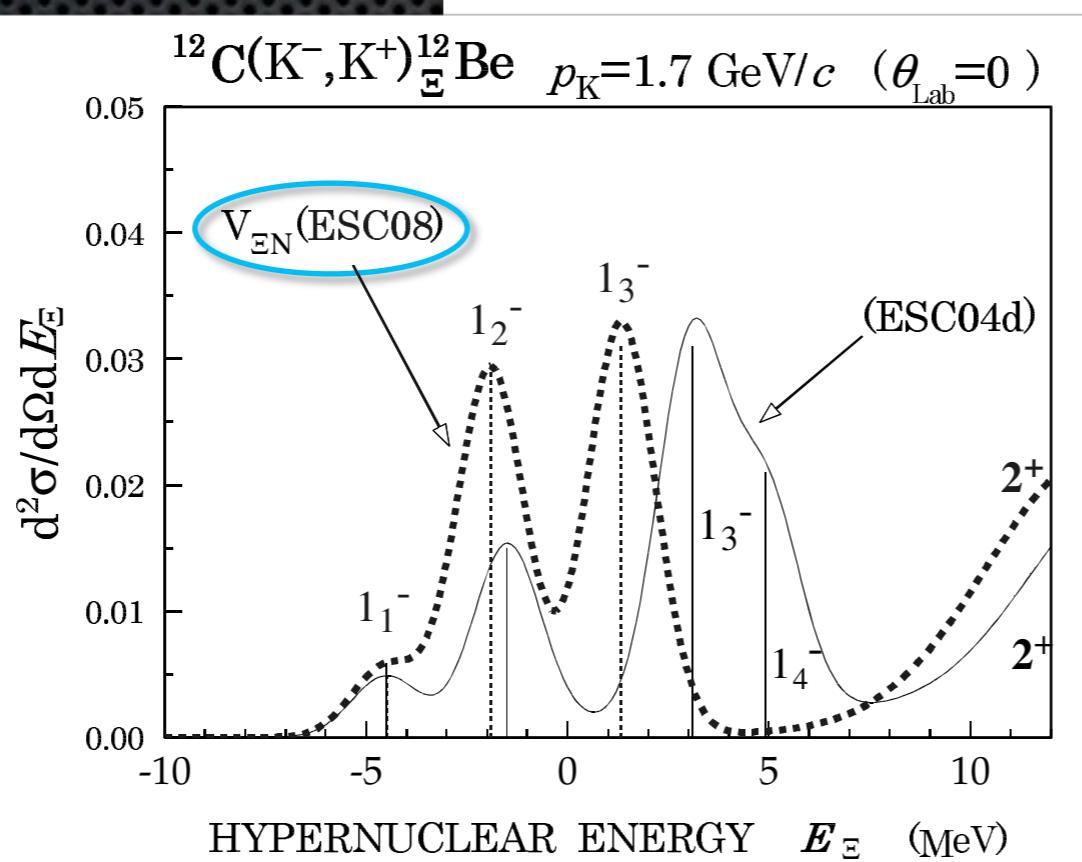


# $\Xi N$ 相互作用

- $U_\Xi \sim -15 \text{ MeV}$

Table 3:  $U_\Xi(\rho_0)$  and partial wave contributions.

model	T	$^1S_0$	$^3S_1$	$^1P_1$	$^3P_0$	$^3P_1$	$^3P_2$	$U_\Xi$	$\Gamma_\Xi$
ESC08a	0	6.0	-1.0	-0.3	-2.6	1.3	-0.9	-20.2	5.8
	1	8.5	-28.0	0.6	0.4	-3.7	-0.6		
ESC08a'	0	5.6	-1.1	-0.3	-2.6	1.3	-0.9	-14.5	7.0
	1	8.4	-21.5	0.6	0.4	-3.7	-0.6		
ESC08b	0	2.4	1.9	-0.6	-1.2	-0.1	-0.7	-31.8	0.9
	1	9.1	-37.8	0.6	-0.5	-3.6	-1.3		
ESC04d	0	6.4	-19.6	1.1	1.2	-1.3	-2.0	-18.7	11.3
	1	6.4	-5.0	-1.0	-0.6	-1.4	-2.8		

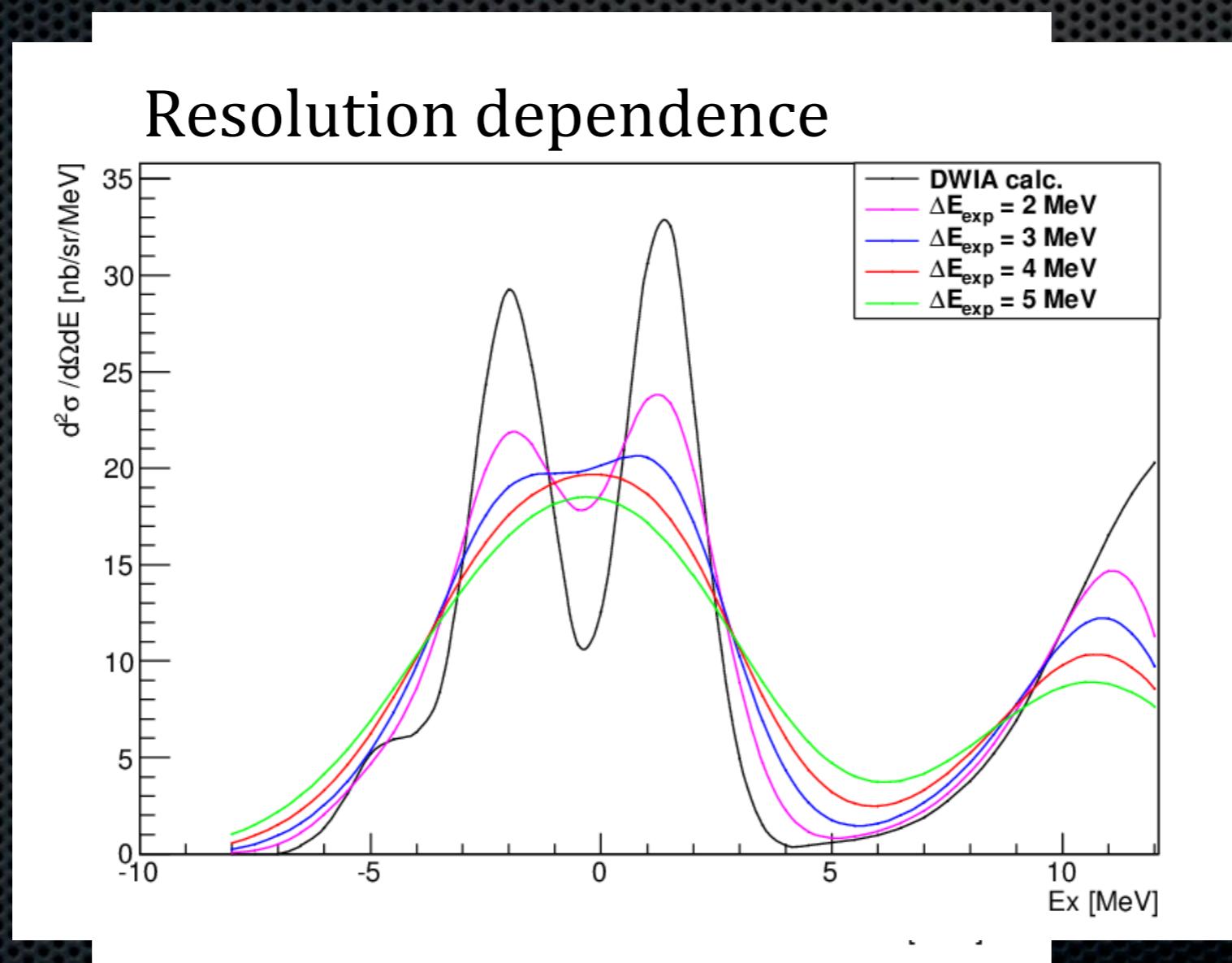


$T = 1$

DWIA spectrum from  
*T. Motoba and S. Sugimoto,*  
*NPA 835 (2010) 223.*

# 必要な分解能

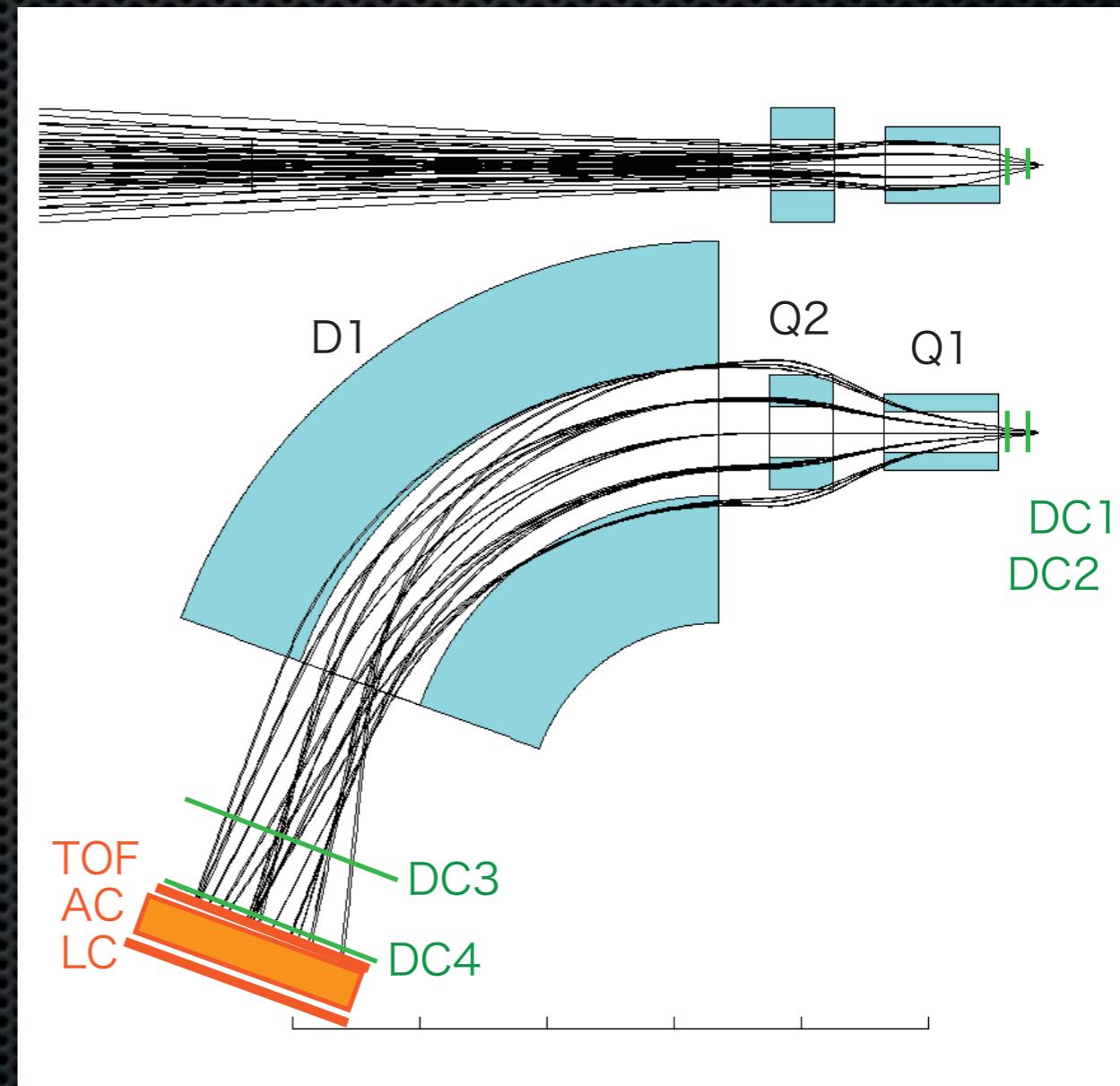
- 2 MeV(FWHM)が必要



# 新しいスペクトロメーター

- S-2Sの建設(H23-H26): 特別推進
  - QCD構成
  - 立体角~60 msr
  - 運動量領域 :  $\pm 10\%$
  - $\Delta p/p < 5 \times 10^{-4}$ (FWHM)
  - $\Delta E = 1.2$  MeV

	立体角 $\Delta\Omega$ (msr)	エネルギー分解能 $\Delta E$ (MeV)
BNL	19	14
SKS+	25	3
S-2S	60	1.2

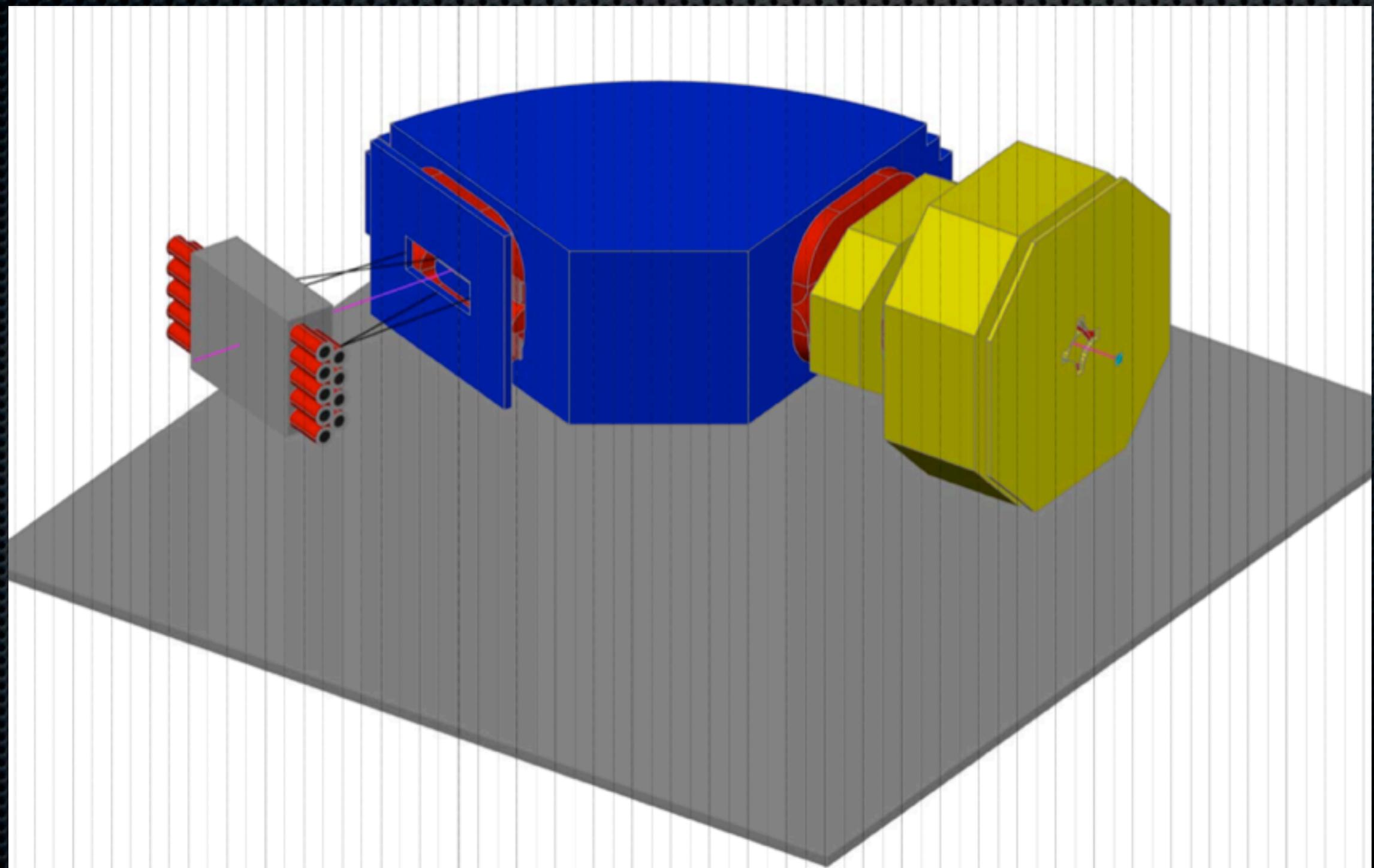


S-2S

$$\Delta p/p = 5 \times 10^{-4}$$

加藤、高橋仁、高橋俊、藤岡、広瀬

60msr

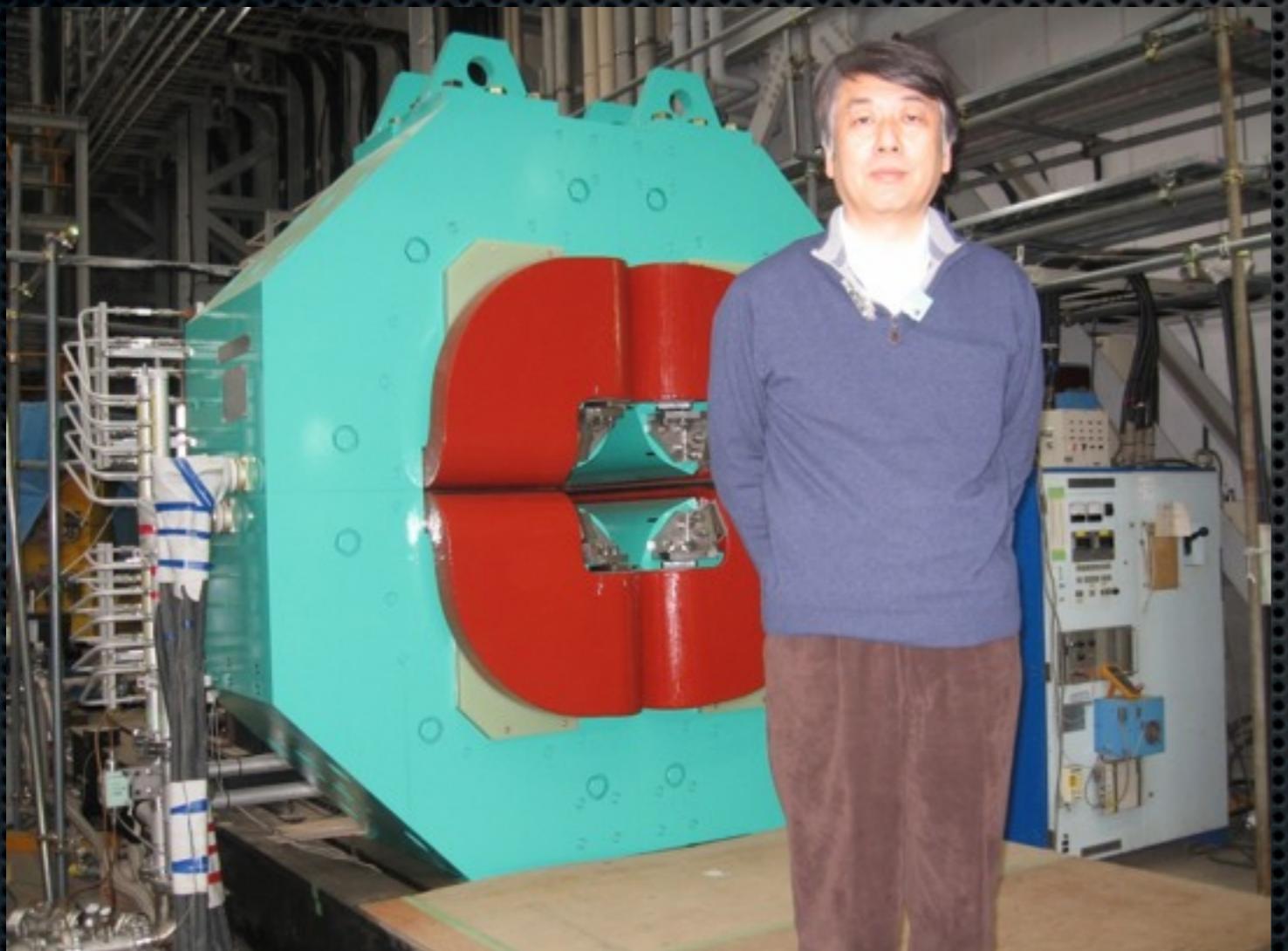


# Q1 of S-2S

Feb. 15, 2013

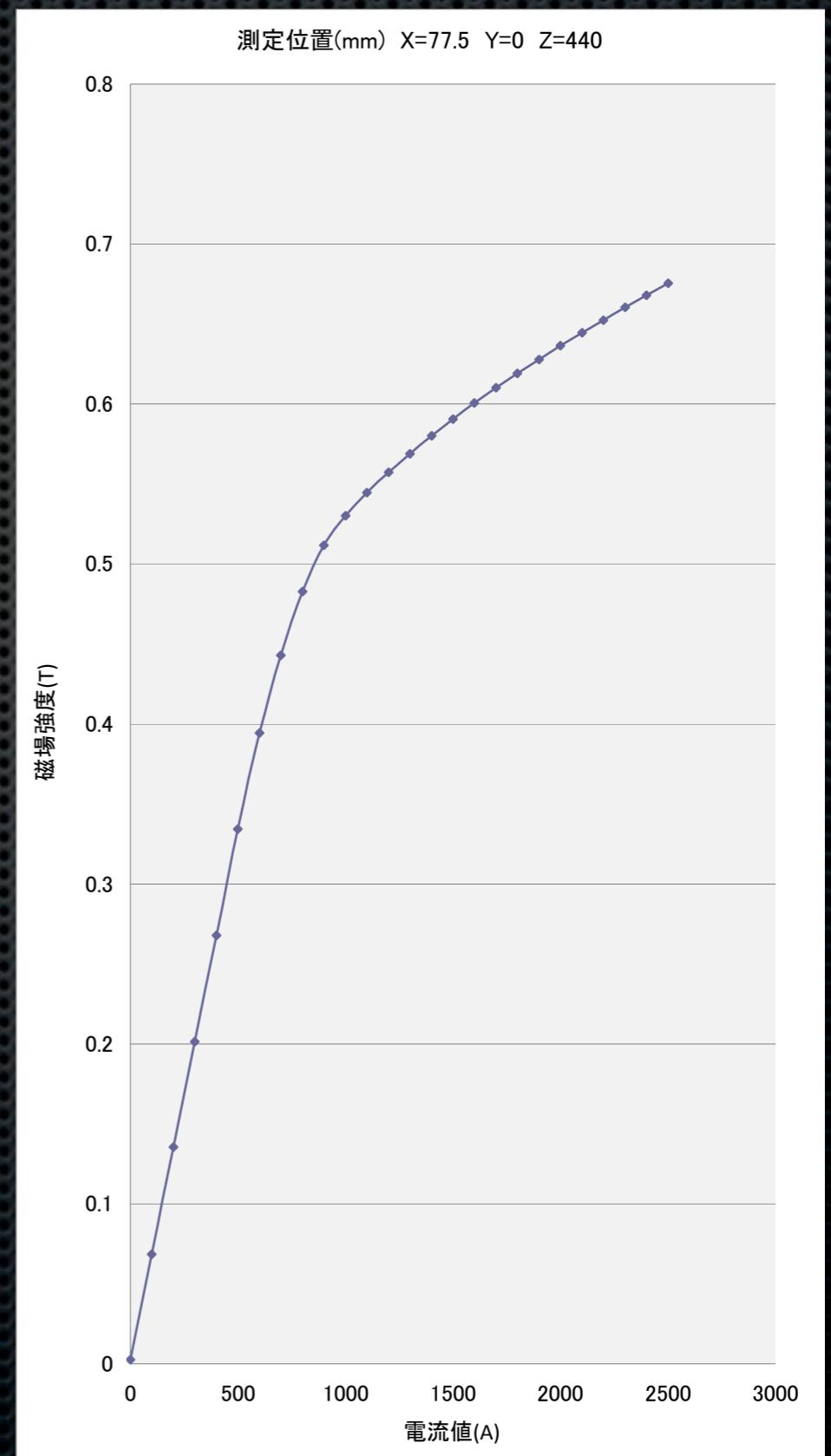


# Q1励磁テスト



KEK 北カウンターホール於

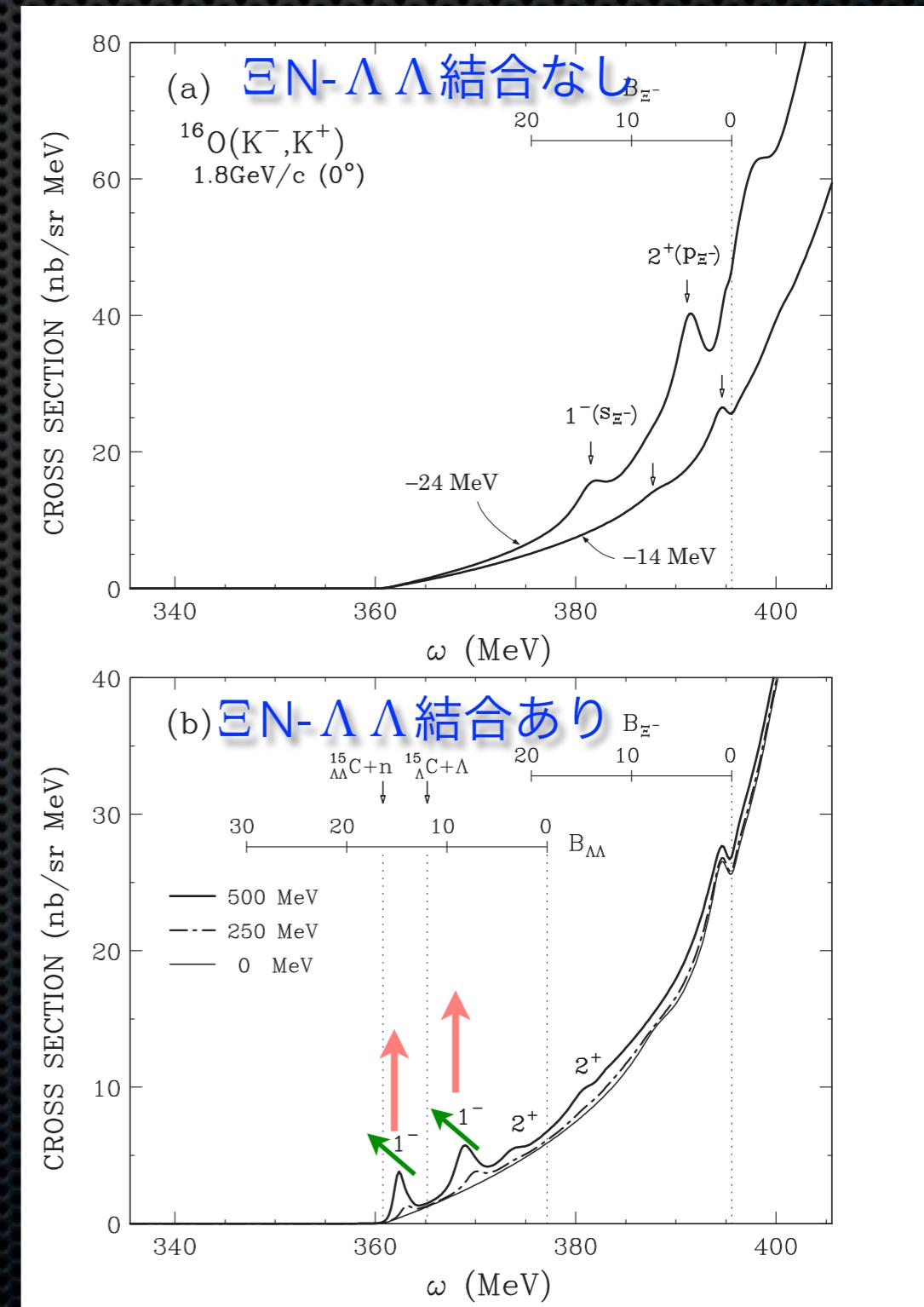
励磁特性



# 期待される成果 H26-H27

- 二重 $\Lambda$ ハイパー核と $\Xi$ ハイパー核の励起
- $\Xi N - \Lambda\Lambda$ 結合が重要な役割
  - 強く励起 ↑
  - 深く束縛 ↘
- $[^{15}\text{N}(1/2^-, 3/2^-) \otimes S_\Xi]_{1^-} \rightarrow [^{14}\text{C}(0^+, 2^+) \otimes S_\Lambda p_\Lambda]_{1^-}$
- $[^{15}\text{N}(1/2^-, 3/2^-) \otimes p_\Xi]_{2^+} \rightarrow [^{14}\text{C}(0^+, 2^+) \otimes p_\Lambda^2]_{2^+}$
- ピークエネルギーの高精度測定(<100 keV)
  - $\Xi N - \Lambda\Lambda$ 結合によるエネルギーのシフト  
~1.17 MeV
- $S_\Lambda - p_\Lambda$ 間相互作用: →初めての測定

原田らの計算例 PLB690 (2010) 363.  $^{16}\text{O}(K^-, K^+)$



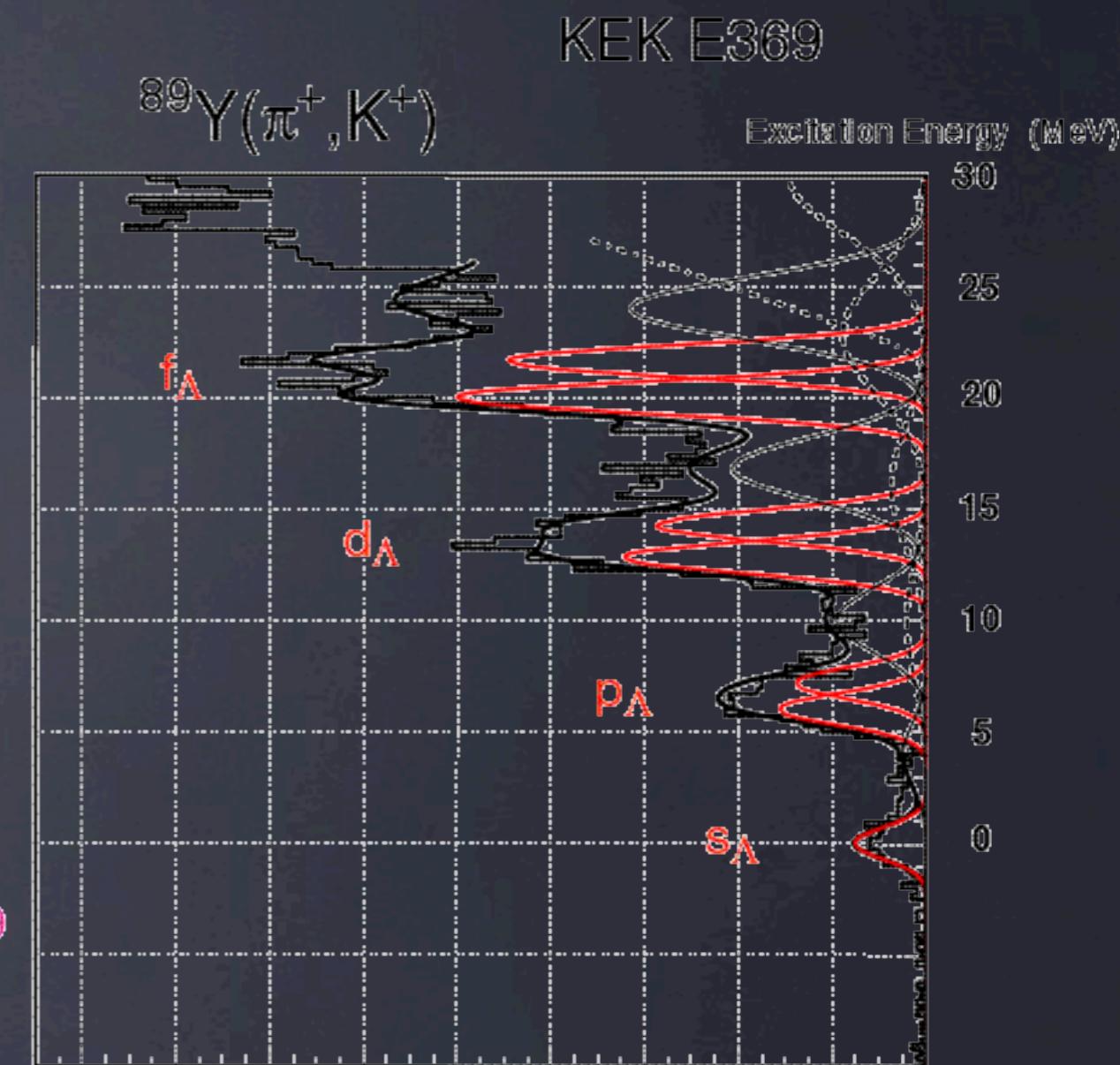
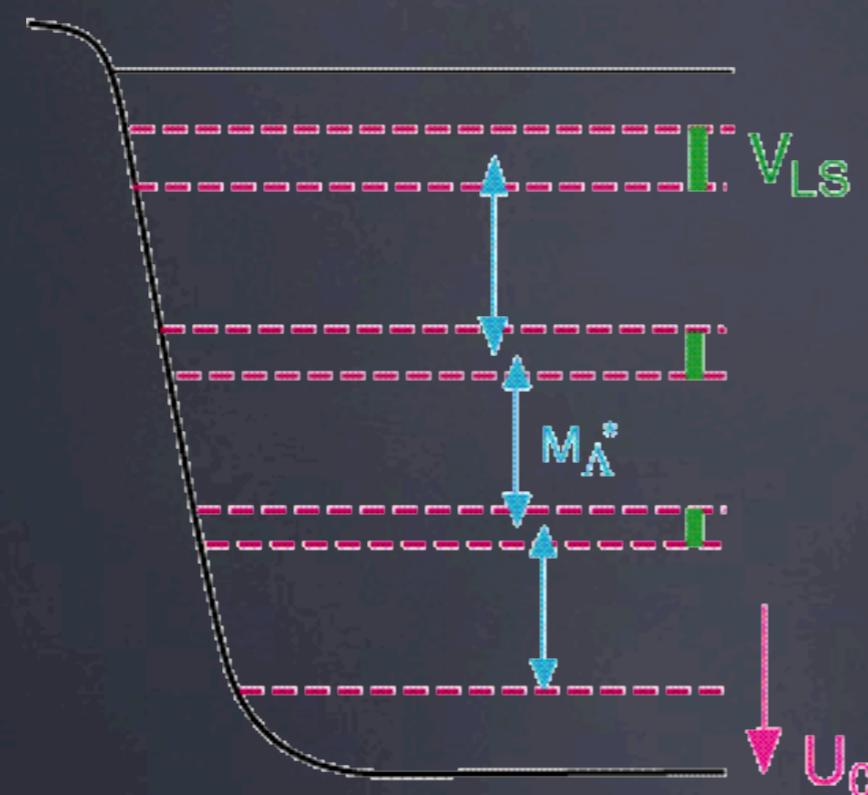
# まとめ

- SKSスペクトロメーター
  - 15年間にわたって世界をリード。多くの若手研究者を輩出。
- ( $\pi, K$ )スペクトロスコピー
  - 広い質量領域でのラムダ一粒子軌道の観測
  - コア励起状態、Genuineハイパー核状態等の観測
  - 元場さんを中心とする理論の助けのおかげ
- $S=-1$ から  $S=-2$ へ
  - ( $K^-, K^+$ )スペクトロスコピーのための  $S-2S$  の建設
  - $\Xi$ ハイパー核の探索。 $\Xi N - \Lambda \Lambda$  結合。



# Single-particle motion of $\Lambda$ in heavy hypernuclei

- $U_0 = -30.5 \text{ MeV}$
- $M\Lambda^* = 0.7 \sim 0.8 \times M\Lambda$



# Effective mass of $\Lambda$ :

## Y.Yamamoto

LDA treatment for  $\Lambda$  hypernuclei

$$\frac{M_\Lambda^*}{M_\Lambda} = \left( \frac{\tilde{M}_\Lambda}{M_\Lambda} \right) \times \left( \frac{\bar{M}_\Lambda}{M_\Lambda} \right)$$

$$\frac{\tilde{M}_\Lambda}{M_\Lambda} = \left[ 1 + 2M_\Lambda \frac{\partial U_\Lambda}{\partial p_\Lambda^2} \right]^{-1}$$

$$\frac{\bar{M}_\Lambda}{M_\Lambda} = 1 - \frac{\partial U_\Lambda}{\partial \epsilon_\Lambda}$$

$$\epsilon_\Lambda(\rho) = \frac{p_\Lambda^2}{2M_\Lambda} + U_\Lambda(\rho, p_\Lambda, \epsilon_\Lambda) \approx \frac{p_\Lambda^2}{2M_\Lambda^*(\rho)} + U_\Lambda^0(\rho)$$

$$\epsilon_\Lambda^0(\rho) = U_\Lambda^0(\rho) = U_\Lambda(\rho, p_\Lambda = 0, \epsilon_\Lambda^0)$$

$$U_\Lambda(r) = \alpha t \sqrt{\pi} r^{-3} \int d\mathbf{r}' U_\Lambda^0(\rho(r')) \exp(-|\mathbf{r} - \mathbf{r}'|^2/t^2) \rho(r')$$

$$M_\Lambda^*(r) = \beta M_\Lambda^*(\rho(r)) + (1 - \beta) M_\Lambda$$

( $\alpha$  and  $\beta$  are additional parameters for fitting)

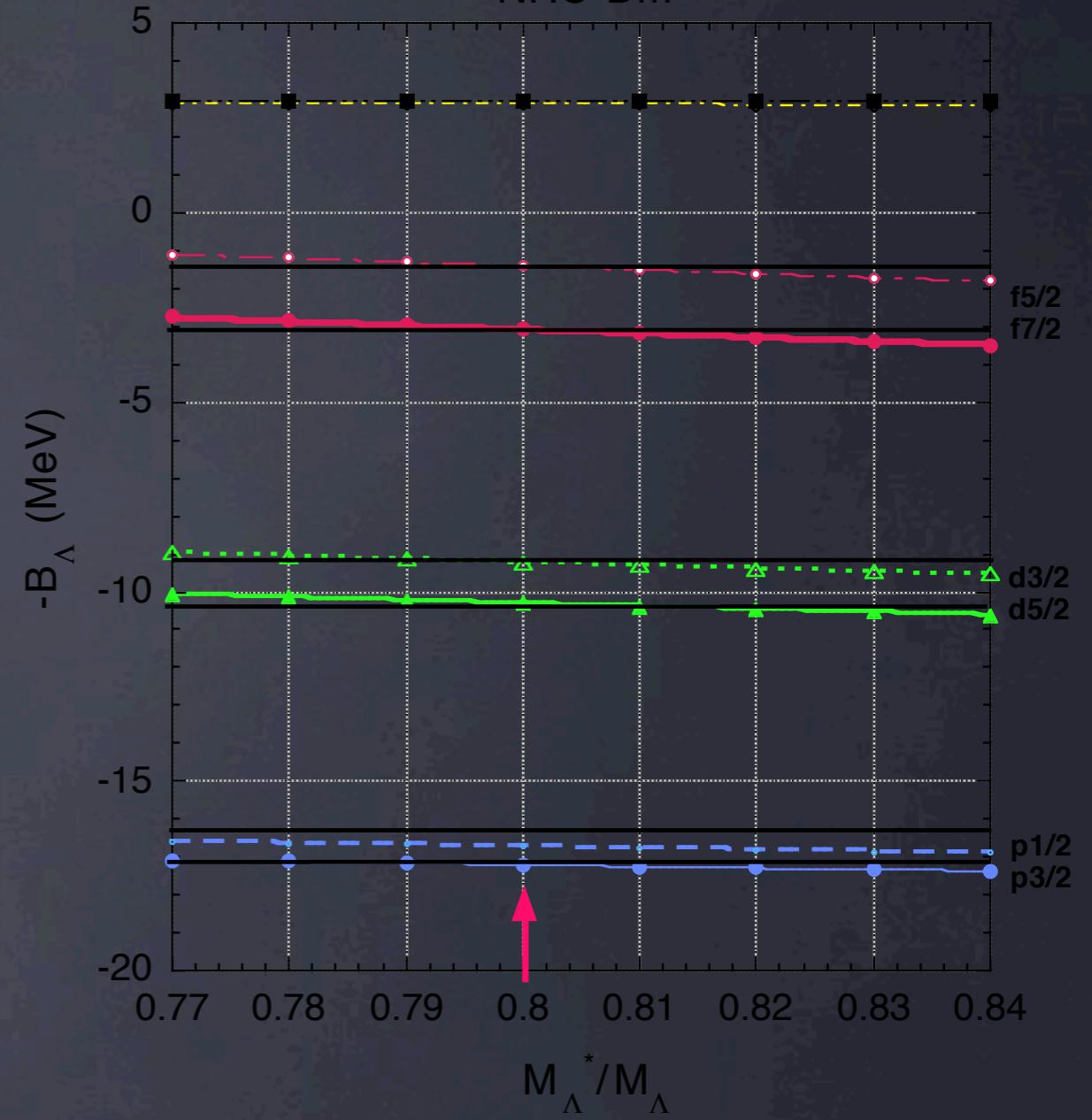
$$U_B^{ls}(r) = K_B \frac{1}{r} \frac{d\rho}{dr} \vec{l} \cdot \vec{\sigma} \quad \text{with } B = N, \Lambda$$

$$\text{where } K_N = -\frac{1}{2}\pi S_{LS} \text{ and } K_\Lambda = -\frac{1}{3}\pi (S_{LS} + S_{ALS})$$

$$S_{LS,ALS} = \frac{3}{\bar{q}} \int_0^\infty r^3 j_1(\bar{q}r) G_{LS,ALS}^3(r) dr$$

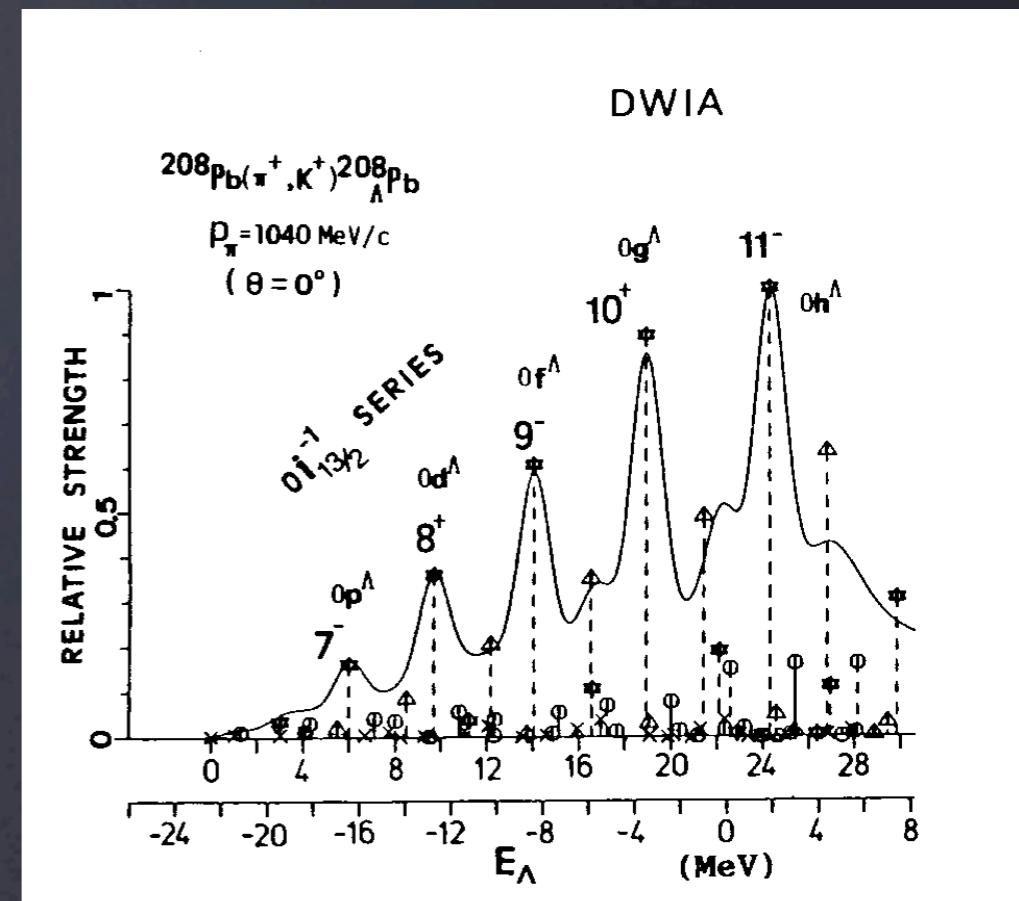
with  $\bar{q} \sim 0.7 \text{ fm}^{-1}$

Calculation by Y. Yamamoto  
NHC-Dm



# Neutron-hole contributions

- | Similar spectrum shape shifted by hole energy difference
- | Strengths are roughly proportional to the spectroscopic factors
- | Shape in each orbit should be the same



# Spectroscopic Factors of $^{89}\text{Y}$

- $1g9/2^-1$ : ~10

- $1f5/2^-1$ : ~3

**Ex=1.71 MeV**

- **Others:**  
**small/broad backgrounds**

TABLE I. Spectroscopic factors obtained for  $^{90}\text{Zr}(\text{p},\text{d})^{89}\text{Zr}$ .

$nlj$	$E_x$ (MeV)	$C^2S$	Simple shell model prediction	Average excitation energy $E_j$ (MeV)
$1g_{9/2}$	0.	9.6		0.18
	2.75	0.29		
	9.03	0.11	<u>10.00<sup>a</sup></u>	10
$2p_{1/2}$	0.60	1.2		1.23
	8.10	0.11	<u>1.31</u>	2
$2p_{3/2}$	1.09	2.1		1.95
	1.87	0.53		
	9.62	0.24	<u>2.87</u>	4
$1f_{5/2}$	1.46	3.5 <sup>b</sup>		3.27 <sup>c</sup>
	2.10	0.66		
	3.0	0.46		
	3.4–7.0	1.7		
	9.86	0.64	<u>6.96</u>	6
$1f_{7/2}$	4.1	0.82		10.3 <sup>c</sup>
	5.2	0.35		
	7.0–11.6	2.8		
	11.6–16.0	2.1		
	16.0–21.0	0.43	<u>6.50</u>	8
$1d_{3/2}$	16.0–21.0	1.2		

<sup>a</sup> $1g_{9/2}$  strength normalized to 10.0.

<sup>b</sup>Contains a small contribution (< 20%) from the 1.51 MeV ( $\frac{9}{2}^+$ ) state.

<sup>c</sup>The spectroscopic strength was assumed always to be at the center of the broad regions.

# Spectroscopic Factors of $^{51}\text{V}$

- $1\text{f}_{7/2}^{-1}$ : ~5  
 $\langle \text{Ex} \rangle = 0.56 \text{ MeV}$
- $2\text{s}_{1/2}^{-1}$ : ~1.2  
 $\text{Ex} = \sim 2.8 \text{ MeV}$
- $1\text{d}_{3/2}^{-1}$ : ~0.9  
 $\text{Ex} = \sim 3.2 \text{ MeV}$

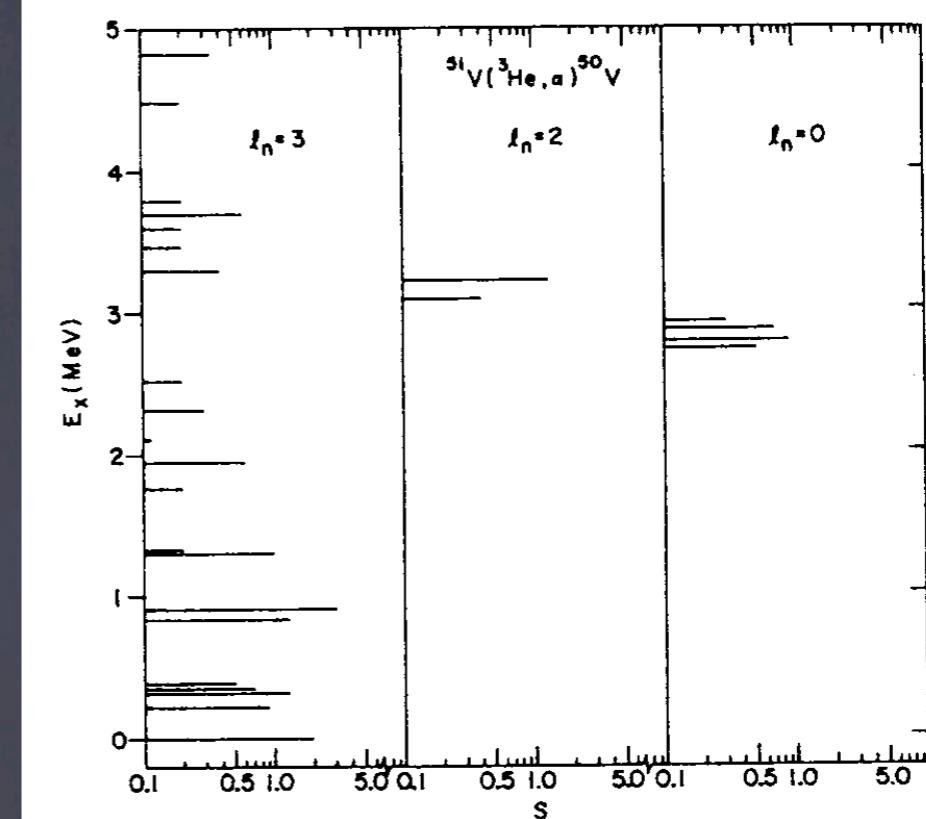
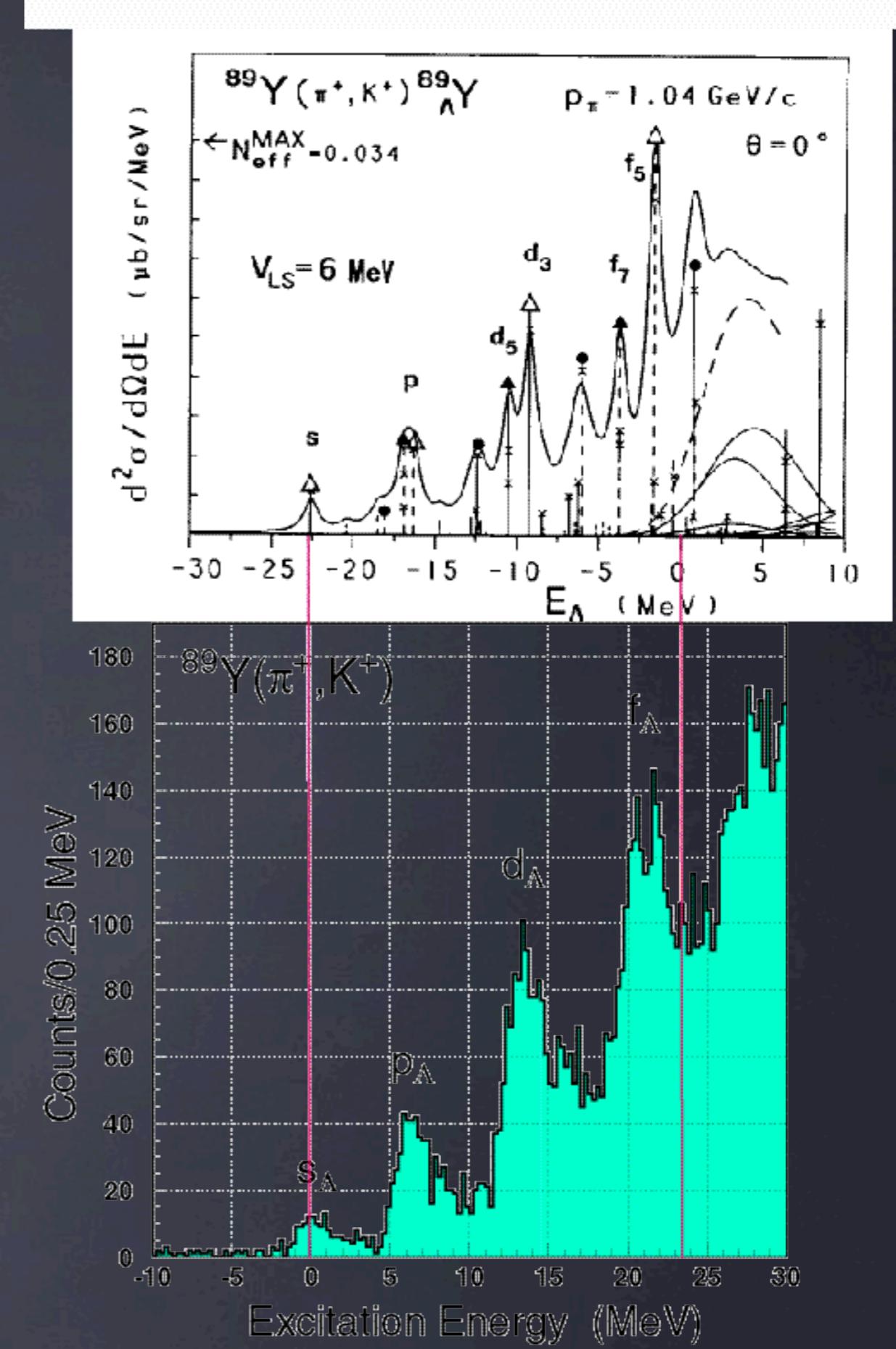


FIG. 13. The spectroscopic factors obtained from a DWBA analysis of the data on the  $^{51}\text{V}(^3\text{He}, \alpha)^{50}\text{V}$  reaction. Results for  $l_n = 3, 2$ , and 0 are indicated by the length of the bar at the appropriate excitation energy.

# Comparison with a calculation

- Overall peak structure: in good agreement, except peak intensity ratios.



# Spin-Orbit Splitting in Nuclei

	Normal nuclei	$\Lambda$ -hypernuclei
G-matrix	~40%	very small
Bruckner S.O. splitting		
Pauli rearrangement	~15%	X
Three-Body Force	~20%	X(?)
Other	~25%	

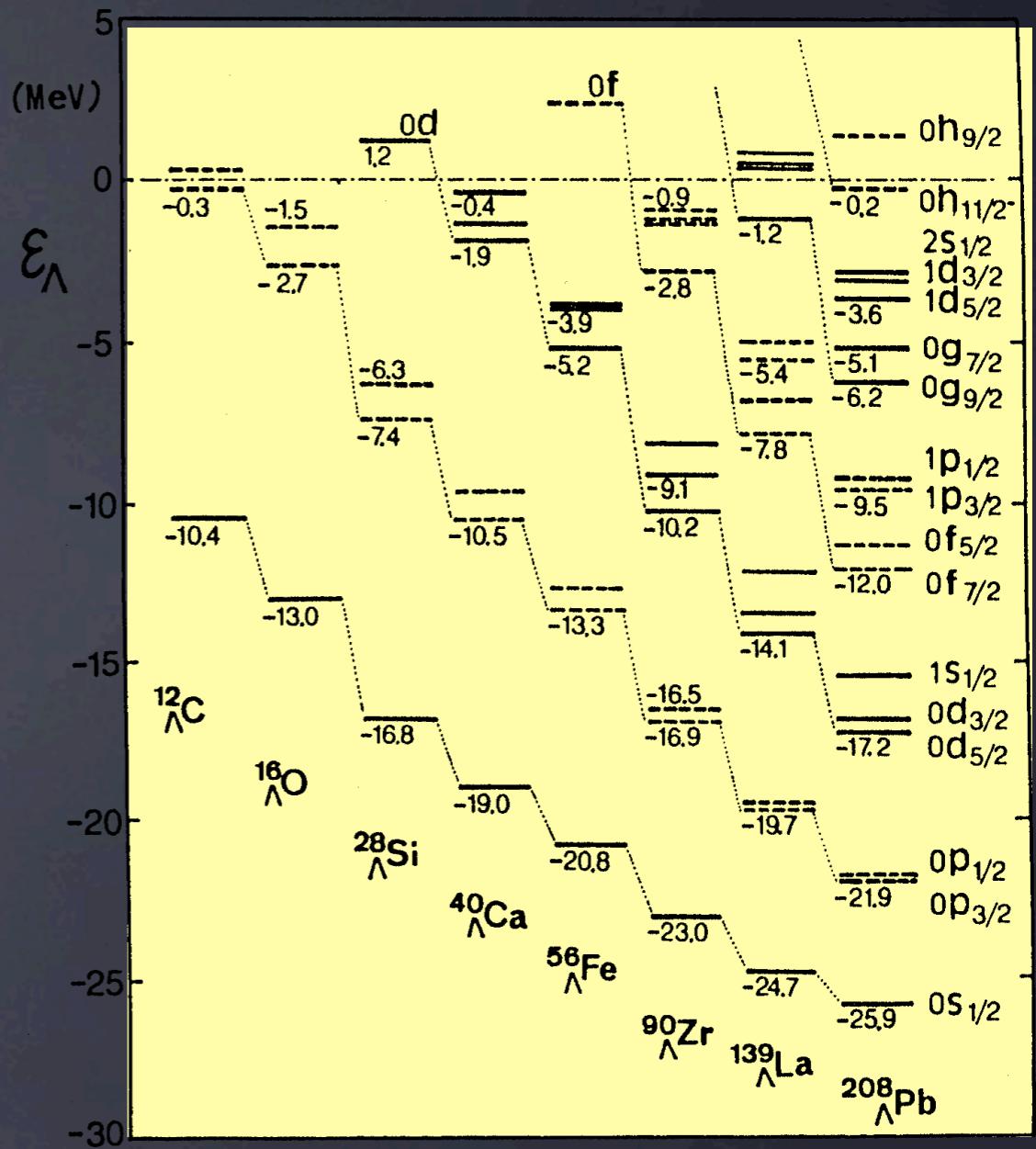
Y.Yamamoto

K.Anodo and H.Bando, PTP 66 (1981)  
227.

# Advantage in heavy hypernuclei

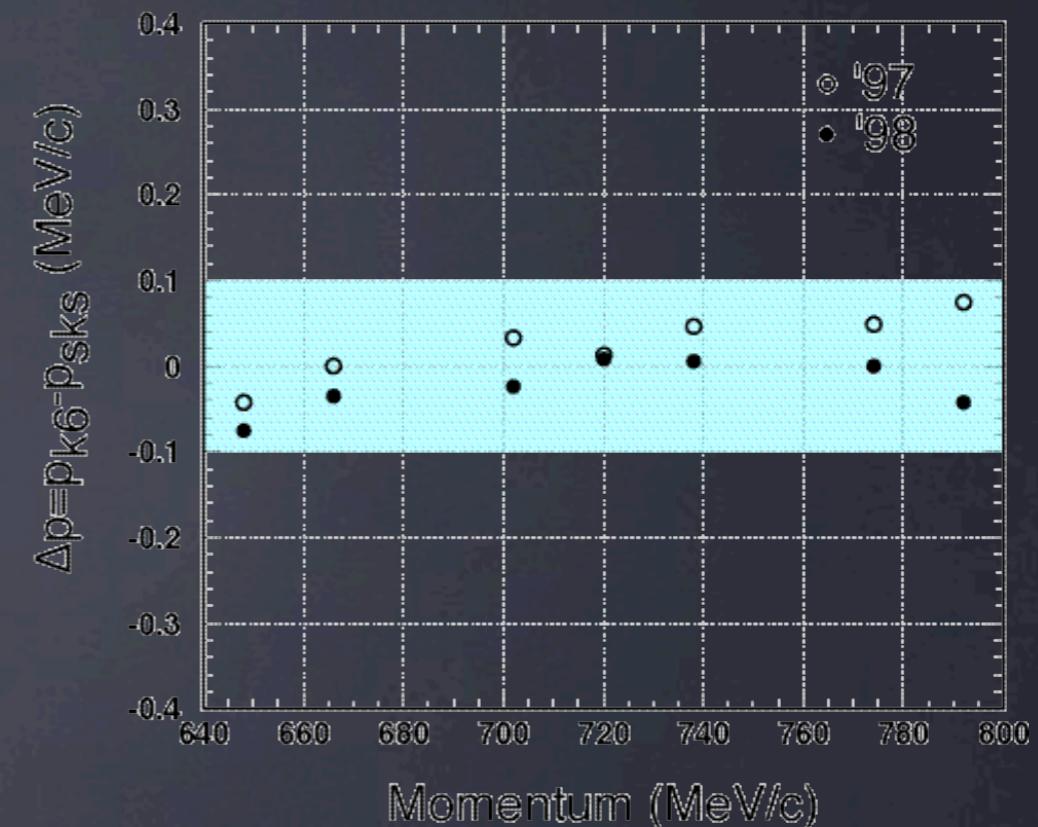


- $\Lambda$ -orbits with large angular momenta are bound.  
*Large LS splittings*  
 $\Delta E \propto V_{LS} (I+1/2) A^{-2/3}$
- Two split peaks are comparably excited in the  $(\pi^+, K^+)$  reaction.

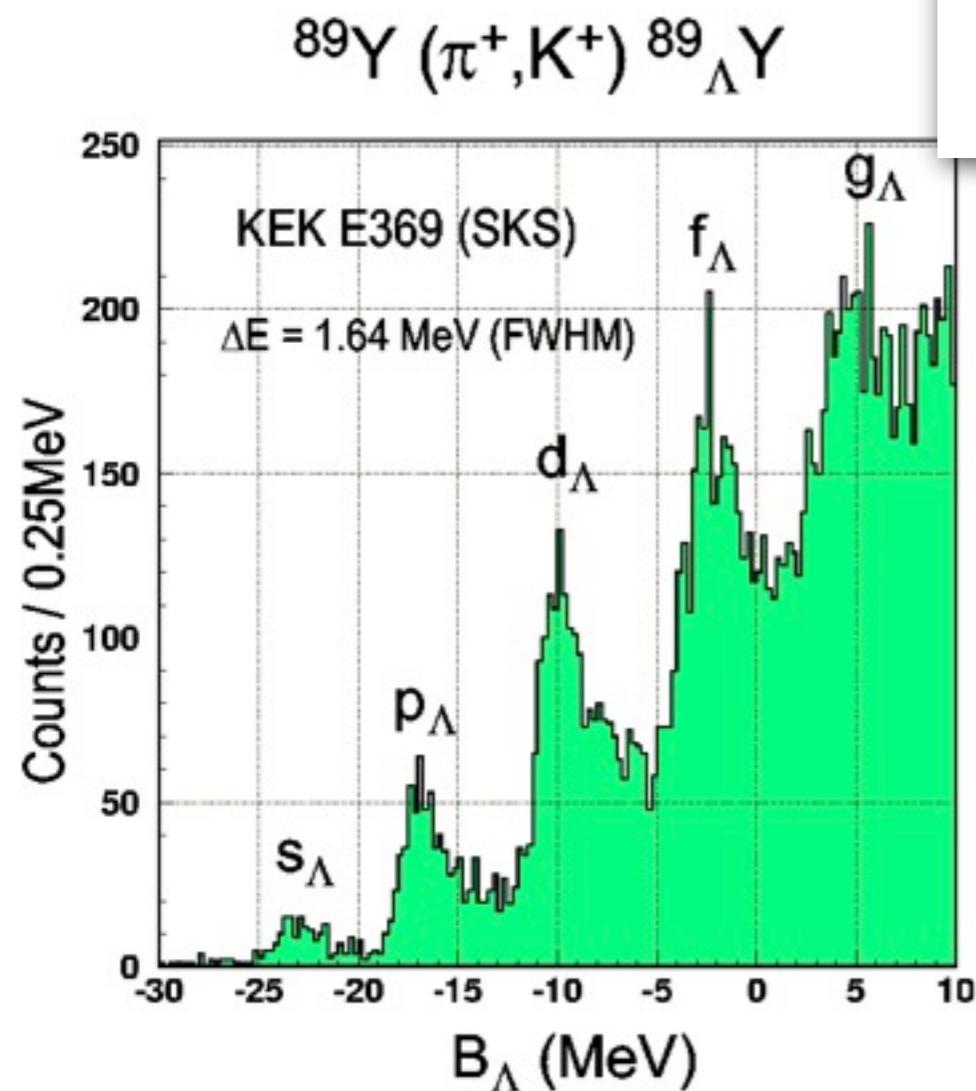


# Energy Scale: linearity and offset

- Beam-through Runs:  $720 \text{ MeV}/c \pm 10\%$ 
  - No target, SKS at 2.2 T
  - $< \pm 0.1 \text{ MeV}/c$
- Energy offset
  - $B(12 \Lambda C_{g.s.}) = 10.75 \pm 0.19 \text{ MeV}$ , in emulsion
  - Long-term fluctuation  $< 0.2 \text{ MeV}$
- Total systematic error  $< \sim \pm 0.2 \pm 0.19 \text{ MeV}$



## YN forces from hypernuclear structure



L Single particle states

-> potential depth

$$U_L = -30 \text{ MeV}$$

$$V_{LN} < V_{NN}$$



forbidden

LN spin-spin, spin-orbit, tensor forces ?  
LN  $\leftrightarrow$  SN coupling force?  
=> precise data necessary

